

THE CRITICAL POINT EQUATION ON A FOUR DIMENSIONAL WARPED PRODUCT MANIFOLD

SEUNGSU HWANG* AND JEONGWOOK CHANG

ABSTRACT. On a compact oriented n -dimensional manifold (M^n, g) , it has been conjectured that a metric g satisfying the critical point equation (2) should be Einstein. In this paper, we prove that if a manifold (M^4, g) is a 4-dimensional oriented compact warped product, then g can not be a solution of CPE with a non-zero solution function f .

1. Introduction

Let M be an n -dimensional compact orientable manifold. Denoting the space of all smooth Riemannian metrics on M by \mathcal{RM} , let \mathcal{M} be the quotient of \mathcal{RM} by the group of all diffeomorphisms of M . Then, for a given smooth structure $g \in \mathcal{M}$, its scalar curvature s_g is an element of the space of $C^\infty(M)$ functions, and the linearization of the scalar curvature is given by

$$s'_g(h) = -\Delta_g \operatorname{tr} h + \delta_g^* \delta_g h - g(h, r_g),$$

where Δ_g is the negative Laplacian of g , r_g its Ricci tensor, δ the divergence operator, and δ^* is the formal adjoint of δ . Also, the L^2 -adjoint operator s'^*_g of s'_g is given by

$$(1) \quad s'^*_g(f) = -g\Delta_g f + D_g df - fr_g$$

and the critical point equation, denoted *CPE hereafter*, is given by ([1])

$$(2) \quad z_g = s'^*_g(f),$$

where z_g is the traceless Ricci tensor, and f a function on M^n with vanishing mean value.

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For the case that the function f in CPE is trivial, it is clear that the metric g is Einstein. Therefore, all considerations in this paper are restricted to a *non-trivial function f only*. For the case that f is a smooth non-trivial function, the following statement has been conjectured ([1]):

CONJECTURE A. *If CPE holds for a non-trivial function f , then the metric g of the manifold M is Einstein.*

If this conjecture is true, it is known that (M, g) is isometric to a standard sphere S^n [10]. It turns out to be difficult to solve Conjecture A even with additional assumptions imposed on the metric. However, there are some partial answers to Conjecture A, such as those in [4] and [9] for example.

The purpose of the present paper is to prove the following Main Theorem as a partial (negative) answer to Conjecture A:

MAIN THEOREM. *Let (M, g) be a 4-dimensional oriented compact warped product given by $(M, g) = (B, \check{g}) \times_{\psi^2} (F, \hat{g})$ with $\psi > 0$. Then g can not be a solution of CPE with a non-zero solution function f .*

REMARK 1. Fisher and Marsden suggested the F-M conjecture in [3], stating that if a smooth function f satisfies $s_g^*(f) = 0$, then a solution metric g is isometric to the standard sphere. For the relationships between Conjecture A and F-M conjecture may be found in [5]. However, counter-examples of F-M conjecture were found (e.g., see [8]), mostly warped product manifolds. Therefore, it naturally arises a question to ask whether there exists a warped product metric which constitutes a counter-example of Conjecture A.

Contrary to the case of F-M conjecture, our Main Theorem shows that no 4-dimensional warped products satisfy CPE unless they are isometric to standard spheres. Combining this result with that of 3-dimensional case in [6], we may conclude that the answer to the question in the previous paragraph is no for $n \leq 4$.

It is natural to ask a similar question for $n \geq 5$. However, the difficulty for answering the question increases greatly, and is beyond our current understanding of the problem.

REMARK 2. For a 4-dimensional warped product manifold $(M, g) = (B, \check{g}) \times_{\psi^2} (F, \hat{g})$ with $\psi > 0$, it is clear that B is complete if M is complete. If we consider the case when the 1-dimensional base space B is incomplete with $\psi \geq 0$ and the fiber F is Einstein, we may conclude that if g is a solution metric of (3) with a non-zero solution function f ,

then (M, g) is isometric to the standard sphere S^4 (Proposition 11). It will be shown in Section 5.

In the present paper, our Main Theorem will be proved as follows. For a 4-dimensional warped product manifold $(M, g) = (B, \check{g}) \times_{\psi^2} (F, \hat{g})$, only the following 3 cases are possible with respect to the dimension of B and F :

Case 1. $\dim B = 1, \dim F = 3$

Case 2. $\dim B = 2, \dim F = 2$

Case 3. $\dim B = 3, \dim F = 1$.

For the case that B is complete, we prove a non-existence theorem (Main Theorem) in Section 3. For the case that B is not complete with $\dim B = 1$, a rigidity result is shown in Section 5.

2. Preliminaries

This section is a brief collection of notations and results, which are needed in our subsequent considerations.

Among the partial answers to the conjecture A given in section 1, the following three theorems, which hold on an n -dimensional manifold (M, g) , are needed in the next sections.

THEOREM 1. *Let g be a solution metric of CPE. If g is conformally flat, then (M, g) is isometric to a standard sphere [9].*

THEOREM 2. *Let (g, f) be a solution of CPE. If f always takes values greater than or equal to -1 , then (M, g) is isometric to a standard sphere [4].*

THEOREM 3. *Let g be a solution metric of CPE. If the metric g is Einstein, (M, g) is isometric to a standard sphere S^n [10].*

Furthermore, using (1) and (2) we have another useful representation of CPE on an n -dimensional manifold (M, g) , which may be written as

$$(3) \quad (1 + f)z_g = D_g df + \frac{s_g f}{n(n-1)}g.$$

Taking the trace of (3), we have $\Delta_g f = -\frac{s_g}{n-1}f$. Note that the scalar curvature s_g of the metric g satisfying CPE is assumed to be constant ([1]). Therefore we have $\int_M f = 0$, and hence f takes both positive and negative values.

Finally, on a 4-dimensional compact warped product manifold $(M, g) = (B, \hat{g}) \times_{\psi^2} (F, \hat{g})$, let

- \tilde{r}, \tilde{s} : the lifts to M of Ricci and scalar curvature of B , respectively
- \hat{r}, \hat{s} : the Ricci and scalar curvature of F , respectively
- X_i : a lifted horizontal orthonormal frame field, $i = 1, \dots, \dim B$
- U_j : a lifted vertical orthonormal frame field, $j = 1, \dots, \dim F$
- \mathcal{V} : the vertical distribution.

Then, the following two propositions hold on M for each of 3 cases mentioned in the last paragraph of section 1:

PROPOSITION 4. We have ([1])

Case 1. $\dim B = 1$ and $\dim F = 3$;

$$\begin{aligned} r(X, X) &= -\frac{3\psi''}{\psi} \\ r(U_i, U_j) &= \hat{r}(U_i, U_j) + \langle U_i, U_j \rangle \left(-\frac{\psi''}{\psi} - 2\frac{\psi'^2}{\psi^2} \right) \\ s &= \frac{\hat{s}}{\psi^2} - 6 \left(\frac{\psi''}{\psi} + \frac{\psi'^2}{\psi^2} \right) \end{aligned}$$

Case 2. $\dim B = 2$ and $\dim F = 2$;

$$\begin{aligned} r(X_i, X_j) &= \tilde{r}(\tilde{X}_i, \tilde{X}_j) - \frac{2}{\psi} \tilde{D}d\psi(\tilde{X}_i, \tilde{X}_j) \\ r(U_i, U_j) &= \hat{r}(U_i, U_j) + \langle U_i, U_j \rangle \left(-\frac{\tilde{\Delta}\psi}{\psi} - \frac{|d\psi|^2}{\psi^2} \right) \\ s &= \tilde{s} + \frac{\hat{s}}{\psi^2} - 4\frac{\tilde{\Delta}\psi}{\psi} - 2\frac{|d\psi|^2}{\psi^2} \end{aligned}$$

Case 3. $\dim B = 3$ and $\dim F = 1$;

$$\begin{aligned} r(X_i, X_j) &= \tilde{r}(\tilde{X}_i, \tilde{X}_j) - \frac{1}{\psi} \tilde{D}d\psi(\tilde{X}_i, \tilde{X}_j) \\ r(U, U) &= -\frac{\tilde{\Delta}\psi}{\psi} \\ s &= \tilde{s} - 2\frac{\tilde{\Delta}\psi}{\psi}. \end{aligned}$$

The next corollary follows from Proposition 4.

COROLLARY 5. If the scalar curvature s of M is constant, then the scalar curvature \hat{s} of F is constant.

Proof. Proposition 4 gives $\hat{s} = \psi^2 s + 6(\psi''\psi + \psi'^2)$ in the Case 1. Since \hat{s} is a function on F , and the right-hand side of this equation is a function on B , \hat{s} should be a constant function on F , i.e., F is of constant scalar curvature. In the Case 2, this corollary follows similarly, since we have $\hat{s} = \psi^2(s - \tilde{s}) + 4\psi\tilde{\Delta}\psi + 2|d\psi|^2$. Finally, we have $\hat{s} = 0$ in the Case 3. \square

3. Proof of main theorem

This section is devoted to the proof of our Main Theorem. Throughout this section we assume that B is complete with $\psi > 0$. The proof of Main Theorem follows directly from the following lemmas. For the Case 1 ($\dim B = 1$), Lemma 6 and 7 give the proof. Lemma 8 and 9 give the proof of the Case 2 ($\dim B = 2$). For the remaining Case 3 ($\dim B = 3$), the proof follows from Lemma 10.

LEMMA 6. *Let $(M, g) = (B, \check{g}) \times_{\psi^2} (F, \hat{g})$ with the 1-dimensional complete base B . Suppose that f is a function of B only. Then g can not be a solution of (3).*

Proof. First note that

$$(4) \quad \langle D_{U_i} df, U_i \rangle = \langle D_{U_i} \check{d}f, U_i \rangle = \frac{\psi' f'}{\psi},$$

where we used the fact that $D_{U_i} \check{d}f = \frac{\psi' \partial_i f}{\psi} U_i$ [1]. Now, from (3) we have

$$(1 + f) \left(\hat{r}(U_i, U_i) - \frac{\psi''}{\psi} - 2 \frac{\psi'^2}{\psi^2} - \frac{s}{4} \right) = \frac{\psi'}{\psi} f' + \frac{sf}{12}.$$

Then it is easy to see that $\hat{r}(U_1, U_1) = \hat{r}(U_2, U_2) = \hat{r}(U_3, U_3)$, or $\hat{r}(\widehat{U}_i, \widehat{U}_i) = \frac{s}{3}$ for all $1 \leq i \leq 3$, and $\hat{r}(U_i, U_j) = 0$ for any $i \neq j$. Thus F is an Einstein manifold. Since F is 3-dimensional, it follows that F should be of constant sectional curvature. In other words, F is isometric to S^3/Γ , with $\Gamma \subset SO(\pm)$. Hence we have $M = S^1 \times_{\psi^2} (S^3/\Gamma)$ with the metric g given by $g = dt^2 + \psi(t)^2 g_0$. This metric g is conformally flat, since g_0 is of constant curvature, c.f. [8]. Then, by Theorem 1, (M, g) is isometric to a standard sphere S^4 . This is clearly a contradiction, since there is no nonvanishing function ψ with $S^4 = S^1 \times_{\psi^2} (S^3/\Gamma)$. \square

LEMMA 7. *Let $(M, g) = (B, \check{g}) \times_{\psi^2} (F, \hat{g})$ with the 1-dimensional complete base B . Then g can not be a solution of (3).*

Proof. From (3), we have

$$(5) \quad 0 = (1 + f)z(X, U_i) = \langle D_X df, U_i \rangle = X \langle df, U_i \rangle - \langle df, D_X U_i \rangle$$

for a lifted horizontal vector field X . Note that $U_i = \frac{1}{\psi} \widehat{U}_i$, where \widehat{U}_i is a lift of vector field on F . Therefore

$$(6) \quad D_X U_i = D_X \left(\frac{1}{\psi} \widehat{U}_i \right) = -\frac{X(\psi)}{\psi^2} \widehat{U}_i + \frac{1}{\psi} D_X \widehat{U}_i = 0,$$

where we used the fact that $D_X \widehat{U}_i = \frac{X(\psi)}{\psi} \widehat{U}_i$ [11] in the last equality. Now, substitution of (6) into (5) gives

$$(7) \quad XU_i(f) = X\langle df, U_i \rangle = 0.$$

Therefore $U_i(f) = \frac{1}{\psi} \widehat{U}_i(f)$ is a function of F only, and it is easy to see that f can be written as

$$(8) \quad f = \psi b + c,$$

where b is a function on F and $c = c(t)$ is a function on B . Substituting (8) into $(1+f)z(X, X) = \langle D_X df, X \rangle + \frac{s}{12}$, we have

$$(9) \quad (1 + \psi b + c) \left(-3 \frac{\psi''}{\psi} - \frac{s}{4} \right) = \psi'' b + c'' + \frac{s}{12} (\psi b + c),$$

where we used the fact that, from $\check{d}f = b d\psi + dc$ and $D_X \widehat{d}f = \mathcal{V} D_X \widehat{d}f$,

$$\begin{aligned} \langle D_X df, X \rangle &= \langle D_X \check{d}f, X \rangle + \langle D_X \widehat{d}f, X \rangle = \langle D_X \check{d}f, X \rangle \\ &= b \langle D_X d\psi, X \rangle + \langle D_X dc, X \rangle = \psi'' b + c''. \end{aligned}$$

Thus, the equation (9) can be rewritten as

$$(10) \quad b \left(-4\psi'' - \frac{s}{3}\psi \right) = (1+c) \left(\frac{3\psi''}{\psi} + \frac{s}{4} \right) + c'' + \frac{s}{12}c.$$

Note that both $-4\psi'' - \frac{s}{3}\psi$ and the right-hand side are functions of B only, while b is a function of F . Thus, in order that the equation (10) holds for any t , either b is constant or $-4\psi'' - \frac{s}{3}\psi$ has to be zero. If b is constant, f is a function of B only, and so g can not be a solution of (3) in virtue of Lemma 6. Now, we may assume that

$$(11) \quad \psi'' + \frac{s}{12}\psi = 0.$$

Since B is complete, ψ has to be defined on the whole of \mathbb{R} . Moreover, since $B = S^1$, ψ has to be periodic. Therefore we may conclude that ψ is zero somewhere on B ; if $\psi \neq 0$,

$$0 = \int_B \left(\frac{\psi'}{\psi} \right)' = \int_B \frac{\psi''}{\psi} - \frac{\psi'^2}{\psi^2} = - \int_B \left(\frac{s}{12} + \frac{\psi'^2}{\psi^2} \right) < 0$$

which is a contradiction. Hence the given warped product metric g can not be a solution of (3). \square

LEMMA 8. Let $(M, g) = (B, \check{g}) \times_{\psi^2} (F, \hat{g})$ with the 2-dimensional complete base B . Suppose that f is a function of B only. Then g can not be a solution of (3).

A proof of Lemma 8 need another section and will be presented in the Section 4.

LEMMA 9. *Let $(M, g) = (B, \check{g}) \times_{\psi^2} (F, \hat{g})$ with the 2-dimensional complete base B . Then g can not be a solution of (3). In other words, Case 2 does not occur.*

Proof. In the beginning, the proof of Lemma 9 may go along the same line as the proof of Lemma 7, concluding that

$$(12) \quad f = \psi b + c,$$

where b is a function on F and c is a function on B . By Proposition 4 and $(1 + f) \sum_{i=1}^2 z(X_i, X_i) = \sum_{i=1}^2 \langle D_{X_i} df, X_i \rangle + \frac{sf}{6}$, we have

$$(13) \quad (1 + f)(\check{s} - \frac{2}{\psi} \check{\Delta} \psi - \frac{s}{2}) = \sum_{i=1}^2 \langle D_{X_i} df, X_i \rangle + \frac{sf}{6}.$$

Now, substitution (12) into (13) gives

$$(14) \quad (1 + \psi b + c)(\check{s} - \frac{2}{\psi} \check{\Delta} \psi - \frac{s}{2}) = b \check{\Delta} \psi + \check{\Delta} c + \frac{s}{6}(\psi b + c),$$

where we used the fact that

$$\begin{aligned} \sum_{i=1}^2 \langle D_{X_i} df, X_i \rangle &= \sum_{i=1}^2 (\langle D_{X_i} \check{d}f, X_i \rangle + \langle D_{X_i} \hat{d}f, X_i \rangle) = \sum_{i=1}^2 \langle D_{X_i} \check{d}f, X_i \rangle \\ &= b \sum_{i=1}^2 \langle D_{X_i} d\psi, X_i \rangle + \sum_{i=1}^2 \langle D_{X_i} dc, X_i \rangle = b \check{\Delta} \psi + \check{\Delta} c \end{aligned}$$

with $D_{X_i} \hat{d}f = \mathcal{V} D_{X_i} \check{d}f$. It is easy to see that (14) can be rewritten as

$$(15) \quad b(\psi \check{s} - 3 \check{\Delta} \psi - \frac{2s}{3} \psi) = \check{\Delta} c - (c + 1)(\check{s} - \frac{2}{\psi} \check{\Delta} \psi - \frac{s}{2}) + \frac{s}{6} c.$$

Note that both $\psi \check{s} - 3 \check{\Delta} \psi - \frac{2s}{3} \psi$ and the right-hand side are functions of B only, while b is a function of F only. Thus, in order that the equation (15) holds, either b is constant or $\psi \check{s} - 3 \check{\Delta} \psi - \frac{2s}{3} \psi$ has to be zero. If b is constant, f is a function of B only, and so g can not be a solution of (3) in virtue of Lemma 8. Thus we may assume that

$$(16) \quad \check{s} = \frac{3 \check{\Delta} \psi}{\psi} + \frac{2}{3} s$$

and

$$(17) \quad \check{\Delta} c - (c + 1)(\check{s} - \frac{2}{\psi} \check{\Delta} \psi - \frac{s}{2}) + \frac{s}{6} c = 0.$$

Substitution of (16) into (17) gives

$$(18) \quad \psi \check{\Delta} c - (c + 1) \check{\Delta} \psi - \frac{s}{6} \psi = 0.$$

Now, integrating both sides of (18) over B , we have

$$\frac{s}{6} \int_B \psi = 0,$$

where we used the fact that $\int_B \psi \check{\Delta} c = \int_B c \check{\Delta} \psi$. This equation together with the condition that $\psi > 0$ leads to a contradiction. \square

LEMMA 10. *Let $(M, g) = (B, \check{g}) \times_{\psi^2} (F, \hat{g})$ with the 3-dimensional complete base B . Then g can not be a solution of (3). In other words, Case 3 does not occur.*

Proof. Suppose that the given warped product metric is a solution of CPE. From the definition of the Laplacian Δ , we have

$$(19) \quad \langle D_U df, U \rangle = \Delta f - \check{\Delta} f = -\frac{s}{3} f - \check{\Delta} f$$

and from CPE we also have

$$(20) \quad (1 + f) \left(-\frac{\check{\Delta} \psi}{\psi} - \frac{s}{4} \right) = \langle D_U df, U \rangle + \frac{s}{12} f.$$

Hence, the following relation holds by combining (19) and (20):

$$(1 + f) \left(-\frac{\check{\Delta} \psi}{\psi} - \frac{s}{4} \right) = -\check{\Delta} f - \frac{s}{4} f$$

which is equivalent to

$$(21) \quad (1 + f) \check{\Delta} \psi + \frac{s}{4} \psi = \psi \check{\Delta} f.$$

Now, integrating both sides of (21) over B , we have

$$\int_B f \check{\Delta} \psi + \frac{s}{4} \int_B \psi = \int_B \psi \check{\Delta} f = \int_B f \check{\Delta} \psi$$

which implies that

$$\int_B \psi = 0.$$

This equation together with the condition that $\psi > 0$ leads to a contradiction. \square

4. Proof of Lemma 8

This section is devoted to the proof of Lemma 8. Suppose that g is a solution of (3), namely CPE, and f is a function of B only. In the following two contentions, we shall prove that this assumption leads to a contradiction. We first consider the following two cases according to the values of f . In the first case that f always takes values greater than or equal to -1 , the proof of the Lemma is completed since (M, g) is isometric to S^4 in virtue of Theorem 2. In the second case that f takes a value less than -1 , this Lemma is proved in a series of two contentions under present conditions¹. In this case, there exists a non-empty set $H = \{x \in B \mid f(x) = -1\}$. After investigating the analytic properties of the tensor $D_X X$ for a tangent vector field X to H in Contention 1, we prove in Contention 2 that the present conditions give a contradiction. Since this contradiction is obtained from the assumption that g is a solution of CPE, we may conclude that g can not be a solution of CPE, proving our Lemma.

CONTENTION 1. Under present conditions, we have $D_X X = -\frac{s}{12}df$ on H , where $H = \{x \in B \mid f(x) = -1\}$, $X \in TB$ is a tangent vector field to H , and $N \in TB$ is a normal vector field to H .

Proof. We may assume that H is non-empty; otherwise g is Einstein by Theorem 2. Also, we note in [4] that a point of H , which is a critical point of f , is a non-degenerate local minimum point of f , and that such non-degenerate critical points are isolated. Therefore, H is a set consisting of finite critical points of f , or hypersurfaces of M , or union of both.

Putting $W = |df|^2$, it was proved in [4] that W is constant in each component of H and does not vanish on H . Therefore, in a small tubular neighborhood of H , we may take orthonormal frame fields $\{X, N\}$, where $N = \frac{df}{W^{1/2}}$ on H . From (3), we have

$$(22) \quad Ddf = \frac{s}{12}g$$

on H . Then, it follows that on H we have

$$\frac{s}{12} = \langle D_X df, X \rangle = X \langle df, X \rangle - \langle df, D_X X \rangle = -\langle df, D_X X \rangle.$$

¹In the proof of Contention 1 and 2, we assume that the metric g of warped product $(M, g) = (B, \hat{g}) \times_{\psi^2} (F, \hat{g})$ is a solution of (3). We also suppose that $\dim B = 2$ and f takes a value less than -1 . Hereafter, this situation will be described by the words “present conditions”.

Hence, $D_X X = -\frac{s}{12} \frac{1}{W^{\frac{1}{2}}} N$ on H . □

CONTENTION 2. Under present conditions, we have a contradiction on H .

Proof. In virtue of Proposition 4, (3) may be reduced to

$$(23) \quad (1 + f) \left(\check{s} - \frac{2}{\psi} \check{\Delta} \psi - \frac{s}{2} \right) = \check{\Delta} f + \frac{sf}{6}$$

$$(24) \quad (1 + f) \left(\frac{\hat{s}}{\psi^2} - 2 \frac{\check{\Delta} \psi}{\psi} - 2 \frac{|d\psi|^2}{\psi^2} - \frac{s}{2} \right) = 2 \frac{\langle d\psi, df \rangle}{\psi} + \frac{sf}{6}$$

since

$$(25) \quad \sum_{i=1}^2 z(X_i, X_i) = \check{s} - \frac{2}{\psi} \check{\Delta} \psi - \frac{s}{2}$$

$$(26) \quad \sum_{i=1}^2 z(U_i, U_i) = \frac{\hat{s}}{\psi^2} - 2 \frac{\check{\Delta} \psi}{\psi} - 2 \frac{|d\psi|^2}{\psi^2} - \frac{s}{2}$$

and $\check{g}(d\psi, df) = \langle d\psi, df \rangle$. Hence, using (23) and (24) we have

$$(27) \quad \check{\Delta} f = \frac{s}{6}$$

$$(28) \quad \frac{\langle d\psi, df \rangle}{\psi} = \frac{s}{12}$$

on H . The relation $\sum_{i=1}^2 z(X_i, X_i) + \sum_{i=1}^2 z(U_i, U_i) = 0$ gives

$$(29) \quad (1 + f) \left(-\check{s} + \frac{2}{\psi} \check{\Delta} \psi + \frac{s}{2} \right) = 2 \frac{\langle d\psi, df \rangle}{\psi} + \frac{sf}{6}.$$

Taking the Lie derivative of (29) with respect to df on H , we have

$$\begin{aligned} & W \left(-\check{s} + \frac{2}{\psi} \check{\Delta} \psi + \frac{s}{2} \right) \\ &= 2 \frac{\langle D_{df} d\psi, df \rangle}{\psi} + 2 \frac{\langle d\psi, D_{df} df \rangle}{\psi} - 2 \frac{\langle d\psi, df \rangle^2}{\psi^2} + \frac{s}{6} W \\ &= 2 \frac{\langle D_N d\psi, N \rangle}{\psi} W + \frac{s}{6} W, \end{aligned}$$

where

$$2 \frac{\langle d\psi, D_{df} df \rangle}{\psi} = \frac{s}{6\psi} \langle df, d\psi \rangle = \frac{s^2}{72} = 2 \frac{\langle d\psi, df \rangle^2}{\psi^2}$$

in virtue of (22) and (28). Therefore

$$(30) \quad -\check{s} + \frac{2}{\psi} \langle D_X d\psi, X \rangle + \frac{s}{3} = 0,$$

where we used the fact that

$$(31) \quad \check{\Delta}\psi = \langle D_X d\psi, X \rangle + \langle D_N d\psi, N \rangle.$$

On the other hand, in order to calculate $z(X, X)$ we take the Lie derivative of (3) with respect to df on H . Then

$$(32) \quad \begin{aligned} Wz(X, X) &= \langle D_{df} D_X df, X \rangle + \langle D_X df, D_{df} X \rangle + \frac{s}{12} W \\ &= W^{\frac{1}{2}} \langle D_N D_X df, X \rangle + |D_X df|^2 + \frac{s}{12} W \\ &= W^{\frac{1}{2}} \langle D_N D_X df, X \rangle + \frac{s^2}{144} + \frac{s}{12} W \end{aligned}$$

on H , where we used the fact that W is constant on H , $D_{df} X = W^{-\frac{1}{2}} D_N X = W^{-\frac{1}{2}} D_X N = D_X df$, and $D_X df = \frac{s}{12} X$ on H . However, the relation $D_N D_X df = D_X D_N df + R(X, N)df$ gives

$$(33) \quad \begin{aligned} \langle D_N D_X df, X \rangle &= \langle D_X D_N df, X \rangle + \langle R(X, N)df, X \rangle \\ &= X \langle D_N df, X \rangle - \langle D_N df, D_X X \rangle - W^{\frac{1}{2}} K(X, N) \\ &= \frac{s}{12W^{\frac{1}{2}}} \langle D_N df, N \rangle - W^{\frac{1}{2}} \check{r}(N, N) \\ &= \frac{s^2}{144W^{\frac{1}{2}}} - W^{\frac{1}{2}} \check{r}(N, N), \end{aligned}$$

where we used the fact that $\langle D_N df, X \rangle = 0$, $D_X X = -\frac{s}{12} \frac{1}{W^{1/2}} N$ on H , (22), and $\check{r}(N, N) = K(X, N)$. Now substituting (33) into (32), we have

$$(34) \quad z(X, X) = \frac{s^2}{72W} - \check{r}(N, N) + \frac{s}{12}.$$

Note that from Proposition 4 we have

$$(35) \quad z(X, X) = r(X, X) - \frac{s}{4} = \check{r}(X, X) - 2 \frac{\langle D_X d\psi, X \rangle}{\psi} - \frac{s}{4}.$$

Substituting (34) into (35) on H we have

$$(36) \quad \check{s} - 2 \frac{\langle D_X d\psi, X \rangle}{\psi} - \frac{s}{3} - \frac{s^2}{72W} = 0.$$

The equation (36) clearly contradicts the equation (30). This contradiction comes from the assumption that g is a solution of CPE. Hence, g can not be a solution of CPE, proving our Lemma 8. \square

5. A rigidity result

This section is devoted to the proof of the following rigidity result.

PROPOSITION 11. *Let M be a compact oriented 4-dimensional manifold containing an open dense subset U which is (for the induced metric) a warped product on a 1-dimensional basis with Einstein fibre F . If g is a solution metric of (3), then (M, g) is isometric to the standard sphere S^4 .*

We take U as large as possible. If the base B is complete, then M is a warped product and Lemma 7 applies. If not, B is an open interval $(0, a)$ with $U = (0, a) \times F$. It is easy to see that M is the quotient of $[0, a] \times F$. Note that the following relation holds in order for M to be a complete manifold:

$$(37) \quad \psi(0) = \psi(a) = 0.$$

Note also that the following relation holds in order for M to be smooth (c.f. [7]):

$$(38) \quad \psi'(0) = -\psi'(a) = \sqrt{\frac{\hat{s}}{6}}.$$

Thus (37) and (38) become the initial conditions for the warping function ψ . On the other hand, in virtue of Proposition 4, ψ satisfies

$$(39) \quad 6\psi''\psi + 6\psi'^2 + \psi^2s - \hat{s} = 0$$

with constants s and \hat{s} . For a solution ψ of (39), we observe that $\psi''(0) = 0$, since the differentiation of (39) gives $3\psi''\psi + 9\psi''\psi' + s\psi\psi' = 0$. It is easy to see that $\psi_0 = \sqrt{\frac{2\hat{s}}{s}} \sin \sqrt{\frac{s}{12}}t$ with $a = \pi\sqrt{\frac{s}{12}}$ is a solution of the ordinary differential equation (39) with the initial conditions (37) and (38). The following lemma shows that the solution of (39) is unique.

LEMMA 12. *ψ_0 is the unique solution of (39) under the conditions (37) and (38). Therefore, $(M, g) = [0, a] \times_{\psi_0^2} S^3$.*

Proof. Let ψ be another solution of (39) with the initial conditions (37) and (38), and let $F = \frac{\psi}{\psi_0}$. It is easy to see that F is well-defined on $[0, a]$ with $F(0) = 1$, since $\psi_0(0) = \psi(0) = 0$ and $\psi'_0(0) = \psi'(0) = \sqrt{\frac{\hat{s}}{6}}$. Now, in order to prove our Lemma, we claim that $F \equiv 1$, implying that $\psi = \psi_0$. First we observe that $F'(0) = 0$ since $\psi'' = \psi''_0 F + 2\psi'_0 F' + \psi_0 F''$

and $\psi_0''(a) = \psi''(0) = 0$. Now, if we substitute $\psi = \psi_0 F$ into (39), F satisfies

$$(40) \quad F'' = -\frac{1}{\psi_0^2 F} \left(4\psi_0 \psi_0' F F' + \psi_0^2 F'^2 + \frac{1}{6} \hat{s}(F^2 - 1) \right),$$

where we used the fact that ψ_0 satisfies (39). Define a function $\xi : \mathbf{R} \times \mathbf{R} \times \mathbf{R} \rightarrow \mathbf{R}$ by

$$\xi(t, y_1, y_2) = -\frac{1}{\psi_0^2 y_1} \left(4\psi_0 \psi_0' y_1 y_2 + \psi_0^2 F'^2 + \frac{1}{6} \hat{s}(F^2 - 1) \right),$$

where ξ is obtained simply by replacing F and F' in (40) by y_1 and y_2 respectively. In order to prove the uniqueness of solutions of (40) satisfying $F(0) = 1$ and $F'(0) = 0$, it suffices to show that $\xi(t, y_1, y_2)$ is continuous with respect to t and is Lipschitz with respect to y_1 and y_2 in the rectangle $R = [0, a] \times [r, R_1] \times [-R_2, R_2]$ for all $0 < r < 1 < R_1$ and $R_2 > 0$. First, it is easy to see that ξ is continuous for t by letting $\xi(0, y_1(0), y_2(0)) = 0$. Secondly, ξ is Lipschitz, since ξ is smooth with respect to y_1, y_2 in R , and $|\xi(t, y_1(t), y_2(t)) - \xi(t, \bar{y}_1(t), \bar{y}_2(t))| < M(|y_1(t) - \bar{y}_1(t)| + |y_2(t) - \bar{y}_2(t)|)$ for some $M > 0$. Therefore, the ordinary differential equation (40) has the unique solution $F \equiv 1$, proving our claim. \square

It is well known (see [1], Corollary 9.107) that the warped product $M = B \times_{\psi^2} F$ is Einstein if and only if \hat{g}, \hat{g}, ψ satisfy that (F, \hat{g}) is Einstein, and

$$(41) \quad -\frac{\psi''}{\psi} - 3\frac{\psi'^2}{\psi^2} + \frac{\hat{s}}{3\psi^2} = \frac{s}{4}$$

and

$$(42) \quad -3\frac{\psi''}{\psi} = \frac{s}{4}.$$

It is an easy exercise for $\psi = \psi_0$ to satisfies (41), (42). Also note that F is Einstein from the assumption. Therefore we may conclude that the warped product U (or M) is Einstein. However, in virtue of Theorem 3, M must be isometric to a standard sphere.

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SEUNGSU HWANG, DEPARTMENT OF MATHEMATICS, CHUNG-ANG UNIVERSITY, SEOUL 156-756, KOREA

E-mail: seungsu@cau.ac.kr

JEONGWOOK CHANG, DEPARTMENT OF MATHEMATICS, SCHOOL OF MATHEMATICS, INFORMATICS AND STATISTICS, KUNSAN NATIONAL UNIVERSITY, KUNSAN 573-701, KOREA

E-mail: jchang@kunsan.ac.kr