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# Effects of external field and thickness asymmetry on dynamic magnetization switching of exchange-coupled trilayer 

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Exchange-coupled trilayers, which are composed of two ferromagnetic layers separated by a thin non-magnetic spacer, have received considerable attention in recent years because of their potential advantages in ultra high density magnetic random access (MRAM) applications [1-3]. Notable advantages include an enhanced thermal stability when the exchange-coupled trilayers use as a free layer structure due to increased magnetic volume over a single magnetic free layer [3], and an reduced critical current density for current induced magnetization switching (CIMS) due to a reduction of the effective magnetization [3]. In this work, the dynamic magnetization switching behavior of exchange-coupled trilayers was investigated by micromagnetic simulation. A particular emphasis was places on the effects of the external field and the thickness asymmetry of the two ferromagnetic layers on the dynamic switching process.

Micromagnetic simulation was performed using a commercial program code of MicroMagus (Sura Instruments, Germany) based on the Landau-Lifschtz equation. The exchange-coupled trilayer is composed of two ferromagnetic layers separated by a very thin non-magnetic spacer. An elliptical magnetic thin film with lateral dimensions of 200 $\mathrm{nm} \times 100 \mathrm{~nm}$ was considered. The total thickness of the two magnetic layers $\left(t_{1}+t_{2}\right)$ was fixed at 4 nm , however, the thickness of each magnetic layer was varied. The thickness of the spacer was fixed at 0.6 nm . The magnetic parameters used were; a saturation magnetization $\left(M_{s}\right)$ of $820 \mathrm{emu} / \mathrm{cc}$, an exchange constant $\left(A_{e x}\right)$ of $1.0 \times 10^{-6} \mathrm{erg} / \mathrm{cm}$, an induced anisotropy constant $\left(K_{u}\right)$ of $4100 \mathrm{erg} / \mathrm{cm}^{3}$, an interlayer exchange constant $\left(J_{e x}\right)$ of $-0.05 \mathrm{erg} / \mathrm{cm}^{2}$, and a damping constant (a) of 0.003 . Subpicosecond time steps were used to simulate dynamics of the magnetization reversal.

Figures (a) and (b) show the results for the dynamic magnetization switching behavior for a sample with a thickness asymmetry, $\Delta t=0.4 \mathrm{~nm}\left(t_{1}=2.2 \mathrm{~nm}\right.$ and $\left.t_{2}=1.8 \mathrm{~nm}\right)$. Fig. (a) shows the time evolution of the normalized magnetization components in the length (x) direction, $M_{\mathrm{x} 1}$ (for thick layer) and $M_{\mathrm{x} 2}$ (for thin layer). The magnetizations of the two ferromagnetic layers are initially in the antiparallel state $M_{1}$ and $M_{2}$ point to the $+x$ and $-x$ directions, respectively. Two different characteristic times can be defined: the coherent rotation time $\tau_{R}$ and the oscillation time $\tau_{\mathrm{o}}$. The magnetization $\mathrm{M}_{\mathrm{x}}$ initially does not react to the external field at $\tau \leq 5 \mathrm{~ns}$, after which, both $M_{\mathrm{x} 1}$ and $M_{\mathrm{x} 2}$ start to switching abruptly. Following this magnetization reversal, both magnetizations oscillate with exactly same frequency as shown in insert for a long period time. The oscillation frequency is quite small $(1.05 \mathrm{GHz}$ at the switching field of 235 Oe ) and decreases as the external field increases as shown in Fig. (b). Figs. (c) and (d) show two characteristic times of $\tau_{\mathrm{R}}$ and $\tau_{\mathrm{O}}$ as a function of the external field at several values of $\Delta t$. As expected, both times of $\tau_{R}$ and $\tau_{O}$ decrease as the external field increases indicating the fast switching at high external field. The switching time was also decreased as the thickness asymmetry was increased.


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## References

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