## Fast current-induced motion of a transverse domain wall induced by interfacial Dzyaloshinshkii-Moriya interaction

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Recently, ferromagnet/heavy metal bilayers attract considerable attention because they allow us to investigate various spin-orbit coupling effects combined with spin transport and magnetization dynamics. A representative example is the spin-orbit spin transfer torque (SOT) [1] that enables very fast current-induced magnetization switching even without the second ferromagnetic layer [2]. Another interesting magnetic property is the Dzyaloshinskii-Moriya interaction (DMI), emerging when all of spin-orbit coupling, exchange interaction, and inversion asymmetry are present. The DMI is the antisymmetric component of the exchange interaction [3,4], which favors non-collinear magnetic textures. Effects of the interfacial DMI on the domain wall motion in perpendicularly magnetized nanowires have been extensively studied [5]. However, the effect of the interfacial DMI on transverse domain wall motion has not been studied yet. In this work, we investigate the effect of the interfacial DMI on static and dynamic properties of a transverse domain wall.

Based on the Euler-Lagrange equation, the equilibrium profile of a transverse domain wall is determined as,

$$\theta(x) = 2 \tan^{-1} \left( \exp \frac{x - q}{\lambda} \right), \tag{1}$$

$$\varphi(x) \equiv \varphi_0 - \chi \sec h \frac{(x-q)}{\lambda}, \qquad (2)$$

where  $\hat{\mathbf{m}} = (\cos\theta, \cos\varphi\sin\theta, \sin\varphi\sin\theta)$ ,  $\theta$  is the polar angle from *x*-axis and  $\varphi$  is the azimuthal angle from *y*-axis, *q* is the domain wall center,  $\varphi_0$  is the domain wall tilt angle,  $\lambda$  is the domain wall width,  $\chi = \delta/\lambda$ ,  $\delta = D/K_d$ , and  $K_d$  is the hard-axis anisotropy energy density. From above equations, one can find that the domain wall distortion appears in cases with finite *D*. We found that the numerically calculated ones are in good agreement with the theoretically predicted ones.

By using the procedure developed by Thiele [6], we derive the equations of motion for the two collective coordinates of a transverse domain wall as following.

$$\alpha \frac{\dot{q}}{\lambda} + \dot{\varphi}_{0} = -\beta \frac{b_{J}}{\lambda} + \gamma c_{J} \left( \frac{\pi}{2} \sin \varphi_{0} - \frac{4}{3} \chi \cos \varphi_{0} \right),$$

$$\frac{\dot{q}}{\lambda} - \alpha \dot{\varphi}_{0} = -\frac{b_{J}}{\lambda} + \gamma \frac{D}{\lambda M_{s}} \left( \frac{\pi}{2} \cos \varphi_{0} + \frac{4}{3} \chi \sin \varphi_{0} \right) + \gamma \frac{K_{d}}{M_{s}} \left( \sin 2\varphi_{0} - \frac{\pi}{2} \chi \cos 2\varphi_{0} \right),$$
(3)

where  $\dot{O} = dO/dt$ ,  $\gamma$  is the gyromagnetic ratio,  $\alpha$  is the damping constant,  $\beta$  is the nonadiabaticity,  $b_J$  is the magnitude of spin-transfer torque,  $c_J$  is the magnitude of SOT and  $M_S$  is the saturation magnetization. Using the small angle approximation, we obtain the domain wall velocity  $v_{DW}$  at the steady state, given as

$$v_{DW} = \frac{-b_J(\pi c_J - 2\beta H_d) + (8/3)\gamma \chi \lambda c_J H_d}{\pi c_J - 2\alpha H_d},$$
(5)

where  $H_d = 2K_d/M_s$ .

We perform semi two-dimensional micromagnetic simulation for confirming a validity of Eq. (5). As shown in the figure 1, when D = 0,  $v_{DW}$  is small, but it increases rapidly with D. With reasonable material parameters, the  $v_{DW}$  reaches 400 m/s at the current density of  $9 \times 10^6$  A/cm<sup>2</sup> which has never been achieved for a transverse domain wall without DMI [7]. This high  $v_{DW}$  can be explained by a DMI-induced domain wall distortion. It generates non-zero SOT in the z-direction, which tilts the domain wall and enhances the domain wall motion.

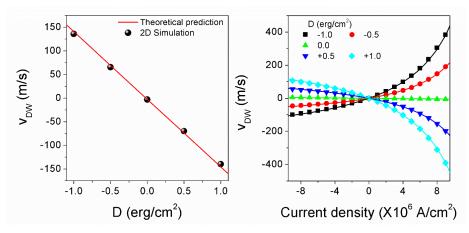


Fig. 1. (a) Effect of the interfacial DMI (*D*) on the domain wall velocity for 1-D micromagnetic simulation results (dots) and theoretical prediction (line). (b)  $v_{DW}$  as a function of the current density at various DMI values.

## References

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