# A Further Study on the NMR Chemical Shift for a 3d ${ }^{1}$ System in a Strong Crystal Field Environment of Octahedral Symmetry with a Threefold Axis of Quantization 

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The hamiltonian representing the various interaction in a strong crystal field environment of octahedral symmetry may be expressed as

$$
\begin{equation*}
H=-\frac{\hbar^{2}}{2 m} \nabla^{2}-\frac{Z e^{2}}{4 \pi e_{0} r}+V(\underset{\sim}{r})+\zeta \underline{\sim} . S+\mu_{B}(\underline{l}+2 \underset{\sim}{S})+H_{m} \tag{1}
\end{equation*}
$$

where

$$
\begin{equation*}
\left.H_{N N}=\frac{\mu_{4}}{4 \pi} g_{N} \mu_{B} \mu_{N}\left\{\frac{2 t_{N} \cdot I}{r_{N} 3}+g_{8}\left(\frac{3\left(r_{N} \cdot S\right.}{r_{N}}\right) \underline{r}_{N} \cdot I \quad \frac{S \cdot I}{r_{N}{ }^{2}}\right]\right\} \tag{2}
\end{equation*}
$$

and when a threefold axis is chosen as the axis of quantization

$$
\begin{equation*}
V(r)=a_{4}\left\{\frac{1}{3} \sqrt{\frac{7}{3}} Y_{4}(\theta, \phi)-\frac{1}{3} \sqrt{\frac{10}{3}}\left\{Y_{4-1}(\theta, \phi)-Y_{42}(\theta, \phi)\right]\right\} \tag{3}
\end{equation*}
$$

Here $r$ and $r_{\underline{N}}$ are the electron radius vectors about the electron-bearing atom and the nucleus with nuclear spin angular momentum $I$, respectively. The quantity $B$ is the applied magnetic field, $\mathrm{V}(r)$ is the crystal field potential of octahedral symmetry when a threefold axis is chosen as the axis of quantizaiton and $\mathrm{a}_{4}{ }^{2)}$ is the required crystal field parameter for the 3d electron system. The other symbols have their usual meaning. In this paper $g_{s}$, the free electron $g$ value is taken to be equal to exactly 2 .

When a threefold axis is chosen as the axis of quantization, the axial wave functions with $\mathrm{t}_{2}$ symmetry may be expressed, in $\mid \mathrm{m}_{ \pm}>$notation, as ${ }^{3}$

$$
\begin{align*}
& \phi_{0}^{ \pm}-\left|3 d_{2} 2^{ \pm}\right\rangle \\
& \phi_{1}^{ \pm}=\sqrt{\frac{2}{3}}\left|2^{ \pm}\right\rangle-\sqrt{\frac{1}{3}}\left|-1^{+}\right\rangle \\
& \phi_{5}^{ \pm}=\sqrt{\frac{2}{3}}\left|-2^{ \pm}\right\rangle+\sqrt{\frac{1}{3}}\left|1^{ \pm}\right\rangle \tag{4}
\end{align*}
$$

To determine the NMR chemical shift arising from the 3 d orbital angular momentum and the 3d ellectron spin dipolarnuclear spin dipolar-nuclear spin angular momentum interaction for a $3 \mathrm{~d}^{1}$ system in a strong crystal field environment of octahedral symmetry when a threefold axis is chosen as the axis of quantization, the principal values of the NMR screening tensor $\alpha_{x}, \alpha_{y y}$ and $\alpha_{u s}$, are determined by considering the
magnetic field interaction as pararell to the $x, y$ and $z$ axes and averaged assuming a Boltzmann distribution. It follows that the contribution to the NMR chemical shift, $\Delta \mathrm{B}$, is given by

$$
\Delta B=\frac{1}{3} B\left(\sigma_{x x}+\sigma_{y y}+\sigma_{x s}\right)
$$

where

$$
\begin{equation*}
\sigma_{a \sigma}=\left(\frac{\left.\partial^{2}<H_{N N}\right\rangle}{\partial \mu_{a} \partial B_{a}}\right)_{\sim}^{\mu}=B-0 \tag{5}
\end{equation*}
$$

where

$$
\underset{\sim}{\mu}=\mathrm{g}_{N} \mu_{N^{\prime}} \cdot \underline{I}
$$

$<\mathrm{H}_{n f}>$ refers to the Boltzmann average of the hyperfine interaction represented by equation (2). The NMR chemical shift arising from the 3d electron orbital angular momnetum and the 3d electron spin dipolar-nuclear spin angular momentum interaction, when a three-fold axis is chosen as the axis of quantization, is given by

$$
\begin{align*}
& \Delta B / B \\
& \qquad=-\frac{\mu_{0}}{4 \pi} \frac{\mu_{8}^{t}}{K T}\left\{\frac{d(R)+[1-\exp (3 \zeta / 2 K T)] K T / \zeta S(\underline{R})}{1+2 \exp (3 \zeta / 2 K T)}\right\} \tag{6}
\end{align*}
$$

where

$$
\begin{aligned}
& d(R)=\frac{16 \sqrt{\pi}}{105} Y_{o 0}(\theta, \Phi) J(t)-\frac{7168}{t^{t}} Y_{3} E(t)-\frac{921600}{t^{7}} Y_{4} M(t) \\
& S \stackrel{(R)}{\sim}=-\frac{16 \sqrt{\pi}}{63} Y_{00}(\theta, \Phi) P(t)+\frac{7168}{t^{2}} Y, G(t)-\frac{614400}{t^{\top}} Y_{s} M(t)
\end{aligned}
$$

where

$$
\begin{aligned}
& Y_{s}= \sqrt{\frac{\pi}{21}}\left\{\frac{1}{3} \sqrt{\frac{7}{3}} Y_{4 t}(\theta, \Phi)-\frac{1}{3} \sqrt{\frac{10}{3}}\left[4_{t-3}(\theta, \Phi)-Y_{42}(\theta, \Phi)\right]\right\} \\
& Y_{t}= \sqrt{\frac{\pi}{26}}\left\{\frac{4 \sqrt{2}}{9} Y_{t 0}(\theta, \Phi)+\frac{1}{9} \sqrt{\frac{35}{3}}\left(Y_{t-3}(\theta, \Phi)-Y_{43}(\theta, \Phi)\right]\right. \\
&\left.+\frac{1}{9} \sqrt{\frac{77}{6}}\left[Y_{t-t}(\theta, \Phi)+Y_{4 t}(\theta, \Phi)\right]\right\} \\
&\left.M(i)=\beta^{t}\left\{1-e^{-z} \sum_{n=0}^{n} \frac{t^{n}}{n!}\right)\right\}
\end{aligned}
$$

Table 1. The NMR Chemical Shift for a 3d' System in a Strong Crystal field Environment of Octahedral Symmetry when Threefold Axis is Chosen as the Quantization Axis. $\left(\zeta=154 \mathrm{~cm}^{-1}, \beta=4 / 3 \mathrm{a}_{a}\right.$ and $\left.\mathrm{T}=300 \mathrm{~K}\right)$

| $\mathrm{R}(\mathrm{nm})$ | $<001\rangle$ | <100> | <010> | <110> | <111> |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0.05 | -1417.465 | - 1466.051 | -1445.746 | -1455.899 | $-1428.524$ |
| 0.10 | -402.246 | -463.234 | -395.726 | -429.480 | -394.346 |
| 0.15 | -61.262 | -87.121 | -30.603 | -58.862 | -43.453 |
| 0.20 | -1.378 | -9.935 | 18.533 | 4.299 | 9.577 |
| 0.25 | 4.137 | 0.724 | 12.292 | 6.508 | 8.617 |
| 0.30 | 3.007 | 1.226 | 5.618 | 3.422 | 4.493 |
| 0.35 | 1.833 | 0.782 | 2.480 | 1.631 | 2.246 |
| 0.40 | 1.104 | 0.461 | 1.159 | 0.810 | 1.180 |
| 0.45 | 0.679 | 0.276 | 0.587 | 0.431 | 0.661 |
| 0.50 | 0.429 | 0.171 | 0.321 | 0.246 | 0.392 |

Table 2. A Comparison of the Exact Value of $\Delta \mathrm{B} / \mathrm{B}$ (ppm) Using Equation (6) with the Multipolar Terms for Specific R-Values (2a) Along the <001> Axis

| $R(\mathrm{~nm})$ | $1 / \mathrm{R}^{s}$ | $1 / \mathrm{R}^{7}$ | sum of all <br> multipolar terms | From eq.(6) |
| :---: | ---: | ---: | :---: | ---: |
| 0.05 | 96.937 | -16.877 | 80.060 | -1417.465 |
| 0.1 | 161.408 | -56.111 | 105.296 | -402.247 |
| 0.2 | 40.608 | -23.663 | 16.945 | -1.378 |
| 0.3 | 7.004 | -3.651 | 3.353 | 3.008 |
| 0.4 | 1.690 | -0.581 | 1.109 | 1.104 |
| 0.5 | 0.554 | -0.124 | 0.430 | 0.430 |
| 0.6 | 0.223 | -0.035 | 0.188 | 0.188 |
| 0.7 | 0.103 | -0.012 | 0.091 | 0.091 |

(2b) Along the $\langle 100\rangle$ Axts

| $\mathrm{R}(\mathrm{nm})$ | $1 / \mathrm{R}^{s}$ | $1 / \mathrm{R}^{\prime}$ | sum of all <br> multipolar terms | From eq.(6) |
| :---: | ---: | ---: | :---: | ---: |
| 0.05 | 36.351 | -4.878 | 31.473 | -1466.052 |
| 0.1 | 60.528 | -16.219 | 44.308 | -463.235 |
| 0.2 | 15.228 | -6.840 | 8.388 | -9.935 |
| 0.3 | 2.626 | -1.055 | 1.571 | 1.226 |
| 0.4 | 0.352 | -0.075 | 0.277 | 0.277 |
| 0.5 | 0.129 | -0.018 | 0.111 | 0.111 |
| 0.6 | 0.056 | -0.006 | 0.050 | 0.050 |
| 0.7 | 0.038 | -0.003 | 0.035 | 0.035 |

(2c) Along the $<010>$ Axis

| $\mathbf{R}(\mathrm{nm})$ | $1 / R^{3}$ | $1 / \mathrm{R}^{\prime}$ | sum of all <br> multipolar terms | From eq.(6) |
| :---: | ---: | ---: | ---: | ---: |
| 0.05 | 36.351 | 15.327 | 51.778 | -1445.747 |
| 0.1 | 60.528 | 51.289 | 111.816 | -395.727 |
| 0.2 | 15.228 | 21.228 | 36.857 | 18.534 |
| 0.3 | 2.626 | 3.337 | 5.964 | 5.619 |
| 0.4 | 0.634 | 0.531 | 1.164 | 1.159 |
| 0.5 | 0.208 | 0.113 | 0.321 | 0.321 |
| 0.6 | 0.083 | 0.032 | 0.115 | 0.115 |
| 0.7 | 0.039 | 0.011 | 0.049 | 0.049 |

(2d) Along the <110> Axis

| $R(\mathrm{~nm})$ | $1 / \mathrm{R}^{s}$ | $1 / \mathrm{R}^{*}$ | sum of all <br> multipolar terms | From eq.(6) |
| :---: | ---: | ---: | ---: | ---: |
| 0.05 | 36.351 | $\mathbf{5 . 2 7 4}$ | 41.425 | -1455.899 |
| 0.10 | 60.527 | 17.535 | 78.062 | -429.481 |
| 0.20 | 15.228 | 7.395 | 22.625 | 4.299 |
| 0.30 | 2.626 | 1.141 | 3.767 | 3.422 |
| 0.40 | 0.634 | 0.181 | 0.815 | 0.810 |
| 0.50 | 0.208 | 0.039 | 0.246 | 0.246 |
| 0.60 | 0.083 | 0.011 | 0.094 | 0.094 |
| 0.07 | 0.039 | 0.004 | 0.042 | 0.042 |

## (2e) Along the <111> Axis

| $\mathrm{R}(\mathrm{nm})$ | $1 / \mathrm{R}^{3}$ | $1 / \mathrm{R}^{\prime}$ | sum of all <br> multipolar terms | From eq.(6) |
| :---: | ---: | ---: | ---: | ---: |
| 0.05 | 70.021 | -1.021 | 69.000 | -1428.524 |
| 0.10 | 116.591 | -3.394 | 113.197 | -394.346 |
| 0.20 | 29.333 | -1.431 | 27.901 | 9.578 |
| 0.30 | 5.059 | -0.221 | 4.838 | 4.493 |
| 0.40 | 1.220 | -0.035 | 1.185 | 1.181 |
| 0.50 | 0.400 | -0.008 | 0.392 | 0.392 |
| 0.60 | 0.161 | -0.002 | 0.159 | 0.159 |
| 0.70 | 0.074 | -0.001 | 0.054 | 0.054 |

$$
E(t)=\left\{1-e^{-t}\left(\frac{8}{11} \frac{t^{n}}{9!}+\sum_{n-4}^{n} \frac{t^{n}}{n!}\right)\right\}
$$

$$
J(t)=\beta^{\prime} e^{-t}\left(\frac{2}{9} \frac{t^{4}}{4!}+\sum_{n=0}^{n} \frac{t^{n}}{n!}\right)
$$

$$
G(t)=\beta^{4} e^{-t}\left(\frac{16}{33} \frac{t^{\prime}}{9!}\right)
$$

$$
\begin{equation*}
P(t)-\beta^{2} e^{-t}\left(-\frac{4}{45} \frac{t^{4}}{4!}+\sum_{n=0} \cdot \frac{t^{n}}{n!}\right) \text { with } t=2 \beta R \tag{6a}
\end{equation*}
$$

For the case of the free atom, the expression (6) may take the form given by

$$
\Delta B / B=-\frac{\mu_{0}}{4 \pi} \frac{\mu_{i}^{2}}{K T} \frac{8 \beta^{2}}{315}\left\{\frac{3-5(1-\exp (3 \zeta / 2 K T)) K T / \zeta}{1+2 \exp (3 \zeta / 2 K T)}\right)(6
$$

Table 3. The Temperature Dependence of $\Delta \mathrm{B} / \mathrm{B}(\mathrm{ppm})$ at Various of $\mathbf{R}$ Expressed in Terms of the Coefficients in a Strong Crystal Field Environment of Octahedral Symmetry when Threefold Axis is Chosen as the Quantization Axis. $\left(\zeta=154 \mathrm{~cm}^{-1}, \beta=4 / 3 \mathrm{a}_{0}\right)$

| $\mathbf{R}(\mathrm{nm})$ | axis | $\mathrm{A}(\mathrm{ppm})$ | $\mathrm{B}(\mathrm{ppm} \cdot \mathrm{K})$ | $\mathrm{C}\left(\mathrm{ppm} \cdot \mathrm{K}^{1}\right)$ |
| :--- | ---: | ---: | ---: | ---: |
| 0.10 |  | 3.016 | -151104.964 | 8864269.841 |
| 0.20 |  | 23.407 | -4091.950 | -1013419.422 |
| 0.30 | $\langle 001\rangle$ | 4.760 | 398.385 | -275368.236 |
| 0.40 |  | 1.045 | 253.580 | -71226.016 |
| 0.50 |  | 0.312 | 113.164 | -23486.317 |
| 0.10 |  | -39.474 | -167323.153 | 12085475.757 |
| 0.20 |  | 6.136 | -4655.555 | -52275.372 |
| 0.30 | $<100\rangle$ | 1.597 | 214.869 | -98555.688 |
| 0.40 |  | 0.367 | 116.394 | -26658.403 |
| 0.50 |  | 0.118 | 44.347 | -8471.374 |
| 0.10 |  | -71.901 | -137105.688 | 12027322.672 |
| 0.20 |  | -7.539 | 8087.561 | -76797.118 |
| 0.30 | $<010\rangle$ | -0.513 | 2181.078 | -102339.329 |
| 0.40 |  | 0.032 | 429.026 | -27259.999 |
| 0.50 |  | 0.040 | 113.999 | -8943.946 |
| 0.20 | $<110\rangle$ | -0.701 | 1716.030 | -64539.811 |
| 0.20 | $<111\rangle$ | 8.646 | 2263.367 | -599140.710 |

When $R$ is large

$$
\begin{aligned}
& d(R)=-\frac{1}{R^{4}}\left(\frac{224}{\beta^{2}} Y_{3}\right)-\frac{1}{R^{\prime}}\left(\frac{7200}{\beta^{4}} Y_{4}\right) \\
& S(\underset{\sim}{R})=\frac{1}{R^{2}}\left(\frac{224}{\beta^{2}} Y_{4}\right)-\frac{1}{R^{2}}\left(\frac{4800}{\beta^{4}} Y_{4}\right)
\end{aligned}
$$

When a threefold axis is chosen as the axis of quantization, the calculated NMR chemical shift for a 3d' system in a strong crystal field environment of octahedral symmetry using equation (6) along the $x, y$ and $z$ axes is listed Table 1. Here we choose $\beta=4 / 3 a_{0}$, the spin-orbit coupling constant $\zeta=154 \mathrm{~cm}^{-1}$ and the temperature $\mathrm{T}=300 \mathrm{~K}$. As shown in Table 1 , the NMR chemical shift for specific R-values for a $3 d^{1}$ system in a strong crystal field environment of octahedral symmetry along the $x, y$ and $z$ axis has different values, and the NMR chemical shift, $\Delta \mathrm{B} / \mathrm{B}(\mathrm{ppm})$, decreases in magnitude rapidly as R increases. $\Delta \mathrm{B} / \mathrm{B}$ changes sign when $\mathrm{R}=0.15 \mathrm{~nm}$ along the $\langle 010\rangle,\langle 110\rangle$ and $\langle 111\rangle$ and $R \approx 0.2$ nm along the $<0.01\rangle$ and $\langle 100\rangle$ axes, the values being negative for smaller $R$ values and positive for greater $R$ values.

When a threefold axis is chosen as the axis of quantization, a comparison of the exact values of $\Delta \mathrm{B} / \mathrm{B}(\mathrm{ppm})$ calculated by equation (6) with the multipolar terms for specific $R$-values shows that when $R=0.20 \mathrm{~nm}$ the constant term contributes dominantly to the NMR chemical shift and when $\mathrm{R}=0.20 \mathrm{~nm}$ the major contribution arises from $1 / R^{3}$, and along the $\langle 001\rangle$, $<100\rangle$ and $\langle 111\rangle$ axes the $1 / \mathrm{R}^{5}$ term gives values opposite in sign to that of the $1 / R^{7}$, while along the $\langle 010\rangle$, and $\langle 110\rangle$ axes the $1 / R^{s}$ and $1 / R^{7}$ terms gives values equal in sign. It is interesting to note that the exact solution given by equation (6) for a $3 \mathrm{~d}^{1}$ system in a strong crystal field environment of octahedral symmetry is in good agreement with the multipolar results when $\mathrm{R}>0.40 \mathrm{~nm}$. In addition, Table 2 shows
that the NMR chemical shift arising from the 3d electron orbital angular momentum and 3d electron spin dipolar-nuclear angular momentum interaction in a $3 \mathrm{~d}^{1}$ system in a strong crystal field environment of octahedral symmetry is large for significant distances between the NMR nucleus and the 3d electron bearing atom. For less than 0.40 nm it should not be neglected.

It is usual to correlate the temperature dependence of $\Delta \mathrm{B} / \mathrm{B}(\mathrm{ppm})$ to the expression

$$
\begin{equation*}
\Delta B=a<S_{*}>+b\left(\chi_{k}-\chi_{1}\right) \tag{7}
\end{equation*}
$$

The fitting of the data at a ligand position over a temperature range 200 to 400 K from equation (6) could be expressed as

$$
\begin{equation*}
\Delta B / B=A+B / T+C / T^{2} \tag{8}
\end{equation*}
$$

where the $I / T$ term arises fromt be Fermi contact term and the $1 / \mathrm{T}^{2}$ term from pseudo contact term. Therefore the temperature dependence of $\Delta \mathrm{B} / \mathrm{B}(\mathrm{ppm})$ for a $3 \mathrm{~d}^{\prime}$ system in a strong crystal field environment of octahedral symmetry, when a threefold axis is chosen as the axis of quantization, at ligand along the $z, x$ and $y$ axes, may be expressed, respectively, as

$$
\begin{align*}
& \Delta B / B=23.407-4091.95 / T-1013420 / T^{2} \\
& \Delta B / B=6.136-4655.56 / T-52275.3 / T^{2} \\
& \Delta B / B=-7.539+8087.56 / T-76797.1 / T^{2} \tag{9}
\end{align*}
$$

These results show that each term in equation (9) contributes significantly to the value of the NMR chemical shift and equation (9) is of the form given by equation (8). However, the ratio $\mathrm{A} / \mathrm{B}$ differs markedly from the expected value of $-32 / 105$ and so we may not use equation (7) to interpret the origin of the NMR chemical shift. The major contribution to the NMR chemical shift however arises from the $1 / \mathrm{T}^{2}$ term but the contributions of other two terms are significant as shown in Table 3.

It is necessary to mention at this moment that this work has already been carried out, since it is found however that some of the result are not correctly expressed," we rederive the general formula for a $3 \mathrm{~d}^{\prime}$ system in a strong crystal field environment of octahedral symmetry when a threefold axis is chosen as the axis of quantization and recalculate the NMR chemical shift. We extend further this work to investigate the temperature dependence of $\Delta \mathrm{B} / \mathrm{B}$ from equation (6).
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