# A Further Study on the NMR Chemical Shift for a $4 d^{1}$ System in a Strong Crystal Field Environment of Octahedral Symmetry 

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The various interactions considered in this paper in calculating the NMR chemical shift arising from the 4d orbital angular momentum and the 4 d electron spin dipolarnuclear spin angular momentum interaction may be expressed as

$$
\begin{equation*}
H=\frac{\hbar^{2}}{2 m} \nabla^{\prime}-\frac{Z e^{2}}{4 \pi e_{0} r}+V(r)+\xi \underline{\underline{S}} \underset{\sim}{S}+\mu_{\mathrm{a}}(\underline{l}+2 \underline{S}) \underline{B}+H_{h} \tag{1}
\end{equation*}
$$

where

$$
\begin{aligned}
& V(v) \\
& =a_{4}\left\{\frac{1}{2} \sqrt{\frac{7}{3}} Y_{4 *}(\theta, \phi)+\frac{1}{2} \sqrt{\frac{5}{6}}\left(Y_{4-1}(\theta, \phi)+Y_{4}(\theta, \phi)\right)\right\}^{\prime \prime}
\end{aligned}
$$

and

$$
\begin{equation*}
H_{A r}=\frac{\mu_{0}}{4} g_{N} \mu_{a} \mu_{N}\left\{\frac{2 \underline{l}_{N} \cdot \tilde{I}}{r_{N}^{2}}+g_{s}\left[\frac{3\left(\underline{r}_{N} \cdot S_{-}\right) \underline{\tau}_{N} \cdot \underline{I}}{r_{N}^{*}}-\frac{S \cdot I}{r_{M}^{3}}\right]\right\} \tag{2}
\end{equation*}
$$

Here $r$ and $\underline{g}_{N}$ are, respectively, the electron radius vectors about the electron-bearing atom and nucleus with nuclear spin angular momentum $I$. The quantity $B$ is the applied magnetic field. $V(\underline{r})$ is the crystal field potential of octahedral symmetry and $\mathrm{a}_{4}{ }^{2}$ is the required crystal field parameter for 4 d electron system. The other symbols have their usual meaning. In this work, $g_{n}$, the free electron $g$ value, is taken to be equal to exactly 2 .

To determine the NMR shift arising from the $4 d$ electron orbital angular momentum and the 4 d electron spin dipolarnuclear spin angular momentum interaction for a $4 \mathrm{~d}^{1}$ system in a strong crystal field environment of octahedral symmetry the principal values $\sigma_{\alpha}$ of the NMR screening tensor are calculated by considering the magnetic field interaction parallel to the $x, y$ and $z$ axes and averaged assuming a Boltzmann distribution. The contribution to the NMR chemical shift, $\Delta \mathrm{B}$, is given by ${ }^{3}$

$$
\begin{equation*}
\Delta B=B \sigma_{\alpha \alpha} / 3 \tag{3}
\end{equation*}
$$

where

$$
\sigma_{a c}=\left(\frac{\left.\delta^{2}<H_{N N}\right\rangle}{\delta \mu_{a} \delta B_{a}}\right)_{\mu m B-4}
$$

with

$$
\mu=g_{N} \mu_{N} I
$$

The term $\left\langle\mathrm{H}_{n}\right\rangle$ refers to the Boltzmann average of the hyper-fine interaction represented by equation (2). The NMR chemical shift for a $4 \mathrm{~d}^{1}$ system in a strong crystal field environment of octahedral symmetry is given by

$$
\begin{equation*}
\frac{\Delta B}{B}=-\frac{\mu_{0}}{4 \pi} \frac{\mu_{B}^{*}}{K T}\left\{\frac{d(R)+[1-\exp (3 \zeta / 2 K T)) K T / \zeta S(R)}{1+2 \exp (3 \zeta / 2 K T)}\right\} \tag{4}
\end{equation*}
$$

where

$$
\begin{aligned}
& d(\underset{\sim}{R}) \\
& \quad=\frac{12 \sqrt{\pi}}{147} Y_{o 0}(\theta, \Phi) E(t)+\frac{17280}{t^{5}} Y_{4} H(t)-\frac{22809600}{7 t^{\tau}} Y_{6} J(t) \\
& S(\underline{R}) \\
& \approx=-\frac{20 \sqrt{\pi}}{147} Y_{o 0}(\theta, \Phi) M(t)-\frac{7680}{t^{4}} Y_{t} N(t) N(t)-\frac{15206400}{7 t^{7}} Y_{t} J(t)
\end{aligned}
$$

where

$$
\begin{aligned}
& Y_{4} \\
& =\sqrt{\frac{\pi}{21}}\left\{\frac{1}{2} \sqrt{\frac{7}{3}} Y_{44}(\theta, \Phi)+\frac{1}{2} \sqrt{\frac{5}{6}}\left[Y_{4-4}(\theta, \Phi)+Y_{44}(\theta, \Phi)\right]\right\} \\
& Y_{4} \\
& =\sqrt{\frac{\pi}{26}}\left\{\frac{1}{2} \sqrt{\frac{1}{2}} Y_{n}(\theta, \Phi)-\frac{\sqrt{7}}{4}\left[Y_{4-c}(\theta, \Phi)+Y_{4}(\theta, \Phi)\right]\right\}
\end{aligned}
$$

Table 1. The NMR Chemical Shift for $\mathbf{a ~}^{4} \mathrm{~d}^{\text {d }}$ System in a Strong Crystal Field of Octahedral Symmetry. ( $\boldsymbol{\xi}=500 \mathrm{~cm}^{-1}, \beta=3.2679 / \mathrm{a}_{0}$ and $\mathrm{T}=300 \mathrm{~K}$ )

| $\mathrm{R}(\mathrm{nm})$ | $\langle 001\rangle$ | $\langle 100\rangle$ | $\langle 010\rangle$ | $\langle 110\rangle$ | $\langle 11\rangle\rangle$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 0.05 | -4701.679 | -4701.679 | -4701.679 | -1920.887 | -2032.734 |
| 0.10 | -551.519 | -551.519 | -551.519 | 273.451 | -234.919 |
| 0.15 | -46.595 | -46.595 | -46.595 | 48.658 | -38.796 |
| 0.20 | -7.386 | -7.386 | -7.386 | 7.639 | -5.401 |
| 0.25 | -1.909 | -1.909 | -1.909 | 1.703 | -0.907 |
| 0.30 | -0.657 | -0.657 | -0.657 | 0.506 | -0.170 |
| 0.35 | -0.273 | -0.273 | -0.130 | 0.184 | -0.024 |
| 0.40 | -0.130 | -0.068 | -0.068 | -0.068 | 0.078 |
| 0.45 | -0.038 | -0.038 | -0.038 | 0.037 | 0.005 |
| 0.50 |  |  | 0.109 | 0.010 |  |

$$
\begin{aligned}
& E(t)=\beta^{3} e^{-t}\left(\frac{1}{3} \frac{t^{n}}{6!}+\sum_{n=1}^{s} \frac{t^{n}}{n!}\right) \\
& H(t)=\beta^{2}\left\{1-e^{-t}\left(\frac{8}{9} \frac{t^{n t}}{11!}+\sum_{n=0}^{10} \frac{t^{n}}{n!}\right)\right\} \\
& J(t)=\beta^{3}\left\{1-e^{-t}\left(\frac{13}{11} \frac{t^{n \prime}}{13!}+\sum_{n}^{12} \frac{t^{n}}{n!}\right)\right\} \\
& M(t)=\beta^{2} e^{-t}\left(-\frac{2}{13} \frac{t^{t}}{6!}+\sum_{n=t} \frac{t^{n}}{n!}\right) \\
& N(t)=\beta^{2} e^{-t}\left(\frac{4}{3} \frac{t^{11}}{11!}\right) \text { with } t=2 \beta R
\end{aligned}
$$

For the case of the free atom we may take $\mathrm{R} \rightarrow 0$ to find

$$
\begin{equation*}
\frac{\Delta B}{\bar{R}} \rightarrow-\frac{\mu_{0}}{4 \pi} \frac{2 \beta^{2}}{147}\left\{\frac{3-5(1-\exp (3 \zeta / 2 K T)) K T / \xi}{1+2 \exp (3 \zeta / 2 K T)}\right\} \tag{5}
\end{equation*}
$$

When R is large
Table 2. A comparison of the Exact Value of $\Delta \mathrm{B} / \mathrm{B}(\mathrm{ppm})$ Using Equation (4) with the Multipolar Terms for Specific R-values (2a) Along the <100> Axis

| $R(\mathrm{~nm})$ | $1 / \mathrm{R}^{s}$ | $1 / \mathrm{R}^{\mathbf{y}}$ | sum of all <br> multipolar terms | From eq.(4) |
| :---: | ---: | ---: | ---: | ---: |
| 0.05 | -1714.588 | -242.831 | -1957.519 | -4701.679 |
| 0.1 | -275.418 | -183.123 | -458.541 | -551.519 |
| 0.2 | -3.160 | -4.219 | -7.239 | -7.386 |
| 0.3 | -0.409 | -0.249 | -0.658 | -0.658 |
| 0.4 | -0.097 | -0.032 | -0.130 | -0.130 |
| 0.5 | -0.032 | -0.007 | -0.039 | -0.039 |
| 0.6 | -0.009 | -0.001 | -0.010 | -0.010 |
| 0.7 | -0.006 | 0 | -0.006 | 0.006 |

(2b) Along the $\langle 110\rangle$ Axis

| $\mathbf{R}(\mathrm{nm})$ | $\mathrm{l} / \mathrm{R}^{3}$ | $1 / \mathrm{R}^{\prime}$ | sum of all <br> multipolar terms | From eq.(4) |
| :---: | ---: | ---: | ---: | ---: |
| 0.05 | 428.673 | 394.600 | 823.272 | -1920.888 |
| 0.1 | 68.854 | 297.576 | 366.430 | 273.452 |
| 0.2 | 0.790 | 6.857 | 7.647 | 7.640 |
| 0.3 | 0.102 | 0.405 | 0.507 | 0.507 |
| 0.4 | 0.024 | 0.054 | 0.078 | 0.078 |
| 0.5 | 0.008 | 0.011 | 0.019 | 0.019 |
| 0.6 | 0.003 | 0.003 | 0.006 | 0.006 |
| 0.7 | 0.001 | 0.001 | 0.003 | 0.003 |

(2c) Along the $<111>$ Axis

| $R(\mathrm{~nm})$ | $1 / \mathrm{R}^{3}$ | $1 / \mathrm{R}^{3}$ | sum of all <br> multipolar terms | From eq.(4) |
| :---: | ---: | :---: | ---: | ---: |
| 0.05 | 1143.126 | -431.699 | 711.426 | -2032.734 |
| 0.1 | 183.612 | -325.553 | -141.941 | -234.919 |
| 0.2 | 2.107 | -7.501 | -5.395 | -5.401 |
| 0.3 | 0.273 | -0.443 | -0.170 | -0.170 |
| 0,4 | 0.065 | -0.06 | 0.006 | 0.006 |
| 0.5 | 0.021 | -0.012 | 0.009 | 0.009 |
| 0.6 | 0.009 | -0.003 | 0.005 | 0.005 |
| 0.7 | 0.004 | -0.001 | 0.003 | 0.003 |

Table 3. The Temperature Dependence of $\Delta \mathrm{B} / \mathrm{B}(\mathrm{ppm})$ at Various of $\mathbf{R}$ Expressed in Terms of the Coeficients $\left(\zeta=500 \mathrm{~cm}^{-1}\right.$, and $\beta=3.2679 \mathrm{a}_{0}$ )

| $R(\mathrm{~nm})$ | axis | $\mathrm{A}(\mathrm{ppm})$ | $\mathrm{B}(\mathrm{ppm} \cdot \mathrm{K})$ | $\mathrm{C}\left(\mathrm{ppm} \cdot \mathrm{K}^{2}\right)$ |
| :---: | :---: | ---: | ---: | ---: |
| 0.10 |  | -578.1688 | 7293.1479 | 251237.4250 |
| 0.20 |  | -18.2736 | 4718.6856 | -437249.2688 |
| 0.30 | $<001\rangle$ | -2.3893 | 754.5414 | -70816.4113 |
| 0.40 |  | -0.5644 | 189.3803 | -17813.7882 |
| 0.50 |  | -0.1852 | 63.9512 | -6031.2148 |
| 0.20 | $\langle 110\rangle$ | 43.8326 | 1292.8309 | -136384.6603 |
| 0.20 | $\langle 111\rangle$ | 11.6756 | -7536.8862 | 727913.8958 |

$$
\begin{align*}
& d(R)=\frac{1}{R^{6}}\left(\frac{540}{\beta^{2}} Y_{4}\right)-\frac{1}{R^{\prime}}\left(\frac{178200}{\beta^{4}} Y_{4}\right) \\
& S(\underset{\sim}{R})=-\frac{1}{R^{6}}\left(\frac{240}{\beta^{2}} Y_{4}\right)-\frac{1}{R^{2}}\left(\frac{118800}{\beta^{4}} Y_{4}\right) \tag{6}
\end{align*}
$$

The calculated NMR chemical shift using equation (4) is listed in Table 1. Here we choose $\beta=3.2679 / \mathrm{a}_{0}$, the spin-orbit coupling constant $\zeta=500 \mathrm{~cm}^{-1}$ and the temperature $\mathrm{T}=300^{\circ} \mathrm{K}$. As expected, for a $4 \mathrm{~d}^{\mathbf{1}}$ system in a strong crystal field einvornment of ochahedral symmetry the NMR chemical shift for specific $R$-values is consistent with each other. As shown in Table 1, the NMR chemical shift $\Delta B / B$, decreases in magnitude rapidly as R increases. Along the $<100>$ axis, $\Delta \mathrm{B} / \mathrm{B}$ is negative for all values of R , while along the $\langle 110\rangle$ and $<111>$ axes, $\Delta \mathrm{B} / \mathrm{B}$ changes sign when $\mathrm{R} \simeq 0.05 \mathrm{~nm}$ and $\mathrm{R} \simeq$ 0.40 nm , respectively, the values being negative for smaller $R$ values and positive for greater $R$ values. A comparison of the exact values of $\Delta \mathrm{B} / \mathrm{B}(\mathrm{ppm})$ given by equation (4) with the multipolar terms for specific values R values shows that the $1 / R^{3}$ and $1 / R^{\boldsymbol{p}}$ terms contribute almost equally to the NMR chemical shift and along the $\left\langle 111>\right.$ axis, the $1 / R^{s}$ term gives values opposite in sign to that of the $1 / \mathrm{R}^{\mathbf{2}}$ term, while along the other axes, the signs of the $1 / \mathrm{R}^{\mathrm{s}}$ and $1 / \mathrm{R}^{\mathbf{y}}$ terms are equal.

It is interesting to note that the exact solution given by equation (4) for a $4 \mathrm{~d}^{1}$ system in a strong crystal field environment of octahedral symmetry is in good agreement with the multipolar results when $\mathrm{R} \geqslant 0.3 \mathrm{~nm}$. For distances less than 0.30 nm it should not be neglected.

The temperature dependence of the NMR chemical shift may be expressed as ${ }^{3}$

$$
\begin{equation*}
\Delta B / B=A+\frac{B}{T}+\frac{C}{T^{1}} \tag{7}
\end{equation*}
$$

where $A / B=-32 / 105$ for the $4 f^{\prime}$ ion in a crystal field environment of octahedral symmetry. In equation (7) the $1 / \mathrm{T}$ and $1 / \mathrm{T}^{2}$ terms arise from the Fermi and the pseudo contact interactions. ${ }^{6}$ It is found that the NMR chemical results over the temperature range from 200 to 400 K from the exact solution of $\Delta \mathrm{B} / \mathrm{B}(\mathrm{ppm})$ given by equation (4) may be fitted almost precisely to an expression of the form given by equation (7). The temperature dependence of the NMR chemical shift at the ligand may be expressed as

$$
\begin{equation*}
\Delta B / B=-18.2736+4718.68 / \mathrm{T}-437249 / \mathrm{T}^{2} \tag{8}
\end{equation*}
$$

The equation (8) indicates that the $1 / \mathrm{T}$ term and the $1 / \mathrm{T}^{2}$ term contribute significantly to the value of the NMR chemical shift. Although equation (8) is of the form given by equation (7) the
ratio $\mathrm{A} / \mathrm{B}$ differs markedly the expected value and hence the temperature dependence data may not be used to interpret the NMR chemical shift. We also observe from Table 3 that values of $A, B$ and $C$ depend markedly on the location of the NMR nucleus and the dominant contribution to $\Delta \mathrm{B} / \mathrm{B}$ arises from the $1 / \mathrm{T}^{2}$ term but the $1 / \mathrm{T}$ term is certainly significant.

At this moment it is required to mention that the NMR chemical shift has alreadly been examined and has derived a general mformula for a $4 \mathrm{~d}^{1}$ system in a strong crystal field environment of octahedral symmetry. ${ }^{7}$ It is however found that this formula is not correct, therefore we reexamine the NMR chemical shift for a $4 \mathrm{~d}^{\prime}$ system in a strong crystal field environment of octahedral symmetry and derive the correct expression for $\Delta \mathrm{B} / \mathrm{B}$. We also study further the temperature dependence of $\Delta \mathrm{B} / \mathrm{B}$.
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## References

1. R.M. Golding, R.O. Pascual and C. Suvanprakorn, "Advances Solution Chemistry, Plenum, 129 (1981).
2. R.M. Golding, R.O. Pascual and S. Ahn, J. Afagn, Recon., 46, 406 (1982).
3. L.C. Stubbs, Unpublished Thesis (Ph. D. The University of New South Wales, (1978).
4. R.M. Golding and L.C. Stubbs, Proc. R. Soc. Lond. A. 354, 223 (1977).
5. B. Bleaney, J. Magn, Reson., 8, 91 (1972).
6. J.W.M. DeBoer, P.J.D. Sakkers, C.W. Hitbers and E. DeBoer, J. Magn. Reson., 25, 455 (1977).
7. S. Ahn, S.W. Oh, and E. Park, Bull. Korean Chem. Soc., 4, 103 (1983).

# A Thermal and a Photochemical Mutarotation of $\alpha$-(D)-Glucose in Dimethyl Sulfoxide 

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Mutarotation of glucose in aqueous solvent has been extensively investigated,' but in nonaqueous solvents a few research works have been reported. ${ }^{2}$ Recently, Pincock ${ }^{3}$ et al. have reported the kinetic study of heterogeneous catalysis of glucose mutarotation by standard alumina suspended in dimethyl sulfoxide. In the alumina catalyzed mutarotation, of course, a thermal mutarotation was assumed to be negligiblle at roorn temperature. This led us to investigate a thermal and a photochemical mutarotation of $\alpha-(\mathrm{D})$-glucose in dimethyl

Table 1. Thermodynamic Constants for Mutarotation of $\alpha-(\mathrm{D})$ Glucose in DMSO and $\mathrm{H}_{2} \mathrm{O}$ at $30^{\circ} \mathrm{C}$. (a) reference (1) data, (b) this work except $\Delta H$

| Solvent | $\Delta \mathrm{H}$ <br> cal. $\mathrm{mol}^{-1} \mathrm{deg}^{-1}$ | $\Delta \mathrm{G}$ <br> cal. mol $^{-1}$ | $\Delta \mathrm{S}$ <br> e.s.u. |
| :---: | :---: | :---: | :---: |
| $\mathrm{H}_{2} \mathrm{O}$ (a) | -270 | -327 | 0.19 |
| DMSO (b) | -270 | -315 | 0.15 |

Table 2. Activation Energies for Mutarotation of $\alpha(\mathrm{D})$-Glucose in DMSO and $\mathrm{H}_{2} \mathrm{O}$ at $30^{\circ} \mathrm{C}$. (a) Kcal. mol ${ }^{-1}$ (b) cal. mol ${ }^{-1}$. $\mathrm{deg}^{-1}$. (c) references 7 and 8. (d) this work

| Solvent | E. ${ }^{*}$ | $\Delta G^{* *}$ | $\Delta H^{* *}$ | $\Delta S^{* 0}$ |
| :--- | :---: | :--- | :--- | :--- |
| $\mathrm{H}_{2} \mathrm{O}$ (c) | 17.24 | 24.66 | 17.24 | -24.9 |
|  | 16.09 | 22.08 | 16.09 | -19.8 |
| DMSO (d) <br> (thermal) | 21.28 | 23.20 | 20.0 | -10.55 |
| DMSO (d) <br> (photo.) | 10.28 | 22.44 | 9.62 | -42.3 |

sulfoxide. Although the mechanisms of glucose mutarotation in DMSO are not thoroughly understood at present, we wish to report the mechanism including thermodynamic parameters, temperature dependence of quantum yield and the relationship between the rate constants and the quantum yields.

Irradiation ${ }^{4}$ of $\alpha-(\mathrm{D})$-glucose in DMSO causes the photomutarotation and the reaction mixture reaches equilibrium at the optical rotation, $[\alpha]=51^{\circ}$. During the irradiation, a thermal mutarotation was also observed especially at the elevated temperature and the equilibrium mixture showed optical rotation, $[\alpha]=54^{\circ}$.

While the thermal mutarotation was reversible, the photomutarotation was irreversible, ${ }^{5}$ so that the two different rate equations ${ }^{6}$ were used to evaluate the rate constants (Figure 1 and 2). Although the photo-mutarotation is irreversible, photo-mutarotation reaches equilibrium at the optical rotation, $[\alpha]=51^{\circ}$. This equilibrium state may come from that the thermal equilibrium control overrides the irreversibie

Table 3. Temperature Dependence of Quantum Yields. (a) Calculated Values Using $k$, (photo). (b) Mmol Disappearance us. Milli Einstein Absorbed

| Temp $\left({ }^{\circ} \mathrm{C}\right)$ | $\mathrm{k}_{r}\left(\mathrm{~min}^{-1}\right)$ <br> thermal | $\mathrm{k}_{r}\left(\mathrm{~min}^{-1}\right)$ <br> photo | $\phi_{\text {must }}(\mathrm{b})$ |
| :---: | :---: | :---: | :---: |
| 30 | $1.06 \times 10^{-4}$ | $3.75 \times 10^{-4}$ | $0.21(0.20)$ |
| 40 | $2.91 \times 10^{-4}$ | $8.69 \times 10^{-4}$ | $0.48(0.48)$ |
| 50 | $1.05 \times 10^{-3}$ | $1.55 \times 10^{-3}$ | $0.86(0.84)$ |
| 60 | $2.09 \times 10^{-3}$ | $2.21 \times 10^{-3}$ | $1.28(1.25)$ |
| 70 | $5.75 \times 10^{-3}$ | $3.30 \times 10^{-3}$ | $1.83(1.91)$ |

