

Relationship between Carrier Concentration and Superconducting Transition Temperature in Bi-Sr-Ca-Cu-O Superconductor

Myung Chul Kim[†], Soon Ja Park^{*}

Department of Materials Science and Engineering, KunSan National University, KunSan 573-360, Korea

^{*}Department of Inorganic Materials Engineering, Seoul National University, Seoul 151-742, Korea

(Received Jun. 1, 1992)

요 약. 초산염(acetate)을 출발원료로 하여 Bi₂Sr₂Ca₂Cu₃O_y계 단일상 고온초전도체를 합성하려 하였으며 그 초전도상의 형성과정을 초전도전이온도 T_c 및 전하나르계 농도와 상관관계로부터 설명하고자 하였다. 초전도전이온도 T_c는 전기저항밀도 및 자화율을 측정하여 결정하였다. 시료로는 상기의 초전도체 출발조성물의 하소분말에 대해 850°C, 860°, 그리고 870°C에서 40시간 동안 공기 중 열처리한 것을 이용하였다. 전하나르계 농도는 홀계수를 측정하여 구하였다. 열처리과정 중의 초전도체 형성과정을 세라믹벌크내에 형성된 초전도체의 양 및 전하나르계 농도 분포의 상관관계로부터 설명하였다.

ABSTRACT. Bi₂Sr₂Ca₂Cu₃O_y-based superconductor phases were synthesized by the wet process using acetate precursors. Superconducting transition temperature (T_c) was determined from both measurements of electrical resistivity and magnetic susceptibility for the samples which were sintered at the temperatures of 850°C, 860°C, and 870°C for 40 hours, respectively. The values of carrier concentration from Hall measurements were compared with T_c data as a function of the sintering temperature. The formation mechanism of the superconducting phase was tentatively discussed on a basis of the distribution profile concept of the carrier concentration and the amount of superconducting phases in a ceramic bulk. This explanation may be supported by the experimental results of correlation between the relative amount of superconducting phases and the difference of T_c values between superconducting onset temperature and cutoff temperature at each sintering temperature.

Key Words : Bi-Sr-Ca-Cu-O Superconductor, Superconducting transition, formation mechanism.

1. Introduction

Since the first reporting on the high T_c superconductor by Bednorz and Muller¹, there have been lots of studies^{2, 3} on the superconducting mechanism of La₂CuO₄ system including YBCO and Bi(Tl)-Sr-Ca-Cu-O which have been found immediately. Among the methods for explaining the conduction mechanism of high T_c superconductor one of the most probable concept is the correlation of the

superconducting transition temperature, T_c, with the electron/hole concentration. Therefore many authors have reported⁴⁻⁸ the Hall coefficient as an indicator of the carrier concentration in relation with T_c for the purpose of explaining the conduction mechanism in high T_c superconductor.

It has been well known that synthesis of a exclusively single phase in Bi-Sr-Ca-Cu-O-based superconductor system is very difficult^{9, 10}. Even though there has been much effort to prepare the bulk

superconductor with a single 2223 phase, the synthesis of a purely high T_c phase has been little reported on the sintered bodies¹¹. Moreover, the reason is not yet clearly understood until now.

In this study the authors have investigated the correlation between carrier concentration from Hall coefficient measurement and T_c for the sintered bodies of Bi-Sr-Ca-Cu-O system superconductor and made an effort to explain the formation mechanism of a superconducting phase.

2. Experimental Procedure

The metal acetates were selected as highly reactive precursors both to attain the homogeneity in mixing and calcination and to make it possible to obtain a single phase of high T_c superconductor. Bismuth acetate (Kanto Chemical, special grade), strontium acetate (Wako Purechemical, special grade), calcium acetate (Wako Purechemical, special grade) and copper acetate (Wako Purechemical, special grade) were verified as above 99.9% purity.

Appropriate amounts of each reagent were weighed and then dissolved in deionized water to give a homogeneous solution of metal acetates. The solution was moderately evaporated and dried in a beaker on the hot stirrer. The homogeneous mixture of metal acetates was placed in an alumina crucible and then heated to the decomposition temperatures between 150°C and 300°C. After decomposition, the obtained black fine powders were ground in an agate mortar and calcined in muffle furnace at 800°C for 24 hours. The calcined powders were pelletized by dry press and sintered in air at 850°C~870°C for 40 hours and slowly cooled in furnace. The composition for all test samples were selected as 2:2:2:3 for Bi:Sr:Ca:Cu.

The crystallographic phases formed during sintering were determined by X-ray diffractometer (JEOL). Hall effects at room temperature were measured by RESITEST 8200 (Toyo Technica Co. Ltd.,) and silver paste was applied as an electrode material. Hall effect measurement for thin film as a standard four probe

DC measurement was adopted on the rectangularly sliced flat thin samples of the sintered bodies in a magnetic field of 1.5 T at room temperature. Superconducting transition temperature, T_c was determined by both electrical resistivity and magnetic susceptibility measurement. The electrical resistivity measurement was performed using DC 4 point probe method in a helium cryostat and the magnetic susceptibility was measured with a QUANTUM DESIGN (MPMS SQUID) susceptometer to characterize the exact T_c value and the relative amount of each superconducting phase. The Meissner effect was measured in a field of 2mT during warming up after cooling.

3. Results and Discussion

Decomposed acetate mixtures were in an amorphous form, but after sintering and annealing process it resulted in the well developed platelike crystals of superconductor phase. A significant amount of high T_c phase could not be detected in all samples by XRD as shown in Fig. 1, but the resistivity drop and Meissner effect were observed at 108 K for all tested samples.

Fig. 2 shows the temperature dependence of the resistivity for $\text{Bi}_2\text{Sr}_2\text{Ca}_2\text{Cu}_3\text{O}_y$ samples which were sintered at the different temperatures of 850°C, 860°C, and 870°C for 40 hours as indicated. The specific resistivity at temperature above $T_{c, \text{off}}$ increases drastically, whereas below that no significant difference in resistivity is observed. Among the above samples the increase of firing temperature leads to the lowering of resistivity and the increase of T_c value.

Fig. 3 represents the magnetic susceptibility measured by SQUID magnetometer for the samples in Fig. 2. $T_{c, \text{on}}$ was decided by the data analysis method of Koike et. al⁵ in ρ vs T plot and χ vs T plot. The midpoint of the transition curve in ρ vs T plot was selected as $T_{c, \text{on}}$ and this onset temperature of superconduction was well identified with the defined value in χ vs T plot where the extrapolated line of the steepest part as shown in Fig. 3 reaches the normal

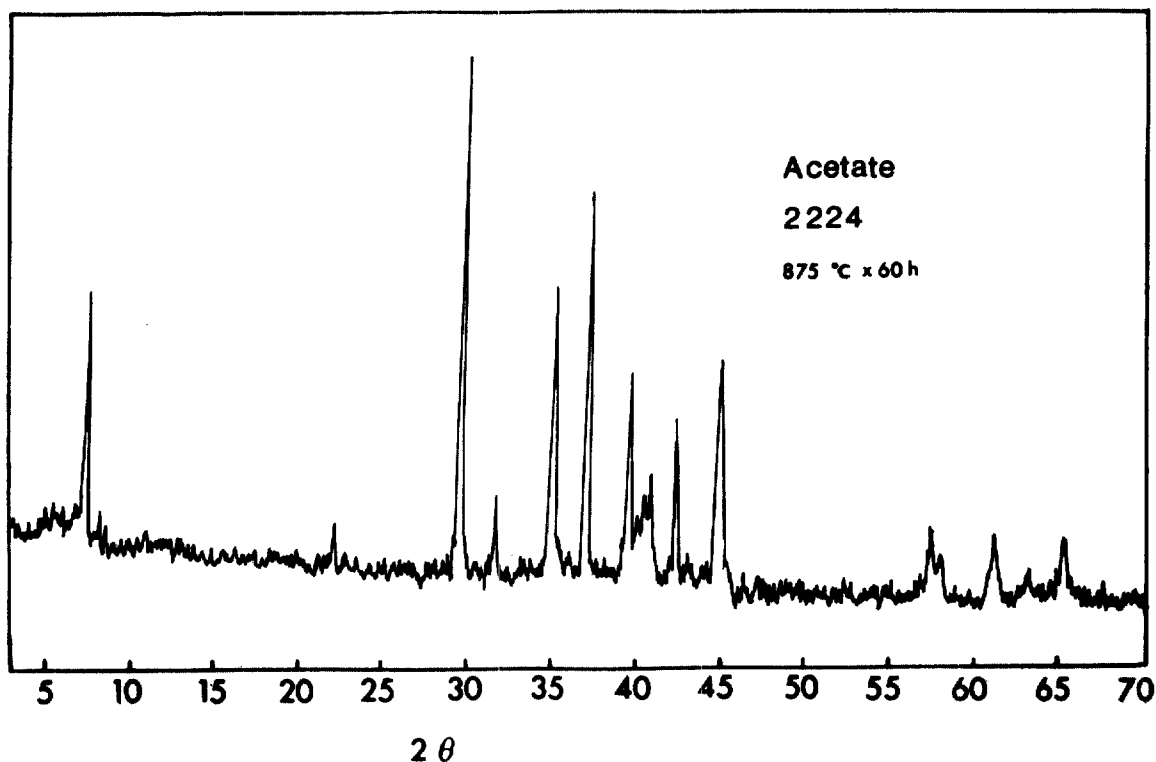


Fig. 1. X-ray peaks for the $\text{Bi}_2\text{Sr}_2\text{Ca}_2\text{Cu}_3\text{O}_y$ samples which sintered at 875°C for 60 hours in air. There exists only a slight 2223 phase.

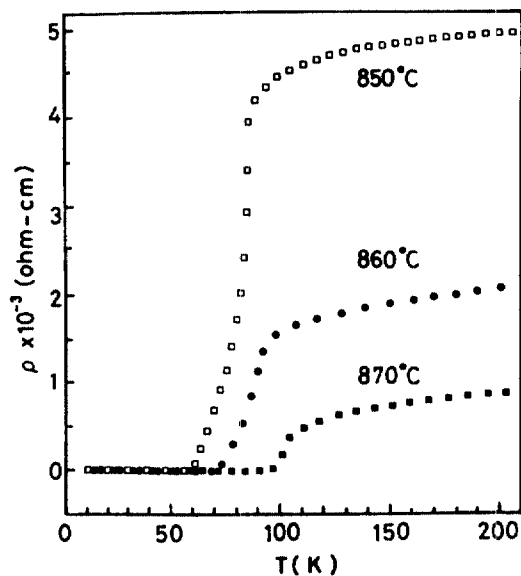


Fig. 2. Electrical resistivity with temperature for the $\text{Bi}_2\text{Sr}_2\text{Ca}_2\text{Cu}_3\text{O}_y$ samples which sintered at indicated temperatures for 40 hours in air.

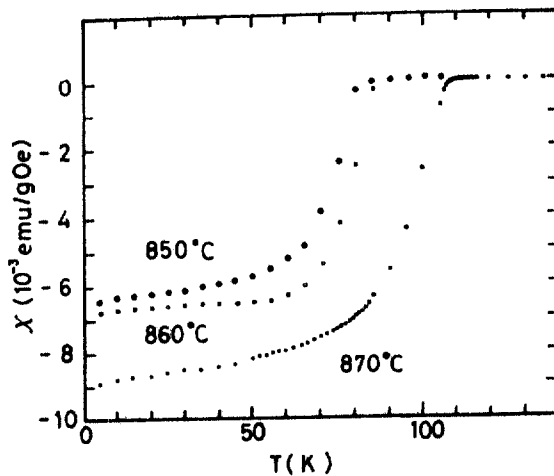


Fig. 3. The magnetic susceptibility with temperature for the samples of Fig. 2. The susceptibility values were determined by the SQUID susceptometer and these values were normalized.

state value of χ . These experimental results show the $T_{c, on}$ and $T_{c, off}$ explicitly and the relative amount of the superconducting phase among the tested samples can be estimated from the level of the normalized susceptibility values. From the above results, it is noted that $T_{c, on}$ of and $T_{c, off}$ each sample increase with the increase of firing temperature. As a result the gap of superconducting temperature T_c between $T_{c, on}$ and $T_{c, off}$ gradually decreases with the increase of sintering temperature in the tested temperature region. The normalized values of magnetic susceptibility roughly increase with the increase of firing temperature as shown in Fig. 4. The increase of firing temperature leads to the improvement of T_c level, the decrease of T_c gap, and the continuous increase of the amounts of the high T_c superconducting phase in the bulk.

In starting materials from the calcined powder, it may be postulated that there are mixtures of several kinds of phases such as nonsuperconducting phase, lower T_c phase, middle T_c phase, and/or higher T_c phase, etc. From this viewpoint, it is interesting to investigate the behavior of the carrier concentration and the firing temperature for the purposes of understanding the mechanism of the increase of T_c level and the formation of superconducting phase in the phase mixtures mentioned above. Fig. 4 shows the relationship between the values of $T_{c, on}$ and $T_{c, off}$ and the relative amount of superconducting phase estimated from χ value in Fig. 3. Besides the variation of the hole concentration, $1/R_{Hc}$ with temperature at room temperature, estimated from Hall coefficient R_H are shown together. The carrier concentration in the bulk monotonically increases with the increase of firing temperature in the tested temperature range as shown in Fig. 4, which is due to the decrease of electrical resistivity with the increase of firing temperature. Essentially Bi-Sr-Ca-Cu-O-based superconductor is a p-type conductor of hole carrier from Hall effect measurement and so the increase in the hole carrier concentration must lead to the lowering of electrical resistivity.

From the forgoing discussion we can draw such a structural image that several kinds of superconducting

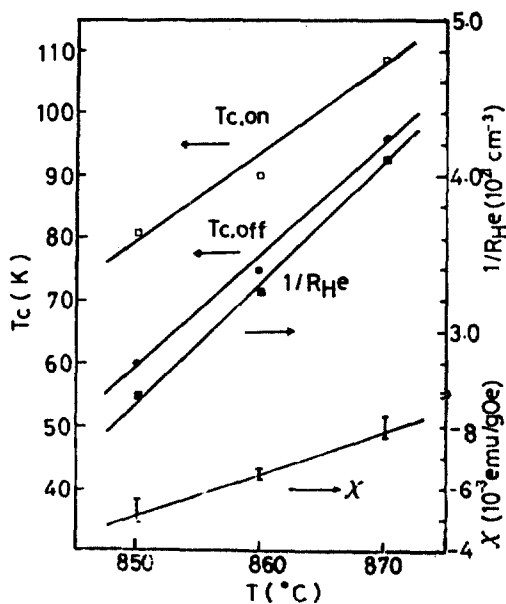


Fig. 4. The correlation between T_c , hole concentration $1/R_{He}$ at room temperature estimated from Hall coefficient R_H and the relative amount of superconducting phases in the $\text{Bi}_2\text{Sr}_2\text{Ca}_2\text{Cu}_3\text{O}_y$ which are sintered at temperatures of 850°C, 860°C, and 870°C for 40 hours in air. The relative amount of superconducting phases was estimated from the normalized values of magnetic susceptibility of Fig. 3.

phases and nonsuperconducting phases are mixed in a bulk and there is a local composition fluctuation within the individual phases of the starting calcined powders. It can be supposed that the phase with lower T_c value has the lower level of carrier concentration, because the T_c level of a superconducting phase is well proportional to the carrier concentration. The amount of a superconducting phase or a nonsuperconducting phase in the bulk may be distributed over the carrier concentration values with some range. It seems that the bulk phase does not take a single value of carrier concentration in the sintered body of high T_c superconductor because of the inevitably existing chemical inhomogeneity. We can consider the correlation between the carrier concentration distribution and the relative amount of superconducting phases and/or nonsuperconducting phases in a bulk. When

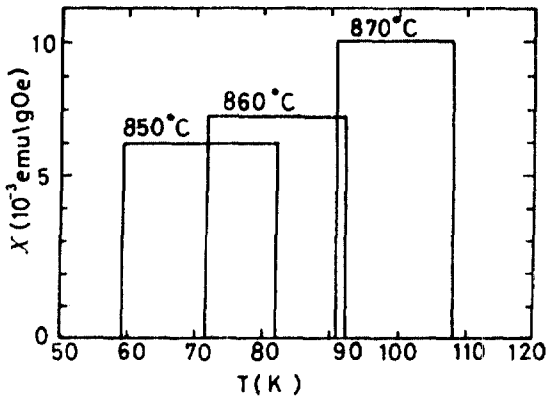


Fig. 5. The normalized value of magnetic susceptibility with $T_{c, on}$ and $T_{c, off}$ for the $\text{Bi}_2\text{Sr}_2\text{Ca}_2\text{Cu}_3\text{O}_y$ samples which were sintered at the temperatures of 850°C, 860°C, and 870°C for 40 hours in air, respectively. The relatively normalized values of χ were roughly estimated from Fig. 3. and schematically plotted against the distribution of T_c values between $T_{c, on}$ and $T_{c, off}$ to express the distribution profile of the amount and T_c values of superconducting phases.

this starting material is fired at sintering temperature the superconducting phases in the bulk will be formed and the carrier concentration of each phase will be distributed over some range according to the firing schedule. The distribution profile and the amount of the carrier concentration may be fluctuated and overlapped among the various phases in the bulk during the firing process. This hypothetically suggested distributional behavior of phases and their carrier concentration in a sintered body may be caused by the compositional and/or chemical variation which is inevitably introduced in the preparation process. If we use the highly reactive precursor such as e.g. acetates in the preparation, the inhomogeneity may be considerably diminished and so the distribution profile become sharper. In reality the above explanation will be supported by the relationship of χ vs $T_{c, on}$ and $T_{c, off}$ plot as shown in Fig. 5. In this figure the normalized value of magnetic susceptibility estimated from Fig. 3 are plotted against the distribution of T_c values be-

tween $T_{c, on}$ and $T_{c, off}$ to express the distribution profile of the amounts and T_c values of superconducting phases. So the amounts of the superconducting phases are distributed over the some range of T_c values. Therefore the successively rapid drop in the transition region in $\rho - T$ or $\chi - T$ curve is observed and it may be accounted for the existence of various superconducting phases with their T_c values in a bulk. From the above discussion it can be said that there may be lots of phases with a different carrier concentration and T_c value in a superconductor bulk. Accordingly, it is concluded that to get the single 2223 phase requires preparation of a homogeneous powder, extremely accurate furnace temperature control and very long annealing.

4. Conclusion

The formation mechanism of superconducting phases in $\text{Bi}_2\text{Sr}_2\text{Ca}_2\text{Cu}_3\text{O}_y$ bulk can be discussed on a basis of the correlation of carrier concentration, $1/R_{4e}$ and superconducting transition temperature, T_c . It may be considered that there are distribution profiles between the carrier concentration and the amount of superconducting phases with their T_c values in a bulk and the distribution will be fluctuated and overlapped among the various phases during the sintering process. This-like distribution profile may be accounted for the difficulty of the synthesis of a singly pure high T_c phase in Bi-Sr-Ca-Cu-O superconductor system and so to obtain the single 2223 phase requires accurate furnace temperature control and very long annealing.

References

1. Bednorz J. G. and Muller K. A., *Z. Phys. B. Cond. Matter.*, **64**, 186(1986).
2. D. H. Lee and J. Ihm, *Solid State Comm.*, **62**(12), 811-5(1987).
3. M. Avignon and K. H. Bennemann, *Solid State Comm.*, **68**(10), 999-1004(1989).
4. D. M. Eagles, *Physica C*, **153-155**, 701-2(1988).

5. Y. Koike, Y. Iwabuchi, S. Hosoya, N. Kobayashi, and T. Fukase, *Physica C*, **159**, 105-110(1989)
6. D. M. Eagles, *Solid State Comm.*, **69(3)**, 229-234(1989).
7. W. Umland and F. Tietz, *Solid State Comm.*, **69(10)**, 995-7(1989).
8. A. Davidson, P. Santhanam, A. Palevski, and M. J. Brady, *Phys. rev. B*, **38(4)**, 2828-31(1988).
9. A. Maqsood, M.S. Awan and N. Amin, *J. Mat. Sci.*, **26**, 4893-6(1991).
10. L. R. Yuan, K. Kurosawa, T. Matsushita, *Jap. J. appl. Phys.*, **30(9A)**, L1545-8(1991).
11. D. H. Chen, C. Y. Shei, S. R. Sheen and C. T. Chang, *Jap. J. Appl. Phys.*, **30(6)**, 1198-1203(1991).