

A CLASS OF CONDITIONAL ANALYTIC FEYNMAN INTEGRALS

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ABSTRACT. In this paper we establish the existence of the conditional Feynman integral of certain functions which are not in the Banach algebra \mathbf{S} of functions on Wiener space which are a kind of stochastic Fourier transform of complex Borel measures on $L^2[a, b]$. This result is used to provide the fundamental solution for the Schrödinger equation for the forced harmonic potential.

1. Introduction

In [3] Cameron and Storvick treated a Banach algebra \mathbf{S} of functions on Wiener space which are a kind of stochastic Fourier transform of complex Borel measures on $L^2[a, b]$. They established the existence of the analytic Feynman integral for all functions F in \mathbf{S} and give a formula for their Feynman integrals. In [10] Johnson has shown that \mathbf{S} is isometrically isomorphic to the class $\mathcal{F}(H)$ of Fresnel integrable functions as given by Albeverio and Hoegh-Krohn [1].

In [6], Chung and Skoug introduced the concept of a conditional analytic Feynman integral of a function F on Wiener space given a function X . For a certain choice of X , they established the existence of the conditional Feynman integral for all functions F in the Banach algebra \mathbf{S} and used the concept of conditional Feynman integral to provide a method of getting the fundamental solution to the Schrödinger equation and to obtain the kernel of various operator valued Feynman integrals [7].

In this paper we establish the existence of the conditional Feynman integral of certain functions which are not in the Banach algebra \mathbf{S} . This

Received October 10, 1995. Revised December 18, 1995.

1991 AMS Subject Classification: 28C20.

Key words and phrases: conditional Wiener integral, conditional Feynman integral.

¹ Research supported by the Sogang University Faculty Research Grant.

^{1,2,3} Research supported by BSRIP, MOE, 1994.

result is used to provide the fundamental solution for the Schrödinger equation for the forced harmonic potential.

2. Definitions and Preliminaries

Let $(C_0[0, T], \mathcal{B}(C_0[0, T]), m_w)$ denote Wiener space where $C_0[0, T]$ is the space of all continuous functions x on $[0, T]$ with $x(0) = 0$. For the partition $\tau = \tau_n = \{t_1, t_2, \dots, t_n\}$ of $[0, T]$ with $0 = t_0 < t_1 < t_2 < \dots < t_n = T$, let $X_\tau : C_0[0, T] \rightarrow \mathbf{R}^n$ be defined by $X_\tau(x) = (x(t_1), \dots, x(t_n))$. Let F be a complex-valued (\mathbf{C} -valued) integrable function on $C_0[0, T]$. Let \mathcal{F}_τ be the σ -algebra generated by the set $\{X_\tau^{-1}(B) : B \in \mathcal{B}(\mathbf{R}^n)\}$. Then, by the definition of conditional expectation, the conditional expectation of F given \mathcal{F}_τ , written $E[F|X_\tau]$, is any real valued \mathcal{F}_τ -measurable function on $C_0[0, T]$ such that

$$\int_E F dm_w = \int_E E[F|X_\tau] dm_w \quad \text{for } E \in \mathcal{F}_\tau.$$

It is well known that there exists a Borel measurable and P_{X_τ} -integrable function Ψ on $(\mathbf{R}^n, \mathcal{B}(\mathbf{R}^n), P_{X_\tau})$ such that $E[F|X_\tau] = \Psi \circ X_\tau$, where P_{X_τ} is the probability distribution of X_τ defined by $P_{X_\tau}(A) = m_w(X_\tau^{-1}(A))$ for $A \in \mathcal{B}(\mathbf{R}^n)$. Following Yeh [16], the function $\Psi(\xi)$, written $E[F|X_\tau = \vec{\xi}]$, is called *the conditional Wiener integral of F given X_τ* .

We note that a real valued function Y on $[0, T] \times C_0[0, T]$ defined by

$$Y(t, x) \equiv y(t) = x(t) - \frac{t}{T}x(T)$$

is a pinned Wiener process on $(C_0[0, T], \mathcal{B}(C_0[0, T]), m_w)$ and $[0, T]$ with $y(0) = 0$ and $y(T) = 0$. This process $\{y(t), 0 \leq t \leq T\}$ induces the Gaussian measure, called the pinned Wiener measure m_p , on $C_0^0[0, T] = \{x \in C_0[0, T] | x(T) = 0\}$, which is uniquely determined by mean function $E[y(t)] = 0$ for every $t \in [0, T]$ and covariance function $k(s, t) = E[y(s)y(t)] = \min\{s, t\} - \frac{st}{T}$.

The following proposition gives a convenient formula for evaluating conditional Wiener integral of functions involving quadratic potential $V(s, \xi)$ (see [4]).

PROPOSITION 2.1. Let F be an integrable function on $C_0[0, T]$ defined by $F(x) = \exp\{\int_{t_1}^T V(s, x(s))ds\}$. Then for $0 < t_1 < T$ and $\xi, \xi_1 \in \mathbf{R}$,

$$E[F(x)|x(t_1) = \xi_1, x(T) = \xi] = \int_{C_0^o[0, T-t_1]} F(y + g)dm_p(y),$$

where $g(t) = \frac{\xi - \xi_1}{T - t_1}t + \xi_1, t \in [0, T - t_1]$.

In particular, if $t_1 = 0$, then

$$E[F(x)|x(T) = \xi] = \int_{C_0^o[0, T]} F(y + h)dm_p(y)$$

where $h(t) = \frac{t}{T}\xi, t \in [0, T]$.

A subset E of $C_0[0, T]$ is said to be scale-invariant measurable ([5],[11]) provided ρE is Wiener measurable for each $\rho > 0$, and a scale-invariant measurable set N is said to be scale-invariant null provided $m(\rho N) = 0$ for each $\rho > 0$. A property that holds except on a scale-invariant null set is said to hold scale-invariant almost everywhere.

DEFINITION 2.1. Let X be an \mathbf{R}^n -valued scale-invariant measurable function on $C_0[0, T]$ and let F be a scale-invariant measurable function on

$$\int_{C_0[0, T]} F(\lambda^{-1/2}x)dm(x)$$

exists as a finite number for all $\lambda > 0$. For $\lambda > 0$, let

$$J_\lambda(\vec{\eta}) = E(F(\lambda^{-1/2}\cdot)|X(\lambda^{-1/2}\cdot) = \vec{\eta})$$

denote the conditional Wiener integral of $F(\lambda^{-1/2}\cdot)$ given $X(\lambda^{-1/2}\cdot)$. If for almost every $\vec{\eta} \in \mathbf{R}^n$, there exists a function $J_\lambda^*(\vec{\eta})$, analytic in λ on $\mathbf{C}^+ \equiv \{\lambda \in \mathbf{C} : \text{Re } \lambda > 0\}$ such that $J_\lambda^*(\vec{\eta}) = J_\lambda(\vec{\eta})$ for all $\lambda > 0$, then J_λ^* is defined to be the conditional analytic Wiener integral of F given X with parameter λ and for $\lambda \in \mathbf{C}^+$ we write

$$E^{\text{anw}\lambda}(F|X = \vec{\eta}) = J_\lambda^*(\vec{\eta}).$$

If for fixed real $q \neq 0$, the limit

$$\lim_{\lambda \rightarrow -iq} E^{\text{anf}\lambda}(F|X = \vec{\eta})$$

exists for almost every $\vec{\eta} \in \mathbf{R}^n$ where λ approaches $-iq$ through \mathbf{C}^+ , we will denote the value of this limit by $E^{\text{anf}q}(F|X)$ and call it *the conditional analytic Feynman integral of F given X with parameter q* .

3. Conditional analytic Feynman integrals

Let k be the covariance function of the pinned Wiener process $\{y(t) : t \in [0, T]\}$, that is, k is the function on $[0, T] \times [0, T]$ defined by

$$(3.1) \quad k(s, t) = \min\{s, t\} - \frac{st}{T}.$$

Let A be the integral operator on $L^2[0, T]$ (the space of real valued square integrable function on $[0, T]$) defined by

$$(3.2) \quad Af(s) = \int_0^T k(s, t)f(t)dt, \quad s \in [0, T], \quad f \in L^2[0, T].$$

Then it can be shown that the orthonormal eigen-functions $\{e_n\}$ of A and the corresponding eigen-value $\{\alpha_n\}$ are given by, respectively

$$(3.3) \quad e_n(s) = \sqrt{\frac{T}{2}} \sin\left(\frac{n\pi}{T}s\right) \quad \text{and} \quad \alpha_n = \frac{T^2}{n^2\pi^2}.$$

Furthermore, it can be shown that $\{e_n\}$ is a basis of $L^2[0, T]$, and that A is a trace class operator on $L^2[0, T]$. The Karhunen - Loeve theorem [2] shows that the Fourier series representation of the pinned Wiener process $\{y(t) : t \in [0, T]\}$ is given by

$$(3.4) \quad y(t) = \sum_{n=1}^{\infty} z_n e_n(t), \quad 0 \leq t \leq T$$

where z_n 's are independent Gaussian random variables with mean 0 and variance α_n .

We need the following lemmas(see [4]).

LEMMA 3.1. For $\alpha > 0, t \in [0, T]$,

$$(3.5) \quad \sum_{n=1}^{\infty} \frac{T}{n^2\pi^2 + \alpha T^2} \cos\left(\frac{n\pi}{T}t\right) = \frac{\cosh\sqrt{\alpha}(T-t)}{2\sqrt{\alpha}\sinh\sqrt{\alpha}T} - \frac{1}{2\alpha T}.$$

LEMMA 3.2. For $\alpha > 0$, let

$$R(s, t, \alpha) = \sum_{n=1}^{\infty} \frac{\alpha_n}{1 + \alpha\alpha_n} e_n(s)e_n(t), \quad s, t \in [0, T]$$

where α_n and e_n are as in (3.3). Then for each $t \in [0, T]$,

$$(3.6) \quad R(s, t, \alpha) = \begin{cases} \frac{\sinh\sqrt{\alpha}(T-t)\sinh\sqrt{\alpha}s}{\sqrt{\alpha}\sinh\sqrt{\alpha}T}, & 0 \leq s \leq t; \\ \frac{\sinh\sqrt{\alpha}(T-s)\sinh\sqrt{\alpha}t}{\sqrt{\alpha}\sinh\sqrt{\alpha}T}, & t \leq s \leq T. \end{cases}$$

THEOREM 3.3. Let F be a measurable function on $C_0[0, T]$ defined by

$$F(x) = \exp\left\{-\frac{1}{2}\alpha \int_{t_1}^T x^2(s)ds + \alpha\beta \int_{t_1}^T v(s)x(s)ds\right\},$$

where $\operatorname{Re} \alpha \geq 0, \beta \in \mathbf{C}$, and $v \in L^2[0, T]$. Let $X(x) = (x(t_1), x(T))$. Then for all $\lambda \in \mathbf{C}^+$, $E^{\text{anw}\lambda}(F|X)$ exists and for all $(\xi_1, \xi) \in \mathbf{R}^2$ is given by the formula

$$(3.7) \quad \begin{aligned} E^{\text{anw}\lambda}[F|X = (\xi_1, \xi)] &= \left(\frac{\sqrt{\frac{\alpha}{\lambda}}(T-t_1)}{\sinh\sqrt{\frac{\alpha}{\lambda}}(T-t_1)}\right)^{\frac{1}{2}} \cdot \exp\left\{\frac{\lambda(\xi - \xi_1)^2}{2(T-t_1)}\right\} \\ &\cdot \exp\left\{-\frac{\sqrt{\alpha\lambda}}{2}(\xi^2 + \xi_1^2)\coth\sqrt{\frac{\alpha}{\lambda}}(T-t_1) + \frac{\sqrt{\alpha\lambda}\xi\xi_1}{\sinh\sqrt{\frac{\alpha}{\lambda}}(T-t_1)}\right\} \\ &\cdot \exp\left\{\alpha\beta \int_{t_1}^T \frac{v(t)}{\sinh\sqrt{\frac{\alpha}{\lambda}}(T-t_1)} \left(\xi \sinh\sqrt{\frac{\alpha}{\lambda}}(t-t_1) \right. \right. \\ &\quad \left. \left. + \xi_1 \sinh\sqrt{\frac{\alpha}{\lambda}}(T-t)\right) dt\right\} \\ &\cdot \exp\left\{\frac{\alpha^2\beta^2}{2\lambda} \int_0^{T-t_1} \int_0^{T-t_1} R\left(s, t, \frac{\alpha}{\lambda}\right) v(s+t_1)v(t+t_1)dsdt\right\} \end{aligned}$$

Furthermore, $E^{\text{anf}_q}(F|X)$ exists for all real $q \neq 0$ and for all $(\xi_1, \xi) \in \mathbf{R}^2$ is given by the formula

$$\begin{aligned}
 E^{\text{anf}_q}[F|X = (\xi_1, \xi)] &= \left(\frac{\sqrt{\frac{\alpha i}{q}}(T - t_1)}{\sinh \sqrt{\frac{\alpha i}{q}}(T - t_1)} \right)^{\frac{1}{2}} \exp \left\{ \frac{q(\xi - \xi_1)^2}{2i(T - t_1)} \right\} \\
 &\cdot \exp \left\{ -\frac{\sqrt{-iq\alpha}}{2}(\xi^2 + \xi_1^2) \coth \sqrt{\frac{\alpha i}{q}}(T - t_1) + \frac{\sqrt{-iq\alpha}\xi\xi_1}{\sinh \sqrt{\frac{\alpha i}{q}}(T - t_1)} \right\} \\
 (3.8) \quad &\cdot \exp \left\{ \alpha\beta \int_{t_1}^T \frac{v(t)}{\sinh \sqrt{\frac{\alpha i}{q}}(T - t_1)} \left(\xi \sinh \sqrt{\frac{\alpha i}{q}}(t - t_1) \right. \right. \\
 &\qquad \qquad \qquad \left. \left. + \xi_1 \sinh \sqrt{\frac{\alpha i}{q}}(T - t) \right) dt \right\} \\
 &\cdot \exp \left\{ \frac{\alpha^2\beta^2 i}{2q} \int_0^{T-t_1} \int_0^{T-t_1} R(s, t, \frac{\alpha i}{q}) v(s + t_1) v(t + t_1) ds dt \right\}
 \end{aligned}$$

PROOF. Note that for any $z \in \mathbf{C}$ and $\lambda > 0$, $\exp\{\frac{z}{\sqrt{\lambda}} \int_0^T v(s)x(s)ds\}$ is Wiener integrable. Hence $F(\lambda^{-\frac{1}{2}}\cdot)$ is Wiener integrable for $\lambda > 0$, $\text{Re}\alpha \geq 0$ and any $\beta \in \mathbf{C}$. So $F(\lambda^{-\frac{1}{2}}\cdot)$ is conditional Wiener integrable for the given $x(t_1) = \xi_1$ and $x(T) = \xi$. By Proposition 2.1, we have, for $\xi, \xi_1 \in \mathbf{R}$ and $\lambda > 0$

$$\begin{aligned}
 E[F(\lambda^{-\frac{1}{2}}\cdot) | \lambda^{-\frac{1}{2}}x(t_1) = \xi_1, \lambda^{-\frac{1}{2}}x(T) = \xi] \\
 = \int_{C_0^0[0, T-t_1]} \exp \left\{ -\frac{\alpha}{2\lambda} \int_0^{T-t_1} (y(s) + \sqrt{\lambda}g(s))^2 ds \right. \\
 \left. + \frac{\alpha\beta}{\sqrt{\lambda}} \int_0^{T-t_1} v(s + t_1)(y(s) + \sqrt{\lambda}g(s)) ds \right\} dm_p(y)
 \end{aligned}$$

where $g(t) = \frac{\xi - \xi_1}{T - t_1}t + \xi_1, t \in [0, T - t_1]$. The previous quantity equals

(3.9)

$$\int_{C_0^0[0, T-t_1]} \exp \left\{ -\frac{1}{2}\alpha \sum_{n=1}^{\infty} \left[\frac{1}{\lambda}(z_n + \sqrt{\lambda}g_n)^2 - \frac{2\beta}{\sqrt{\lambda}}v_n(z_n + \sqrt{\lambda}g_n) \right] \right\} dm_p$$

where $y(t) = \sum_{n=1}^{\infty} z_n e_n(t)$ is the Fourier series representation of function y in $C_0^0[0, T - t_1]$ as in (3.4), $g(t) = \sum_{n=1}^{\infty} g_n e_n(t)$, and $v(t + t_1) = \sum_{n=1}^{\infty} v_n e_n(t)$. Since z'_n s are independent Gaussian random variables with mean 0 and variance α_n , (3.10) equals

$$\begin{aligned} & \prod_{n=1}^{\infty} \int_{C_0^0[0, T-t_1]} \exp \left\{ -\frac{\alpha}{2\lambda} z_n^2 + \frac{\alpha}{\sqrt{\lambda}} (\beta v_n - g_n) z_n + \alpha \beta g_n v_n - \frac{\alpha}{2} g_n^2 \right\} dm_p \\ &= \prod_{n=1}^{\infty} \left[\left\{ \frac{1}{\sqrt{2\pi\alpha_n}} \int_{\mathbf{R}} \exp \left\{ -\frac{\alpha}{2\lambda} u^2 + \frac{\alpha}{\sqrt{\lambda}} \omega_n u - \frac{u^2}{2\alpha_n} \right\} du \right\} \right. \\ & \quad \left. \cdot \exp \left\{ \alpha \beta g_n v_n - \frac{\alpha}{2} g_n^2 \right\} \right], \end{aligned}$$

where $\omega_n = \beta v_n - g_n$. Hence (3.10) equals

$$\begin{aligned} & \prod_{n=1}^{\infty} \left[\frac{1}{\sqrt{2\pi\alpha_n}} \left\{ \int_{\mathbf{R}} \exp \left\{ -\frac{1}{2} \left(\frac{\alpha}{\lambda} + \frac{1}{\alpha_n} \right) \left(u - \frac{\sqrt{\lambda} \alpha \alpha_n \omega_n}{\alpha \alpha_n + \lambda} \right)^2 \right. \right. \right. \\ & \quad \left. \left. \left. + \frac{\alpha^2 \alpha_n \omega_n^2}{2(\alpha \alpha_n + \lambda)} \right\} du \right\} + \alpha \beta g_n v_n - \frac{\alpha}{2} g_n^2 \right] \\ (3.10) \quad &= \prod_{n=1}^{\infty} \left[\left(1 + \frac{\alpha}{\lambda} \alpha_n \right)^{-\frac{1}{2}} \exp \left\{ \frac{\alpha^2 \beta^2 \alpha_n v_n^2}{2(\alpha \alpha_n + \lambda)} + \frac{\alpha^2 \alpha_n g_n^2}{2(\alpha \alpha_n + \lambda)} - \frac{\alpha^2 \beta \alpha_n g_n v_n}{\alpha \alpha_n + \lambda} \right. \right. \\ & \quad \left. \left. + \alpha \beta g_n v_n - \frac{\alpha}{2} g_n^2 \right\} \right] \\ &= \left[\prod_{j=1}^{\infty} \left(1 + \frac{\alpha}{\lambda} \alpha_n \right) \right]^{-\frac{1}{2}} \exp \left\{ \frac{\alpha^2 \beta^2}{2} \sum_{n=1}^{\infty} \frac{\alpha_n}{\alpha \alpha_n + \lambda} v_n^2 + \frac{\alpha^2}{2} \sum_{n=1}^{\infty} \frac{\alpha_n}{\alpha \alpha_n + \lambda} g_n^2 \right. \\ & \quad \left. - \alpha^2 \beta \sum_{n=1}^{\infty} \frac{\alpha_n}{\alpha \alpha_n + \lambda} v_n g_n + \alpha \beta \sum_{n=1}^{\infty} g_n v_n - \frac{\alpha}{2} \sum_{n=1}^{\infty} g_n^2 \right\}. \end{aligned}$$

Using

$$\prod_{n=1}^{\infty} \left(1 - \frac{z^2}{n^2 \pi^2} \right) = \frac{\sin z}{z}, \quad \alpha_n = \frac{(T - t_1)^2}{n^2 \pi^2},$$

we have

$$(3.11) \quad \prod_{n=1}^{\infty} \left(1 + \frac{\alpha}{\lambda} \alpha_n \right) = \frac{\sinh \sqrt{\frac{\alpha}{\lambda}} (T - t_1)}{\sqrt{\frac{\alpha}{\lambda}} (T - t_1)}.$$

Observing that

$$(3.12) \quad \sum_{n=1}^{\infty} \frac{\alpha_n}{\lambda + \alpha\alpha_n} v_n^2 = \frac{1}{\lambda} \int_0^{T-t_1} \int_0^{T-t_1} R\left(s, t, \frac{\alpha}{\lambda}\right) v(s+t_1)v(t+t_1) ds dt$$

$$(3.13) \quad \sum_{n=1}^{\infty} \frac{\alpha_n}{\lambda + \alpha\alpha_n} g_n^2 = \frac{1}{\lambda} \int_0^{T-t_1} \int_0^{T-t_1} R\left(s, t, \frac{\alpha}{\lambda}\right) g(s)g(t) ds dt$$

$$(3.14) \quad \sum_{n=1}^{\infty} \frac{\alpha_n}{\lambda + \alpha\alpha_n} v_n g_n = \frac{1}{\lambda} \int_0^{T-t_1} \int_0^{T-t_1} R\left(s, t, \frac{\alpha}{\lambda}\right) v(s+t_1)g(t) ds dt$$

$$(3.15) \quad \sum_{n=1}^{\infty} g_n v_n = \int_0^{T-t_1} g(t)v(t+t_1) dt$$

$$(3.16) \quad \sum_{n=1}^{\infty} g_n^2 = \int_0^{T-t_1} g^2(t) dt$$

and using Lemmas 3.1 and 3.2 with replacing T by $T - t_1$, one can show that

$$(3.17) \quad \begin{aligned} & \frac{\alpha^2}{2} \sum_{n=1}^{\infty} \frac{\alpha_n}{\alpha\alpha_n + \lambda} g_n^2 - \frac{\alpha}{2} \sum_{n=1}^{\infty} g_n^2 \\ &= -\frac{1}{2} \left\{ \sqrt{\alpha\lambda}(\xi^2 + \xi_1^2) \coth \sqrt{\frac{\alpha}{\lambda}}(T-t_1) - \frac{2\sqrt{\alpha\lambda}\xi\xi_1}{\sinh \sqrt{\frac{\alpha}{\lambda}}(T-t_1)} \right. \\ & \quad \left. - \lambda \frac{(\xi - \xi_1)^2}{T-t_1} \right\}, \end{aligned}$$

$$(3.18) \quad \begin{aligned} & -\alpha^2\beta \sum_{n=1}^{\infty} \frac{\alpha_n}{\lambda + \alpha\alpha_n} g_n v_n + \alpha\beta \sum_{n=1}^{\infty} g_n v_n \\ &= \alpha\beta \int_{t_1}^T \left(\frac{\xi \sinh \sqrt{\frac{\alpha}{\lambda}}(t-t_1) + \xi_1 \sinh \sqrt{\frac{\alpha}{\lambda}}(T-t)}{\sinh \sqrt{\frac{\alpha}{\lambda}}(T-t_1)} \right) v(t+t_1) dt. \end{aligned}$$

Putting (3.11),(3.12),(3.17) and (3.18) in the last equation in (3.10) we obtain the formula in (3.7). But one can show the right-hand side in (3.7) is an analytic function of λ throughout \mathbf{C}^+ and is a continuous function of λ for $\text{Re}\lambda \geq 0, \lambda \neq 0$. Hence (3.7) and (3.8) are established, which concludes the proof of theorem.

COLLOARY 3.4. Let $\text{Re } \alpha > 0$ and F be an integrable function on $C_0[0, T]$ defined by $F(x) = \exp\{-\frac{\alpha}{2} \int_0^T x^2(s)ds + \alpha\beta \int_0^T v(s)x(s)ds\}$. Let $0 = t_0 < t_1 < \dots < t_n = T$ and $X(t) = (x(t_1), x(t_2), \dots, x(t_n))$. Then we have, for $\vec{\xi} = (\xi_1, \xi_2, \dots, \xi_n) \in \mathbf{R}^n$ and real $q \neq 0$,

$$\begin{aligned}
 E^{\text{anf}_q}[F|X = \vec{\xi}] &= \prod_{k=1}^n \left[\left(\frac{\sqrt{\frac{\alpha i}{q}}(t_k - t_{k-1})}{\sinh \sqrt{\frac{\alpha i}{q}}(t_k - t_{k-1})} \right)^{\frac{1}{2}} \cdot \exp \left\{ \frac{q(\xi_k - \xi_{k-1})^2}{2i(t_k - t_{k-1})} \right\} \right. \\
 &\cdot \exp \left\{ -\frac{\sqrt{-iq\alpha}}{2}(\xi_k^2 + \xi_{k-1}^2) \coth \sqrt{\frac{\alpha i}{q}}(t_k - t_{k-1}) + \frac{\sqrt{-iq\alpha}\xi_k\xi_{k-1}}{\sinh \sqrt{\frac{\alpha i}{q}}(t_k - t_{k-1})} \right\} \\
 &\cdot \exp \left\{ \alpha\beta \int_{t_{k-1}}^{t_k} \frac{v(t)}{\sinh \sqrt{\frac{\alpha i}{q}}(t_k - t_{k-1})} \left(\xi_k \sinh \sqrt{\frac{\alpha i}{q}}(t - t_{k-1}) \right. \right. \\
 &\quad \left. \left. + \xi_{k-1} \sinh \sqrt{\frac{\alpha i}{q}}(t_k - t) \right) dt \right\} \\
 &\cdot \exp \left\{ \frac{\alpha^2\beta^2 i}{2q} \int_0^{t_k - t_{k-1}} \int_0^{t_k - t_{k-1}} R(s, t, \frac{\alpha i}{q})v(s + t_{k-1})v(t + t_{k-1})dsdt \right\} \Big],
 \end{aligned}$$

where $t_0 = 0, \xi_0 = 0$ and $R(s, t, \alpha)$ is as in (3.6), with replacing T by $t_k - t_{k-1}$.

PROOF. Let $V(s, \xi) = \alpha\xi^2 - 2\alpha\beta v(s)\xi$. Since the Wiener process $\{x(s) : 0 \leq s \leq T\}$ is additive, it can be shown that

$$\begin{aligned}
 &E \left[\exp \left\{ -\frac{1}{2} \int_0^T V(s, \lambda^{-\frac{1}{2}}x(s))ds \right\} \middle| \lambda^{-\frac{1}{2}}x(t_k) = \xi_k, k = 1, 2, \dots, n \right] \\
 &= E \left[\exp \left\{ -\frac{1}{2} \sum_{k=1}^n \left\{ \int_{t_{k-1}}^{t_k} V(s, \lambda^{-\frac{1}{2}}x(s))ds \right\} \right\} \middle| \lambda^{-\frac{1}{2}}x(t_k) = \xi_k, \right. \\
 &\quad \left. k = 1, 2, \dots, n \right]
 \end{aligned}$$

$$\begin{aligned}
 &= \prod_{k=1}^n E \left[\exp \left\{ -\frac{1}{2} \int_{t_{k-1}}^{t_k} V(s, \lambda^{-\frac{1}{2}} x(s)) ds \right\} \middle| \lambda^{-\frac{1}{2}} x(t_{k-1}) \right. \\
 &\qquad \qquad \qquad \left. = \xi_{k-1}, \lambda^{-\frac{1}{2}} x(t_k) = \xi_k \right].
 \end{aligned}$$

Hence this, together with Theorem 3.3, gives the desired result.

If we let $v(t) \equiv 0$ in Corollary 3.4, we then have

COROLLARY 3.5. *Let α be a complex number with $\text{Re}\alpha > 0$. Let $0 = t_0 < t_1 < \dots < t_n = T$. Then for $\xi = (\xi_1, \xi_2, \dots, \xi_n) \in \mathbf{R}^n$ and real $q \neq 0$*

$$\begin{aligned}
 &E^{\text{anf}_q} \left[\exp \left\{ -\frac{1}{2} \alpha \int_0^T x^2(s) ds \right\} \middle| x(t_k) = \xi_k, k = 1, 2, \dots, n \right] \\
 &= \prod_{k=1}^n \left[\left(\frac{\sqrt{\frac{\alpha i}{q}}(t_k - t_{k-1})}{\sinh \sqrt{\frac{\alpha i}{q}}(t_k - t_{k-1})} \right)^{\frac{1}{2}} \cdot \exp \left\{ \frac{q(\xi_k - \xi_{k-1})^2}{2i(t_k - t_{k-1})} \right\} \right. \\
 &\quad \cdot \exp \left\{ -\frac{\sqrt{-iq\alpha}}{2} (\xi_k^2 + \xi_{k-1}^2) \coth \sqrt{\frac{\alpha i}{q}}(t_k - t_{k-1}) \right. \\
 &\quad \left. \left. + \frac{\sqrt{-iq\alpha} \xi_k \xi_{k-1}}{\sinh \sqrt{\frac{\alpha i}{q}}(t_k - t_{k-1})} \right\} \right]
 \end{aligned}$$

COROLLARY 3.6[9, P.64]. *The solution of the Schrödinger equation for an external force $f(t)$:*

$$(3.19) \quad \frac{\partial U}{\partial t} = \frac{i\hbar}{2m} \frac{\partial^2 U}{\partial \xi^2} + \frac{i}{\hbar} \left(-\frac{mw^2}{2} \xi^2 + f(t)\xi \right) U$$

satisfying $U(\xi_2, t_2; \xi_1, t_1) \rightarrow \delta(\xi_2 - \xi_1)$ as $t_2 \downarrow t_1$ is given by

$$\begin{aligned}
 (3.20) \quad &U(\xi_2, t_2; \xi_1, t_1) = \sqrt{\frac{mw}{2\pi i \hbar \sin w(t_2 - t_1)}} \exp \left\{ \frac{imw}{2\hbar \sin w(t_2 - t_1)} \right. \\
 &\quad \times \left[\cos w(t_2 - t_1)(\xi_2^2 + \xi_1^2) - 2\xi_2 \xi_1 + \frac{2\xi_2}{mw} \int_{t_1}^{t_2} f(t) \sin w(t - t_1) dt \right]
 \end{aligned}$$

$$\begin{aligned}
 & + \frac{2\xi_1}{mw} \int_{t_1}^{t_2} f(t) \sin w(t_2 - t) dt \\
 & - \frac{2}{m^2 w^2} \int_{t_1}^{t_2} \int_{t_1}^{t_2} f(s) f(t) \sin w(t_2 - t) \sin w(s - t_1) ds dt \Big\}.
 \end{aligned}$$

PROOF. It can be shown that for $(\xi_2, t_2, \xi_1, t_1, q) \in \mathbf{R} \times [0, T] \times \mathbf{R} \times [0, T] \times (\mathbf{R} - \{0\})$ with $t_1 < t_2$

(3.21)

$$\Gamma(\xi_2, t_2; \xi_1, t_1; -iq) = \sqrt{\frac{q}{2\pi it}} \exp\left\{-\frac{q(\xi_2 - \xi_1)^2}{2it}\right\} E^{\text{anf}_q}[F|X = (\xi_1, \xi_2)]$$

satisfies the Schrödinger equation

$$\frac{\partial \Gamma}{\partial t} = \frac{i}{2q} \frac{\partial^2 \Gamma}{\partial \xi^2} - \frac{\alpha}{2} \xi^2 \Gamma + \alpha \beta \xi q(t) \Gamma$$

with the condition $\Gamma(\xi_2, t_2; \xi_1, t_1; -iq) \rightarrow \delta(\xi_2 - \xi_1)$ as $t_2 \downarrow t_1$. But by using (3.8) in Theorem 3.3 with $q = \frac{m}{\hbar}$, $\alpha = \frac{imw^2}{\hbar}$ and $\beta = \frac{1}{mw^2}$ we can easily obtain (3.20) from (3.21). Hence (3.20) is the solution of (3.19).

Let S be the operator on $L^2[0, T]$ defined by

$$(3.22) \quad Sv(s) = \int_s^T v(u) du, \quad s \in [0, T], \quad v \in L^2[0, T].$$

Then S is a bounded linear operator from $L^2[0, T]$ to $L^2[0, T]$.

THEOREM 3.7. Let F be a measurable function on $C_0[0, T]$ defined by

$$F(x) = \int_{L^2[0, T]} \exp\left\{i \int_{t_1}^T v(s)x(s) ds\right\} d\sigma(v), \quad v \in L^2[0, T],$$

where $\sigma \in M(L^2[0, T])$, the space of the \mathbf{C} -valued, countably additive Borel measure on $L^2[0, T]$. Let $X(x) = (x(t_1), x(T))$. Then for all $\lambda \in$

\mathbf{C}^+ , the conditional analytic Wiener integral $E^{\text{anw}}(F|X)$ exists and for all $(\xi_1, \xi) \in \mathbf{R}^2$ is given by the formula

$$(3.23) \quad \begin{aligned} & E^{\text{anw}\lambda}[F|X = (\xi_1, \xi)] \\ &= \int_{L^2[0,T]} \exp \left\{ -\frac{1}{2\lambda} \int_{t_1}^T \left(\int_s^T v(u)du \right)^2 ds + \frac{1}{2\lambda(T-t_1)} \left(\int_{t_1}^T \int_s^T v(u)duds \right)^2 \right. \\ & \quad \left. + i \frac{\xi - \xi_1}{T-t_1} \int_{t_1}^T \int_s^T v(u)duds \right\} d\sigma(v). \end{aligned}$$

Furthermore, $E^{\text{anf}_q}(F|X)$ exists for all real $q \neq 0$ and for all $(\xi_1, \xi) \in \mathbf{R}^2$ is given by the formula

$$(3.24) \quad \begin{aligned} & E^{\text{anf}_q}[F|X = (\xi_1, \xi)] \\ &= \int_{L^2[0,T]} \exp \left\{ -\frac{i}{2q} \int_{t_1}^T \left(\int_s^T v(u)du \right)^2 ds + \frac{i}{2q(T-t_1)} \left(\int_{t_1}^T \int_s^T v(u)duds \right)^2 \right. \\ & \quad \left. + i \frac{\xi - \xi_1}{T-t_1} \int_{t_1}^T \int_s^T v(u)duds \right\} d\sigma(v). \end{aligned}$$

PROOF. We note that

$$\begin{aligned} F(x) &= \int_{L^2[0,T]} \exp \left\{ i \int_{t_1}^T v(s)x(s)ds \right\} d\sigma(v) \\ &= \int_{L^2[0,T]} \exp \left\{ i \int_{t_1}^T Sv(s)dx(s) \right\} d\sigma(v) \\ &= \int_{L^2[0,T]} \exp \left\{ i \int_{t_1}^T w(s)dx(s) \right\} dc S^{-1}(w). \end{aligned}$$

Using Proposition 2.1, Fubini theorem and the change of variable formula we obtain that for $\lambda > 0$ and all $(\xi_1, \xi) \in \mathbf{R}^2$

$$(3.25) \quad \begin{aligned} & E[F(\lambda^{-\frac{1}{2}} \cdot |X(\lambda^{-\frac{1}{2}} \cdot) = (\xi_1, \xi)] \\ &= \int_{C_0^0[0, T-t_1]} \int_{L^2[0,T]} \exp \left\{ i \int_0^{T-t_1} w(s+t_1)d(\lambda^{-\frac{1}{2}}y + g)(s) \right\} d\sigma S^{-1}(w) dm_p(y) \\ &= \int_{L^2[0,T]} \exp \left\{ i \frac{\xi - \xi_1}{T-t_1} \int_0^{T-t_1} w(s+t_1)ds \right\} \\ & \quad \cdot \int_{C_0^0[0, T-t_1]} \exp \left\{ \frac{i}{\sqrt{\lambda}} \int_0^{T-t_1} w(s+t_1)dy(s) \right\} dm_p(y) d\sigma S^{-1}(w) \end{aligned}$$

$$\begin{aligned}
 &= \int_{L^2[0,T]} \exp \left\{ i \frac{\xi - \xi_1}{T - t_1} \int_0^{T-t_1} w(s + t_1) ds \right\} \left[\frac{1}{\sqrt{2\pi\rho^2}} \int_{\mathbf{R}} \exp \left\{ \frac{i}{\sqrt{\lambda}} u \right\} \right. \\
 &\quad \left. \cdot \exp \left\{ -\frac{u^2}{2\rho^2} \right\} du \right] d\sigma S^{-1}(w) \\
 &= \int_{L^2[0,T]} \exp \left\{ i \frac{\xi - \xi_1}{T - t_1} \int_0^{T-t_1} w(s + t_1) ds \right\} \exp \left\{ -\frac{\rho^2}{2\lambda} \right\} d\sigma S^{-1}(w) \\
 &= \int_{L^2[0,T]} \exp \left\{ -\frac{1}{2\lambda} \int_{t_1}^T w(s)^2 ds + \frac{1}{2\lambda(T - t_1)} \left(\int_{t_1}^T w(s) ds \right)^2 \right. \\
 &\quad \left. + i \frac{\xi - \xi_1}{T - t_1} \int_{t_1}^T w(s) ds \right\} d\sigma S^{-1}(w) \\
 &= \int_{L^2[0,T]} \exp \left\{ -\frac{1}{2\lambda} \int_{t_1}^T Sv(s)^2 ds + \frac{1}{2\lambda(T - t_1)} \left(\int_{t_1}^T Sv(s) ds \right)^2 \right. \\
 &\quad \left. + i \frac{\xi - \xi_1}{T - t_1} \int_{t_1}^T Sv(s) ds \right\} d\sigma(v),
 \end{aligned}$$

where $\rho^2 = \int_{t_1}^T w(s)^2 ds - \frac{1}{T-t_1} \left(\int_{t_1}^T w(s) ds \right)^2$ is the variance of $\int_0^{T-t_1} w(s + t_1) dy(s)$. Since $\sigma \in M(L_2[0, T])$, we see that the last expression on the right-hand side of (3.25) is an analytic function of λ throughout \mathbf{C}^+ and is a continuous function λ for $\text{Re}\lambda \geq 0, \lambda \neq 0$. Thus (3.23) and (3.24) are established, which concludes the proof of Theorem 3.7.

THEOREM 3.8. *Let F be a measurable function on $C_0[0, T]$ defined by*

$$F(x) = \exp \left\{ -\frac{i}{2} \int_{t_1}^T x^2(s) ds \right\} \int_{L^2[0,T]} \exp \left\{ i \int_{t_1}^T v(s)x(s) ds \right\} d\sigma(v),$$

where $x \in C_0[0, T], v \in L^2[0, T]$ and $\sigma \in M(L^2[0, T])$. Let $X(x) = (x(t_1), x(T))$. Then for all $\lambda \in \mathbf{C}^+$, the conditional analytic Feynman integral $E^{\text{anf}_q}(F|X)$ exists and for all $(\xi_1, \xi) \in \mathbf{R}^2$ and real $q \neq 0$ is given by the formula

$$\begin{aligned}
 (3.26) \quad E^{\text{anf}_q}[F|X = (\xi_1, \xi)] &= \left(\frac{\sqrt{-\frac{1}{q}}(T - t_1)}{\sinh \sqrt{-\frac{1}{q}}(T - t_1)} \right)^{\frac{1}{2}} \cdot \exp \left\{ \frac{q(\xi - \xi_1)^2}{2i(T - t_1)} \right\} \\
 &\quad \cdot \exp \left\{ -\frac{\sqrt{q}}{2} (\xi^2 + \xi_1^2) \coth \sqrt{-\frac{1}{q}}(T - t_1) + \frac{\sqrt{q}\xi\xi_1}{\sinh \sqrt{-\frac{1}{q}}(T - t_1)} \right\}
 \end{aligned}$$

$$\int_{L_2[0,T]} \exp \left\{ i \int_{t_1}^T \frac{v(t)}{\sinh \sqrt{-\frac{1}{q}}(T-t_1)} \left(\xi \sinh \sqrt{-\frac{1}{q}}(t-t_1) + \xi_1 \sinh \sqrt{-\frac{1}{q}}(T-t) \right) dt \right\} \cdot \exp \left\{ -\frac{i}{2q} \int_0^{T-t_1} \int_0^{T-t_1} R\left(s, t, -\frac{1}{q}\right) v(s+t_1)v(t-t_1) ds dt \right\} d\sigma(v).$$

PROOF. By Proposition 2.1, we have, for $\xi, \xi_1 \in \mathbf{R}$ and $\lambda > 0$

$$\begin{aligned} E[F(\lambda^{-\frac{1}{2}} \cdot) | \lambda^{-\frac{1}{2}} x(t_1) = \xi_1, \lambda^{-\frac{1}{2}} x(T) = \xi] \\ = \int_{C_0^0[0, T-t_1]} \exp \left\{ -\frac{i}{2\lambda} \int_0^{T-t_1} (y(s) + \sqrt{\lambda}g(s))^2 ds \right\} \\ \int_{L_2[0,T]} \exp \left\{ \frac{i}{\sqrt{\lambda}} \int_0^{T-t_1} v(s+t_1)(y(s) + \sqrt{\lambda}g(s)) ds \right\} d\sigma(v) dm_p(y) \\ = \int_{L_2[0,T]} \int_{C_0^0[0, T-t_1]} \exp \left\{ -\frac{i}{2\lambda} \int_0^{T-t_1} (y(s) + \sqrt{\lambda}g(s))^2 ds \right. \\ \left. + \frac{i}{\sqrt{\lambda}} \int_0^{T-t_1} v(s+t_1)(y(s) + \sqrt{\lambda}g(s)) ds \right\} dm_p(y) d\sigma(v) \end{aligned}$$

where $g(t) = \frac{\xi - \xi_1}{T - t_1}t + \xi_1, t \in [0, T - t_1]$ and the second equality is justified by the Fubini theorem. But we know that the integrand of integral on $L_2[0, T]$ in the last equality is given by the right side of (3.7) with $\lambda > 0, \alpha = i$ and $\beta = 1$. Since σ is in $M(L_2[0, T])$ and the right side of (3.7) is an analytic function of λ through \mathbf{C}^+ and a continuous function for $\text{Re}\lambda \geq 0, \lambda \neq 0$, we obtain, for $\lambda \in \mathbf{C}^+$ and $(\xi_1, \xi) \in \mathbf{R}^2$

$$\begin{aligned} (3.27) \quad E^{\text{an}\lambda}[F|X = (\xi_1, \xi)] &= \left(\frac{\sqrt{\frac{i}{\lambda}}(T-t_1)}{\sinh \sqrt{\frac{i}{\lambda}}(T-t_1)} \right)^{\frac{1}{2}} \cdot \exp \left\{ \frac{\lambda(\xi - \xi_1)^2}{2(T-t_1)} \right\} \\ &\cdot \exp \left\{ -\frac{\sqrt{i\lambda}}{2}(\xi^2 + \xi_1^2) \coth \sqrt{\frac{i}{\lambda}}(T-t_1) + \frac{\sqrt{i\lambda}\xi\xi_1}{\sinh \sqrt{\frac{i}{\lambda}}(T-t_1)} \right\} \end{aligned}$$

$$\int_{L_2[0,T]} \exp \left\{ i \int_{t_1}^T \frac{v(t)}{\sinh \sqrt{\frac{i}{\lambda}}(T-t_1)} (\xi \sinh \sqrt{\frac{i}{\lambda}}(t-t_1) + \xi_1 \sinh \sqrt{\frac{i}{\lambda}}(T-t)) dt \right\} \cdot \exp \left\{ -\frac{1}{2\lambda} \int_0^{T-t_1} \int_0^{T-t_1} R\left(s, t, \frac{i}{\lambda}\right) v(s+t_1)v(t+t_1) ds dt \right\} d\sigma(v).$$

Hence as $\lambda \rightarrow -iq, q \neq 0$, the result (3.26) is obtained as desired.

4. Correction to “A class of conditional Wiener integrals”

In the first author’s paper[4] with S.J.Chang, the hypothesis for Theorem 2.2 should have been : Let F be an integrable function on $C_0[0, T]$ defined by $F(x) = \exp\{\int_{t_1}^T V(s, x(s))ds\}$.

The expression appearing in Theorem 3.3 and its proof :

$$\alpha\beta(\xi - \xi_1) \int_0^{T-t_1} \left(\frac{\sinh\sqrt{\alpha}t}{\sinh\sqrt{\alpha}(T-t_1)} + \frac{\xi_1}{\xi - \xi_1} \right) q(t+t_1)dt$$

should read as

$$\alpha\beta \left[\int_0^{T-t_1} \frac{q(t+t_1)}{\sinh \sqrt{\alpha}(T-t_1)} (\xi \sinh \sqrt{\alpha}t + \xi_1 \sinh \sqrt{\alpha}(T-t_1-t))dt \right].$$

References

1. A. Alberverio and R. Hoegh-Krohn, *Mathematical theory of Feynman path integral*, Lecture notes in Math., **523**, Springer-Verlag, Berlin (1976).
2. R. B. Ash, *Topics in Stochastic Processes*, Academic Press, New York, 1975.
3. R. H. Cameron and Storvick, *Some Banach algebras of analytic Feynman integrable functionals, in analytic functions*, Kozubnik, 1979; vol. 1980, Springer-Verlag, Berlin, New York.
4. S. J. Chang and D. M. Chung, *A class of conditional Wiener integrals*, J. of the Korean Math. Soc. **30** (1993), 161-172.
5. D. M. Chung, *Scale-invariant measurability in abstract Wiener spaces.*, Pacific J. Math. **130** (1987), 27-40.
6. D. M. Chung and D. Skoug, *Conditional analytic Feynman integrals and a related Schrödinger integral equation*, SIAM J.Math. Anal. **20** (1989), 950-965.

7. D. M. Chung, C. Park and D. Skoug, *Operator-valued Feynman integral via conditional Feynman integrals*, Pacific J. Math. **146** (1990), 21-42.
8. M. D. Donsker and J. L. Lions, *Volterra variational equations, boundary value problems and function space integrals*, Acta Math. **109** (1962), 147-228.
9. R. P. Feynman and A. R. Hibbs, *Quantum mechanics and path integrals*, McGraw-Hill, New York, 1965.
10. G. W. Johnson, *The equivalence of two approaches to the Feynman integral*, J. Math. Phys. **23** (1982), 2090-2096.
11. G. W. Johnson and D. L. Skoug, *Scale-invariant measurability in Wiener space*, Pacific J. Math. **83** (1979), 157-176.
12. M. Kac, *Probability and Related Topics in Physical Sciences*, Interscience Publishers, New York, 1959.
13. C. Park and D. Skoug, *A simple formula for conditional Wiener integrals with applications*, Pacific J. Math. **135** (1988), 381-394.
14. J. Yeh, *Inversion of Conditional Wiener Integrals*, Pacific J. Math. **59** (1975), 623-638.

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