

**WEIGHTED ORLICZ SPACE  
INTEGRAL INEQUALITIES FOR  
POTENTIAL MAXIMAL OPERATORS**

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ABSTRACT. We characterize a condition for  $\mathcal{M}_\rho$  to be of weak type  $(\Phi_1, \Phi_2)$  in terms of Orlicz norms.

**1. Introduction**

Given a function  $f$  in  $\mathbb{R}^n$ , we define a function  $\mathcal{M}f$  in  $\overline{\mathbb{R}^{n+1}} = \{(x, s) : x \in \mathbb{R}^n, s \geq 0\}$  by setting

$$\mathcal{M}f(x, s) = \sup \left\{ \frac{1}{|Q|} \int_Q |f(y)| dy : x \in Q \text{ and } \text{sidelength}(Q) \geq s \right\}.$$

It is well known that this maximal operator  $\mathcal{M}$  controls the Poisson integral defined by, for  $x \in \mathbb{R}^n$  and  $s \geq 0$ ,

$$P(f)(x, s) = \int_{\mathbb{R}^n} f(y)P(x - y, s) dy,$$

where

$$P(x, s) = \frac{c_n s}{(|x|^2 + s^2)^{n+1/2}}$$

is the Poisson kernel. For a given positive measure  $\nu$  on  $\mathbb{R}^n \times [0, \infty)$ , under what condition on  $\nu$  can we assert that  $\mathcal{M}$  is bounded from  $\mathcal{L}^p(\mathbb{R}^n)$  into weak- $\mathcal{L}^p(\mathbb{R}^n \times [0, \infty), \nu)$ ? Carleson([3]) showed that this was equivalent to the Carleson condition and later Feffermann-Stein([5]) found a sufficient condition, and Ruiz([13]), Ruiz-Torrea([14]) unified all these

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results. Recently, Gallardo([6]) and Chen([4]) obtained characterizations in terms of the Orlicz norm and in [8], we obtained a characterization for the fractional maximal operator. In this paper, we characterize a condition for  $\mathcal{M}_\varphi$  to be of weak type  $(\Phi_1, \Phi_2)$  having four weights in the Orlicz norm. In the next theorem, we shall assume that  $\varphi$  is essentially nondecreasing, i.e., there exists a positive constant  $\rho$  for which

$$\varphi(t) \leq \rho\varphi(s), \quad t \leq s$$

and

$$\lim_{t \rightarrow \infty} \frac{\varphi(t)}{t} = 0.$$

Our result is as follows.

**THEOREM 3.1.** *Let  $\mathcal{M}_\varphi f$  be the potential maximal operator on  $\overline{\mathbb{R}_+^{n+1}}$ .  $\mathcal{M}_\varphi f$  is defined by*

$$(1) \quad \mathcal{M}_\varphi f(x, s) = \sup_{x \in Q} \frac{\varphi(|Q|)}{|Q|} \int_Q |f(y)| dy, \quad l(Q) \geq s,$$

where  $l(Q)$  denotes the sidelength of  $Q$ . Let  $u, v$  be weights on  $\mathbb{R}^n$ ,  $w$  be a weight on  $\overline{\mathbb{R}_+^{n+1}}$  and  $\mu$  be a nonnegative measure on  $\overline{\mathbb{R}_+^{n+1}}$ .  $\Phi_1$  and  $\Phi_2$  are  $N$ -functions with complementary  $N$ -functions  $\Psi_1$  and  $\Psi_2$ , respectively. Assume further that  $\Phi_2 \circ \Phi_1^{-1}$  is convex. Then weak type boundedness, i.e.,

$$(2) \quad \Phi_2^{-1} \left[ \int_{\{(x,s) \in \overline{\mathbb{R}_+^{n+1}} : \mathcal{M}_\varphi f(x,s) > \lambda\}} \Phi_2(\lambda w(x,s)) d\mu(x,s) \right] \leq \Phi_1^{-1} \left[ \int_{\mathbb{R}^n} \Phi_1(C|f(x)|u(x))v(x) dx \right]$$

holds if and only if

$$(3) \quad \int_Q \Psi_1 \left[ \frac{\gamma(\lambda, \tilde{Q})}{C\lambda u(y)v(y)} \frac{\varphi(|Q|)}{|Q|} \right] v(y) dy \leq \gamma(\lambda, \tilde{Q}) < \infty$$

holds for each cube  $Q$ , where

$$(4) \quad \gamma(\lambda, \tilde{Q}) = \Phi_1 \circ \Phi_2^{-1} \left[ \int_{\tilde{Q}} \Phi_2(\lambda w(x,s)) d\mu(x,s) \right], \quad \tilde{Q} = Q \times (0, l(Q)].$$

When  $\varphi(|Q|) = 1$ , the Hardy-Littlewood maximal operator is obtained. The fractional maximal operator  $\mathcal{M}_\alpha$  ( $0 < \alpha < n$ ) is given by  $\varphi(|Q|) = |Q|^{\frac{\alpha}{n}}$ . Maximal operators connected to the Bessel potential operator are defined by  $\varphi(|Q|) = \int_0^{|Q|^{\frac{1}{n}}} \psi(s) ds$ , where  $\psi$  is the derivative of  $\varphi$ .

### 2. Preliminaries

DEFINITION 2.1. Let  $\Phi : [0, \infty) \rightarrow \mathbb{R}$  be a function satisfying

- (i)  $\Phi(s) > 0$  for all  $s \geq 0$ ;
- (ii)  $\lim_{s \rightarrow 0} \Phi(s)/s = 0$ ;
- (iii)  $\lim_{s \rightarrow \infty} \Phi(s)/s = \infty$ .

Then  $\Phi$  is called an  $N$ -function. Each  $N$ -function has the integral representation:  $\Phi(s) = \int_0^s \phi(t) dt$ , where  $\phi(s) > 0$  for  $s > 0$ ,  $\phi(0) = 0$  and  $\phi(s) \rightarrow \infty$  as  $s \rightarrow \infty$ . Further,  $\phi$  is right-continuous and nondecreasing.  $\phi$  is called the density function of  $\Phi$ . Define  $\rho : [0, \infty) \rightarrow \mathbb{R}$  by  $\rho(t) = \sup\{s : \phi(s) \leq t\}$ . Then  $\rho$  is called the generalized inverse of  $\phi$ . Finally, define

$$\Psi(t) = \int_0^t \rho(s) ds$$

and  $\Psi$  is called the complementary  $N$ -function of  $\Phi$ . For further details, see [10].

DEFINITION 2.2. An  $N$ -function  $\Phi$  is said to satisfy the  $\Delta_2$ -condition in  $[0, \infty)$  if  $\sup_{s>0} \frac{\Phi(2s)}{\Phi(s)} < \infty$ .

REMARK 1. If  $\phi$  is the density function of  $\Phi$ , then  $\Phi$  satisfies the  $\Delta_2$ -condition if and only if there exists a constant  $\alpha > 1$  such that  $s\phi(s) < \alpha\Phi(s)$ , for any  $s > 0$ .

REMARK 2. If  $\Psi$  is the complementary  $N$ -function of  $\Phi$ , then  $st \leq \Phi(s) + \Psi(t)$  for all  $s, t \geq 0$ . Further,  $=$  holds if and only if  $\phi(s-) \leq t \leq \phi(s)$  or else  $\rho(t-) \leq s \leq \rho(t)$ .

DEFINITION 2.3. Let  $(X, \mathcal{M}, \mu)$  is a  $\sigma$ -finite measure space and  $\Phi$  be an  $N$ -function. Then the Orlicz space  $\mathcal{L}_\Phi(d\mu)$  and  $\mathcal{L}_\Phi^*(d\mu)$  are defined by

$$\mathcal{L}_\Phi(d\mu) = \left\{ f : \int_X \Phi(|f|) d\mu < \infty \right\}$$

and

$$\mathcal{L}_\Phi^*(d\mu) = \{ f : fg \in \mathcal{L}_1(d\mu) \text{ for all } g \in \mathcal{L}_\Psi \},$$

where  $\Psi$  is the complementary  $N$ -function of  $\Phi$ .

Keeping these definitions and notions, the following properties about the Orlicz space will be used in the proof of the Theorem 3.1.

PROPOSITION 2.4.

(i) The Orlicz space  $\mathcal{L}_\Phi^*(d\mu)$  is a Banach space with the Orlicz norm

$$\|f\|_\Phi = \sup \left\{ \int |fg| d\mu : g \in \mathcal{S}_\Psi \right\},$$

where  $\mathcal{S}_\Psi = \{g \in \mathcal{L}_\Psi : \int \Psi(|g|) d\mu \leq 1\}$ , or with the Luxemburg norm

$$\|f\|_{(\Phi)} = \inf \left\{ \lambda > 0 : \int \Phi\left(\frac{|f|}{\lambda}\right) d\mu \leq 1 \right\}.$$

(ii) (Hölder's inequality) If  $f \in \mathcal{L}_\Phi^*(d\mu)$  and  $g \in \mathcal{L}_\Psi^*(d\mu)$ , then

$$(5) \quad \|fg\|_\Phi \leq 2\|f\|_{(\Phi)}\|g\|_{(\Psi)}.$$

(iii) (Young's inequality)

$$(6) \quad ab \leq \Phi(a) + \Psi(b) \quad \text{for all } a, b > 0.$$

LEMMA 2.5. Let  $\Phi$  be an  $N$ -function with complementary function  $\Psi$ . Let  $x$  and  $y > 0$ . Then

$$(7) \quad \Phi(x) \leq x\phi(x) \leq \Phi(2x),$$

$$(8) \quad \Phi(x) + \Phi(y) \leq \Phi(x + y)$$

and

$$(9) \quad \Phi\left[\frac{\Psi(x)}{x}\right] \leq \Psi(x).$$

LEMMA 2.6. Suppose that  $f$  is an integrable function in  $\mathbb{R}^n$ . For each  $\lambda > 0$ , let  $E_\lambda = \{(x, s) \in \mathbb{R}_+^{n+1} : M_\varphi f(x, s) > \lambda\}$ . Then, if  $E_\lambda$  is not empty, we have

$$(10) \quad E_\lambda \subset \bigcup_j \widetilde{Q}_j^3,$$

where  $Q_j$  is the family of nonoverlapping maximal dyadic cubes satisfying

$$(11) \quad \frac{\lambda}{4^{n\rho}} < \frac{\varphi(|Q_j|)}{|Q_j|} \int_{Q_j} f(y) dy \leq \frac{\lambda}{2^n}$$

for each integer  $j$ . Furthermore, we have that

$$\left\{ x \in \mathbb{R}^n : M_\varphi^d f(x) > \frac{\lambda}{4^{n\rho}} \right\} = \bigcup_j Q_j.$$

*Proof.* Following [7](p.160), we let  $C_\lambda = \{P_j\}$  be the family of the dyadic maximal nonoverlapping cubes satisfying the condition

$$\lambda < \frac{\varphi(|P_j|)}{|P_j|} \int_{P_j} f(y) dy.$$

To show that there is such a family  $C_\lambda$ , observe that

$$\frac{\varphi(|Q|)}{|Q|} \int_Q f(y) dy \rightarrow 0$$

as  $Q \uparrow \mathbb{R}^n$ , since  $f$  is integrable and since  $\lim_{t \rightarrow \infty} \varphi(t)/t = 0$ . If, for some dyadic cube  $Q$ ,

$$\lambda < \frac{\varphi(|Q|)}{|Q|} \int_Q f(y) dy,$$

then  $Q$  is contained in dyadic cubes satisfying this condition, which are maximal with respect to the inclusion. Thus, there is a family of maximal nonoverlapping dyadic cubes  $\{P_j\}$  yield

$$(12) \quad \lambda < \frac{\varphi(|P_j|)}{|P_j|} \int_{P_j} f(y) dy \leq 2^{n\rho} \frac{\varphi(|P'_j|)}{|P'_j|} \int_{P'_j} f(y) dy \leq 2^n \rho \lambda,$$

where  $P'_j$  denotes the only dyadic cube containing  $P_j$ . From this discussion, it is clear that

$$\{x \in \mathbb{R}^n : M_\varphi^d f(x) > \lambda\} = \bigcup_j P'_j.$$

Let  $(x, s) \in E_\lambda$ ; by definition, there is a cube  $R$  containing  $x$  with  $l(R) \geq s$  such that

$$\lambda < \frac{\varphi(|R|)}{|R|} \int_R f(y) dy.$$

Let  $k$  be the unique integer such that  $2^{(k+1)n} < |R| \leq 2^{-kn}$ . There are some dyadic cubes with side length  $2^{-k}$ , and at most  $2^n$  of them,  $\{J_i : i = 1, \dots, 2^n\}$  meeting the interior of  $R$ . It is easy to see that, for one of these cubes, say  $J_1$ ,

$$\frac{\lambda}{2^n} < \frac{\varphi(|R|)}{|R|} \int_{R \cap J_1} f(y) dy.$$

Now, since  $|R| \leq |J_1| < 2^n |R|$ ,  $\varphi(|R|) \leq \rho \varphi(|J_1|)$  and

$$\frac{\lambda}{4^n} |J_1| < \frac{\lambda}{2^n} |R| < \varphi(|R|) \int_{R \cap J_1} f(y) dy \leq \rho \varphi(|J_1|) \int_{J_1} f(y) dy.$$

Hence,

$$\frac{\lambda}{4^n \rho} < \frac{\varphi(|J_1|)}{|J_1|} \int_{J_1} f(y) dy.$$

By letting  $C_{t/4^n \rho} = \{Q_j\}$ , we see that  $J_1 \subset Q_k$ , for some  $k$ , and  $x \in R \subset J_1^3 \subset Q_k^3$ .

On the other hand,  $s \leq l(R) \leq l(Q_j^3)$ . From this, we conclude that

$$E_\lambda \subset \bigcup_j \widetilde{Q}_j^3.$$

Finally, it follows from (12) that

$$\frac{\lambda}{4^n \rho} < \frac{\varphi(|Q_j|)}{|Q_j|} \int_{Q_j} f(y) dy \leq \frac{\lambda}{2^n},$$

for each  $j$ , concluding the proof of the lemma.  $\square$

### 3. Main Results

**THEOREM 3.1.** *Let  $\mathcal{M}_\varphi f$  be the potential maximal operator on  $\overline{\mathbb{R}^{n+1}_+}$ .  $\mathcal{M}_\varphi f$  is defined by*

$$\mathcal{M}_\varphi f(x, s) = \sup_{x \in Q} \frac{\varphi(|Q|)}{|Q|} \int_Q |f(y)| dy, \text{ where } l(Q) \geq s.$$

Let  $u$  and  $v$  be weights on  $\mathbb{R}^n$ ,  $w$  be a weight on  $\overline{\mathbb{R}^{n+1}_+}$  and  $\mu$  be a nonnegative measure on  $\overline{\mathbb{R}^{n+1}_+}$ .  $\Phi_1$  and  $\Phi_2$  are  $N$ -functions with complements  $\Psi_1$  and  $\Psi_2$ , respectively. Assume further that  $\Phi_2 \circ \Phi_1^{-1}$  is convex. Then weak type boundedness, i.e.

$$\Phi_2^{-1} \left[ \int_{\{(x,s) \in \overline{\mathbb{R}^{n+1}_+} : \mathcal{M}_\varphi f(x,s) > \lambda\}} \Phi_2(\lambda w(x,s)) d\mu(x,s) \right] \leq \Phi_1^{-1} \left[ \int_{\mathbb{R}^n} \Phi_1(C|f(y)|u(y))v(y) dy \right]$$

holds if and only if

$$\int_Q \Psi_1 \left[ \frac{\gamma(\lambda, \tilde{Q})}{C\lambda u(x)v(x)} \frac{\varphi(|Q|)}{|Q|} \right] v(x) dx \leq \gamma(\lambda, \tilde{Q}) < \infty$$

holds for each cube  $Q$ , where

$$\gamma(\lambda, \tilde{Q}) = \Phi_1 \circ \Phi_2^{-1} \left[ \int_{\tilde{Q}} \Phi_2(\lambda w(x,s)) d\mu(x,s) \right], \quad \tilde{Q} = Q \times (0, l(Q)].$$

*Proof.* For the necessity, we follow the idea of [2]. Since  $\frac{\Psi_1(\varepsilon)}{\varepsilon}$  is increasing in  $\varepsilon$  and has full range  $\mathbb{R}^+$ , for given  $\lambda > 0$ , we can choose  $\varepsilon$  such that

$$\int_E \Psi_1 \left( \frac{\varepsilon}{u(y)v(y)} \right) \frac{v(y)}{\varepsilon} dy = 2C\lambda \frac{|Q|}{\varphi(|Q|)}, \quad E \subset Q$$

$$f(x) = \frac{1}{C} \Psi_1 \left( \frac{\varepsilon}{u(x)v(x)} \right) \frac{v(x)}{\varepsilon} \chi_E(x).$$

If  $(y, s) \in \tilde{Q}$ , then  $y \in Q$  and  $s \leq l(Q)$ . So

$$\begin{aligned} \mathcal{M}_\varphi f(y, s) &= \sup_{y \in Q} \frac{\varphi(|Q|)}{|Q|} \int_Q |f(y)| dy \\ &\geq \frac{\varphi(|Q|)}{|Q|} \int_Q |f(y)| dy \\ &= \frac{\varphi(|Q|)}{|Q|} \int_Q \frac{1}{C} \Psi_1\left(\frac{\varepsilon}{u(y)v(y)}\right) \frac{v(y)}{\varepsilon} \chi_E(y) dy \\ &= 2\lambda > \lambda. \end{aligned}$$

Thus if  $(y, s) \in \tilde{Q}$ , then  $(y, s) \in E_\lambda = \{(x, r) \in \mathbb{R}_+^{n+1} : \mathcal{M}_\varphi f(x, r) > \lambda\}$ . Hence

$$\begin{aligned} \gamma(\lambda, \tilde{Q}) &= \Phi_1 \circ \Phi_2^{-1} \left[ \int_{\tilde{Q}} \Phi_2(\lambda w(x, s)) d\mu(x, s) \right] \\ &\leq \Phi_1 \circ \Phi_2^{-1} \left[ \int_{\{\mathcal{M}_\varphi f > \lambda\}} \Phi_2(\lambda w(x, s)) d\mu(x, s) \right] \\ &\leq \int_E \Phi_1 \left[ \frac{u(y)v(y)}{\varepsilon} \Psi_1\left(\frac{\varepsilon}{u(y)v(y)}\right) \right] v(y) dy \\ &\leq \int_E \Psi_1\left(\frac{\varepsilon}{u(y)v(y)}\right) v(y) dy \\ &= 2C\lambda \frac{|Q|}{\varphi(|Q|)} \varepsilon. \end{aligned}$$

The third inequality follows from (2). Put

$$\Delta = \int_E \psi_1 \left( \frac{\gamma(\lambda, \tilde{Q})}{4C\lambda u(y)v(y)} \frac{\varphi(|Q|)}{|Q|} \right) \frac{1}{u(y)} dy.$$

Then

$$\begin{aligned} \Delta &\leq \int_E \psi_1 \left( \frac{\varepsilon}{2u(y)v(y)} \right) \frac{1}{u(y)} dy \\ &\leq 2 \int_E \Psi_1 \left( \frac{\varepsilon}{u(y)v(y)} \right) \frac{v(y)}{\varepsilon} dy \\ &= 4C\lambda \frac{|Q|}{\varphi(|Q|)}. \end{aligned}$$



Also

$$\Delta \geq \frac{4C\lambda|Q|}{\varphi(|Q|)\gamma(\lambda, \tilde{Q})} \int_E \Psi_1 \left[ \left( \frac{\gamma(\lambda, \tilde{Q})}{4C\lambda u(y)v(y)} \right) \frac{\varphi(|Q|)}{|Q|} \right] v(y) dy.$$

Thus

$$\int_E \Psi_1 \left[ \left( \frac{\gamma(\lambda, \tilde{Q})}{4C\lambda u(y)v(y)} \right) \frac{\varphi(|Q|)}{|Q|} \right] v(y) dy \leq \gamma(\lambda, \tilde{Q}) < \infty.$$

As  $E \rightarrow Q$ , (3) holds.

For the sufficiency, we follow the idea of [4] and [11]. From Lemma 2.6, for each  $\lambda > 0$ , let  $E_\lambda = \{(x, s) \in \mathbb{R}_+^{n+1} : M_\varphi f(x, s) > \lambda\}$ . Then, if  $E_\lambda$  is not empty, we have

$$E_\lambda \subset \bigcup_j \tilde{Q}_j^3,$$

where  $Q_j$  is the family of nonoverlapping maximal dyadic cubes satisfying

$$\frac{\lambda}{4^{n\rho}} < \frac{\varphi(|Q_j|)}{|Q_j|} \int_{Q_j} f(y) dy \leq \frac{\lambda}{2^n}$$

for each integer  $j$ . Then  $\frac{\lambda}{4^{n\rho}} < \frac{\varphi(|Q_j|)}{|Q_j|} \int_{Q_j} f(y) dy$  and  $\{\tilde{Q}_j^3\}$  is a covering of  $E_\lambda$ . Hence it follows from (6) that

$$\begin{aligned} 2\gamma(\lambda, \tilde{Q}_j^3) &\leq \int_{Q_j} \frac{4^{n\rho}|f(x)|}{\lambda} \frac{\varphi(|Q_j|)}{|Q_j|} 2\gamma(\lambda, \tilde{Q}_j^3) dx \\ &= \int_{Q_j} 23^n 4^n C\rho |f(x)|u(x) \frac{\gamma(\lambda, \tilde{Q}_j^3)}{C\lambda u(x)v(x)} \frac{\varphi(|Q_j^3|)}{|Q_j^3|} v(x) dx \\ &\leq \int_{Q_j} \Phi_1(23^n 4^n C\rho |f(x)|u(x))v(x) dx \\ &\quad + \int_{Q_j} \Psi_1 \left( \frac{\gamma(\lambda, \tilde{Q}_j^3)}{C\lambda u(x)v(x)} \frac{\varphi(|Q_j^3|)}{|Q_j^3|} \right) v(x) dx \\ &\leq \int_{Q_j} \Phi_1(23^n 4^n \rho C |f(x)|u(x))v(x) dx + \gamma(\lambda, \tilde{Q}_j^3). \end{aligned}$$

So

$$\gamma(\lambda, \widetilde{Q}_j^3) \leq \int_{Q_j} \Phi_1(23^n 4^n C \rho |f(x)|u(x))v(x) dx$$

and thus

$$\int_{\widetilde{Q}_j^3} (\Phi_2(\lambda w(x, s))) d\mu(x, s) \leq \Phi_2 \circ \Phi_1^{-1} \left[ \int_{Q_j} \Phi_1(23^n 4^n C \rho |f(x)|u(x))v(x) dx \right].$$

Summing over  $j$  gives

$$\int_{\{(x,s) \in \mathbb{R}_+^{n+1} : \mathcal{M}_\varphi f(x,s) > \lambda\}} \Phi_2(\lambda w(x, s)) d\mu(x, s) \leq \sum_j \Phi_2 \circ \Phi_1^{-1} \left[ \int_{Q_j} \Phi_1(23^n 4^n C \rho |f(x)|u(x))v(x) dx \right].$$

By using that  $\Phi_2 \circ \Phi_1^{-1}$  is convex, this last sum is bounded by

$$\begin{aligned} & \Phi_2 \circ \Phi_1^{-1} \left[ \sum_j \int_{Q_j} \Phi_1(23^n 4^n C \rho |f(x)|u(x))v(x) dx \right] \\ & \leq \Phi_2 \circ \Phi_1^{-1} \left[ C'_n \int_{\mathbb{R}^n} \Phi_1(23^n 4^n C \rho |f(x)|u(x))v(x) dx \right] \\ & \leq \Phi_2 \circ \Phi_1^{-1} \left[ \int_{\mathbb{R}^n} \Phi_1(C_n |f(x)|u(x))v(x) dx \right]. \quad \square \end{aligned}$$

COROLLARY 3.2. From (3), put  $w = u = 1, \Phi_1 = \Phi_2$ . Then (3) gives

$$(13) \quad \left[ \mu(\tilde{Q}) \frac{\varphi(|Q|)}{|Q|} \right] \phi \left[ \frac{\varphi(|Q|)}{C|Q|} \int_Q \psi\left(\frac{\varepsilon}{v(y)}\right) dy \right] \leq C\varepsilon, \text{ for each } \varepsilon > 0.$$

*Proof.* We follow the proof of Corollary 3.2 of [8]. □

COROLLARY 3.3. If  $\Phi_1 \circ \Phi_2^{-1}$  has the  $\Delta'$  condition and (3) holds for  $w = u = 1$ , then there exist constants  $C', C'' > 0$  such that, for any  $\varepsilon > 0$ ,

$$(14) \quad \Phi_1 \circ \Phi_2^{-1} \left( \mu(\tilde{Q}) \frac{\varphi(|Q|)}{|Q|} \right) \phi_1 \left[ \frac{C'}{|Q|} \int_Q \psi_1\left(\frac{\varepsilon}{v(y)}\right) dy \right] \leq C''\varepsilon.$$

*Proof.* We follow the proof of corollary 3.3 of [8].  $\square$

EXAMPLE 3.4. Let  $\varphi(|Q|) = 1$  and  $\varepsilon = \frac{1}{t}$ . From Corollary 3.2, (13) gives  $A_{\Phi}^+$  of [4], i.e.

$$\sup_{Q, t > 0} \phi \left[ \frac{1}{|Q|} \int_Q \psi \left( \frac{1}{tv(x)} \right) dx \right] \frac{t\mu(\tilde{Q})}{|Q|} < \infty.$$

But (13) is weaker than  $A_{\Phi}^+$ , because an  $N$ -function  $\Phi$  of  $A_{\Phi}^+$  in [4] satisfies the  $\Delta_2$ -condition.

EXAMPLE 3.5. From Corollary 3.3, put  $d\mu(x, t) = u(x)dx \otimes d\delta(t)$ , where  $\delta$  is the Dirac mass on  $[0, \infty)$ , concentrated at 0. Also set  $\Phi_1(x) = \frac{x^p}{p}$  and  $\Phi_2(x) = \frac{x^q}{q}$ , where  $1 < p \leq q < \infty$ . Then (14) gives  $(u, v) \in A(\varphi, p, q)$  of [12], i.e.

$$\varphi(|Q|)|Q|^{\frac{1}{q}-\frac{1}{p}} \left[ \frac{1}{|Q|} \int_Q u(x) dx \right]^{\frac{1}{q}} \left[ \frac{1}{|Q|} \int_Q v(x)^{-\frac{1}{p-1}} dx \right]^{1-\frac{1}{p}} \leq A \quad \text{for all } Q.$$

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