



## Chemical Shift and Quadrupolar Interactions in Solids

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**Abstract :** General expressions for solid state NMR lines are described for transitions under static, magic angle spinning, and variable angle spinning conditions in the case where the principal axis system for the anisotropic chemical shift tensor is noncoincident with that of the quadrupole coupling tensor. It is demonstrated that solid state NMR powder pattern simulation program VMAS based on the conventional grid point method of integrating over the Euler angle space is fast enough in comparison with the POWDER simulation package and Gauss-point method.

Keywords: single-crystal, solid state NMR, chemical shift, quadrupole, magic angle, variable angle

## INTRODUCTION

In solid state NMR, information about the electronic environment and bonding of the nuclei can be obtained from the analysis of powder line shapes. The observed spectrum depends on the interactions present in the sample that add to the Zeeman splitting caused by the static magnetic field. Such interactions are dipolar, quadrupolar (Q), chemical shift (CS) and scalar couplings. Small interactions such as scalar couplings are usually negligible compared to the CS and/or Q interactions. When a suitably sized single crystal is available, it is fairly straightforward process to determine  $Q^{1-3}$  or  $CS^{4,5}$  tensor. There is no reason to believe that the principal axis system (PAS) for the anisotropic CS tensor will be coincident with that of the Q tensor. The case where the principal axis system (PAS) for the anisotropic

CS tensor is noncoincident with that of the Q tensor was described by Power et al. more than a decade ago.<sup>6</sup> Nevertheless, one of the Q principle axes is still coincident with one of the CS principle axes.

On the other hand, NMR peak in polycrystalline solids is remarkably broader than narrower lines in solution. This line broadening is primarily due to an orientation dependence of the magnetic field at the nucleus generated by the anisotropic interactions. In order to obtain these anisotropic parameters, the lineshape analysis of static, magic angle spinning (MAS)<sup>7</sup> or variable angle spinning (VAS)<sup>8</sup> NMR spectra at different magnetic fields are required. Analysis of NMR powder patterns arising from quadrupolar nuclei by spectral simulation is complicated because, in general, the lineshape will depend on both the magnitude and orientation of the Q and CS interactions. The dipolar interaction and the lifetime broadening will broaden and smooth out the theoretical powder pattern. Numerically, the line broadenings can be treated by convolution with a normalized Gaussian or Lorentzian function. Analysis of NMR powder patterns arising from quadrupolar nuclei by spectral simulation is complicated because, in general, the line shape will depend on both the magnitudes and relative orientation of the Q and CS interactions.

Considerable literature is available on the nature of the NMR spectra of powdered samples in the presence of quadrupolar interactions only<sup>9</sup> or chemical shift effects only.<sup>10</sup> There have been earlier studies on the effects of both Q and CS interactions on the NMR line shapes of powder sample, where the interactions were axially symmetric and had the same orientation dependence; i.e., the principal components of the Q and CS tensors could be defined in the same system.<sup>11</sup> Such situations occur for nuclei at trigonal or higher symmetry point group sites in the crystal, as these sites possess an  $n$ -fold axis of rotation where  $n > 2$ . Baugher et al. described the NMR powder pattern for systems involving both nonaxially symmetric Q and CS interactions.<sup>12</sup> For a more general case, one has to consider the second rank tensor properties of both the Q and CS tensors. There is no reason to believe that the principal axis system (PAS) for the anisotropic CS tensor will be coincident with that of the Q tensor. This case was independently described by Cheng et al. and Power et al. at the same time.<sup>13,14</sup> In spite of this achievement, no attempts have been reported accounting for the completely general case where the magnitudes and orientations of the principal axes of the two interactions in a single crystal are allowed to be completely indep-

endent. This paper reports general expressions for NMR lines for transitions under static, MAS, and VAS conditions in the case where the principal axis system (PAS) for the anisotropic CS tensor is noncoincident with that of the Q tensor.

## THEORY

The first step in the analysis of single-crystal NMR data is to place the tensors in a common reference frame in a right-handed fashion. It is convenient to express the CS tensor in the frame of the Q tensor since the more dominant Q interaction is considered up to second order. The Euler angles ( $\varphi$ ,  $\theta$ ,  $\psi$ ) relate the magnetic field  $\mathbf{B}_0$  (X, Y, Z) to the goniometer frame ( $x_g$ ,  $y_g$ ,  $z_g$ ). The Euler angles ( $\alpha$ ,  $\beta$ ,  $\gamma$ ) relate the goniometer frame to the PAS of the Q tensor of the nucleus (x, y, z). The orientation of the CS tensor in the Q frame of reference is defined by the three Euler angles ( $\alpha'$ ,  $\beta'$ ,  $\gamma'$ ). The sign conventions and the Eulerian angle definitions of Rose are adopted.<sup>15</sup> To express the resonance frequency in the laboratory frame, we have utilized the standard method of employing Wigner rotation matrices. Up to the second order in the shift, the single crystal NMR frequency  $\nu_m$  that corresponds to the spin transition  $m \rightarrow m - 1$  is given by

$$\begin{aligned} \nu_m = & \nu_L(1 + \delta_{\text{iso}} + \delta f_{20\text{CS}}) - \nu_Q(m - 1/2)f_{20Q} \\ & + (\nu_Q^2/12\nu_L)[\{24m(m-1) - 4I(I+1) + 9\}|f_{2\pm 1Q}|^2 \\ & - \{6m(m-1) - 2I(I+1) + 3\}|f_{2\pm 2Q}|^2], \end{aligned} \quad (1)$$

where  $I$  and  $m$  are the nuclear spin and the magnetic quantum number, respectively, and

$$\begin{aligned} f_{20\text{CS}} = & 0.5(3\cos^2\theta - 1)f_{20\text{CS}'} + (3/2)^{1/2}\sin\theta\cos\theta(e^{i\varphi}f_{2-1\text{CS}'} - e^{-i\varphi}f_{21\text{CS}'}) \\ & + (3/8)^{1/2}\sin^2\theta(e^{i2\varphi}f_{2-2\text{CS}'} + e^{-i2\varphi}f_{22\text{CS}'}), \end{aligned} \quad (2)$$

$$\begin{aligned} f_{20\text{CS}'} = & 0.5(3\cos^2\beta - 1)f_{20\text{CS}''} + (3/2)^{1/2}\sin\beta\cos\beta(e^{i\alpha}f_{2-1\text{CS}''} - e^{-i\alpha}f_{21\text{CS}''}) \\ & + (3/8)^{1/2}\sin^2\beta(e^{i2\alpha}f_{2-2\text{CS}''} + e^{-i2\alpha}f_{22\text{CS}''}), \end{aligned} \quad (3)$$

$$\begin{aligned}
f_{21CS}' &= (\frac{3}{2})^{1/2} \sin \beta \cos \beta e^{-i\gamma} f_{20CS}'' + 0.5 (2 \cos \beta + 1) (1 - \cos \beta) e^{-i\gamma} e^{ia} f_{2-1CS}'' \\
&\quad + 0.5 (2 \cos \beta - 1) (\cos \beta + 1) e^{-i\gamma} e^{-ia} f_{21CS}'' - 0.5 \sin \beta (\cos \beta - 1) e^{-i\gamma} e^{i2a} f_{2-2CS}'' \\
&\quad - 0.5 \sin \beta (1 + \cos \beta) e^{-i\gamma} e^{-i2a} f_{22CS}'',
\end{aligned} \tag{4}$$

$$\begin{aligned}
f_{2-1CS}' &= -(\frac{3}{2})^{1/2} \sin \beta \cos \beta e^{i\gamma} f_{20CS}'' + 0.5 (2 \cos \beta - 1) (\cos \beta + 1) e^{i\gamma} e^{ia} f_{2-1CS}'' \\
&\quad + 0.5 (2 \cos \beta + 1) (1 - \cos \beta) e^{i\gamma} e^{-ia} f_{21CS}'' + 0.5 \sin \beta (1 + \cos \beta) e^{i\gamma} e^{i2a} f_{2-2CS}'' \\
&\quad + 0.5 \sin \beta (\cos \beta - 1) e^{i\gamma} e^{-i2a} f_{22CS}'',
\end{aligned} \tag{5}$$

$$\begin{aligned}
f_{22CS}' &= (\frac{3}{8})^{1/2} \sin^2 \beta e^{-i2\gamma} f_{20CS}'' - 0.5 \sin \beta (\cos \beta - 1) e^{-i2\gamma} e^{ia} f_{2-1CS}'' \\
&\quad + 0.5 \sin \beta (1 + \cos \beta) e^{-i2\gamma} e^{-ia} f_{21CS}'' + \sin^4(\beta/2) e^{-i2\gamma} e^{i2a} f_{2-2CS}'' \\
&\quad + \cos^4(\beta/2) e^{-i2\gamma} e^{-i2a} f_{22CS}'',
\end{aligned} \tag{6}$$

$$\begin{aligned}
f_{2-2CS}' &= (\frac{3}{8})^{1/2} \sin^2 \beta e^{i2\gamma} f_{20CS}'' - 0.5 \sin \beta (1 + \cos \beta) e^{i2\gamma} e^{ia} f_{2-1CS}'' \\
&\quad + 0.5 \sin \beta (\cos \beta - 1) e^{i2\gamma} e^{-ia} f_{21CS}'' + \cos^4(\beta/2) e^{i2\gamma} e^{i2a} f_{2-2CS}'' \\
&\quad + \sin^4(\beta/2) e^{i2\gamma} e^{-i2a} f_{22CS}'',
\end{aligned} \tag{7}$$

$$f_{20CS}'' = 0.5 (3 \cos^2 \beta' - 1 + \eta^{CS} \sin^2 \beta' \cos 2\alpha'), \tag{8}$$

$$f_{2\pm 1CS}'' = \pm \{ (\frac{3}{2})^{1/2} \sin \beta' \cos \beta' - (\eta^{CS}/6^{1/2}) (\sin \beta' \cos \beta' \cos 2\alpha' \mp i \sin \beta' \sin 2\alpha') \} e^{\mp i\gamma}, \tag{9}$$

$$f_{2\pm 2CS}'' = 0.5 \{ (\frac{3}{2})^{1/2} \sin^2 \beta' + (\eta^{CS}/6^{1/2}) [(1 + \cos^2 \beta') \cos 2\alpha' \mp i \cos \beta' \sin 2\alpha'] \} e^{\mp i2\gamma}, \tag{10}$$

$$\begin{aligned}
f_{20Q} &= 0.5 (3 \cos^2 \theta - 1) f_{20Q}' + (\frac{3}{2})^{1/2} \sin \theta \cos \theta (e^{i\varphi} f_{2-1Q}' - e^{-i\varphi} f_{21Q}') \\
&\quad + (\frac{3}{8})^{1/2} \sin^2 \theta (e^{i2\varphi} f_{2-2Q}' + e^{-i2\varphi} f_{22Q}'),
\end{aligned} \tag{11}$$

$$\begin{aligned}
f_{21Q} &= (\frac{3}{2})^{1/2} \sin \theta \cos \theta e^{-i\psi} f_{20Q}' + 0.5 (2 \cos \theta + 1) (1 - \cos \theta) e^{-i\psi} e^{i\varphi} f_{2-1Q}' \\
&\quad + 0.5 (2 \cos \theta - 1) (\cos \theta + 1) e^{-i\psi} e^{-i\varphi} f_{21Q}' - 0.5 \sin \theta (\cos \theta - 1) e^{-i\psi} e^{i2\varphi} f_{2-2Q}' \\
&\quad - 0.5 \sin \theta (1 + \cos \theta) e^{-i\psi} e^{-i2\varphi} f_{22Q}',
\end{aligned} \tag{12}$$

$$\begin{aligned}
f_{2-1Q} &= -(\frac{3}{2})^{1/2} \sin \theta \cos \theta e^{i\psi} f_{20Q}' + 0.5 (2 \cos \theta - 1) (\cos \theta + 1) e^{i\psi} e^{i\varphi} f_{2-1Q}' \\
&\quad + 0.5 (2 \cos \theta + 1) (1 - \cos \theta) e^{i\psi} e^{-i\varphi} f_{21Q}' + 0.5 \sin \theta (1 + \cos \theta) e^{i\psi} e^{i2\varphi} f_{2-2Q}'
\end{aligned}$$

$$+ 0.5 \sin \theta (\cos \theta - 1) e^{i\psi} e^{-i2\varphi} f_{22Q}', \\ (13)$$

$$f_{22Q} = \left(\frac{3}{8}\right)^{1/2} \sin^2 \theta e^{-i2\psi} f_{20Q}' - 0.5 \sin \theta (\cos \theta - 1) e^{-i2\psi} e^{i\varphi} f_{2-1Q}' \\ + 0.5 \sin \theta (1 + \cos \theta) e^{-i2\psi} e^{-i\varphi} f_{21Q}' + \sin^4(\theta/2) e^{-i2\psi} e^{i2\varphi} f_{2-2Q}' \\ + \cos^4(\theta/2) e^{-i2\psi} e^{-i2\varphi} f_{22Q}', \\ (14)$$

$$f_{2-2Q} = \left(\frac{3}{8}\right)^{1/2} \sin^2 \theta e^{i2\psi} f_{20Q}' - 0.5 \sin \theta (1 + \cos \theta) e^{i2\psi} e^{i\varphi} f_{2-1Q}' \\ + 0.5 \sin \theta (\cos \theta - 1) e^{i2\psi} e^{-i\varphi} f_{21Q}' + \cos^4(\theta/2) e^{i2\psi} e^{i2\varphi} f_{2-2Q}' \\ + \sin^4(\theta/2) e^{i2\psi} e^{-i2\varphi} f_{22Q}', \\ (15)$$

$$f_{20Q}' = 0.5 (3 \cos^2 \beta - 1 + \eta^Q \sin^2 \beta \cos 2\alpha) \\ = A, \\ (16)$$

$$f_{2\pm 1Q}' = \pm \left\{ \left(\frac{3}{2}\right)^{1/2} \sin \beta \cos \beta - (\eta^Q / 6^{1/2}) (\sin \beta \cos \beta \cos 2\alpha \mp i \sin \beta \sin 2\alpha) \right\} e^{\mp i\gamma} \\ = \{ \mp B + iC \} e^{\mp i\gamma}, \\ (17)$$

$$f_{2\pm 2Q}' = 0.5 \left\{ \left(\frac{3}{2}\right)^{1/2} \sin^2 \beta + (\eta^Q / 6^{1/2}) [(1 + \cos^2 \beta) \cos 2\alpha \mp i \cos \beta \sin 2\alpha] \right\} e^{\mp i2\gamma} \\ = \{ D \pm iE \} e^{\mp i2\gamma}. \\ (18)$$

The terms  $\nu_L$  and  $\nu_Q$  are the Larmor and the quadrupolar frequencies, respectively. The chemical shift tensor elements,  $\delta_{11}$ ,  $\delta_{22}$ , and  $\delta_{33}$ , are chosen such that

$$|\delta_{33} - \delta_{\text{iso}}| \geq |\delta_{22} - \delta_{\text{iso}}| \geq |\delta_{11} - \delta_{\text{iso}}| \\ (19)$$

$$\delta_{\text{iso}} = (\delta_{33} + \delta_{22} + \delta_{11})/3 \\ (20)$$

$$\delta = \delta_{33} - \delta_{\text{iso}} \\ (21)$$

$$\eta^{\text{CS}} = (\delta_{11} - \delta_{22}) / (\delta_{33} - \delta_{\text{iso}}). \\ (22)$$

The electric field gradient tensor elements,  $V_{11}$ ,  $V_{22}$ , and  $V_{33}$ , are chosen in a similar way to those of the chemical shift tensor.

$$|V_{33}| \geq |V_{22}| \geq |V_{11}| \quad (23)$$

$$\eta^Q = (V_{11} - V_{22}) / V_{33} \quad (24)$$

When a single crystal is mounted on a goniometer in a well defined fashion, i.e. the crystal axes have a known relation to the goniometer axes. The goniometer rotates clockwise the crystal by an angle  $\varphi$  about its rotation axis  $z_g$  from the starting position ( $X // y_g$ ,  $X // z_g$ ,  $X // x_g$ ) which usually has an angle  $\theta = \pi/2$  with respect to the magnetic field. The Euler angles ( $\alpha, \beta, \gamma$ ) are calculated from the matrix, which is composed of the eigenvectors, using  $R(\alpha \beta \gamma)$ .

$$R(\alpha \beta \gamma) = \begin{pmatrix} \cos\alpha \cos\beta \cos\gamma - \sin\alpha \sin\gamma & \sin\alpha \cos\beta \cos\gamma + \cos\alpha & \sin\gamma - \sin\beta \cos\gamma \\ -\cos\alpha \cos\beta \sin\gamma - \sin\alpha \cos\gamma & -\sin\alpha \cos\beta \sin\gamma + \cos\alpha & \cos\gamma \sin\beta \sin\gamma \\ \cos\alpha & \sin\beta \sin\alpha \sin\beta & \cos\beta \end{pmatrix} \quad (25)$$

Then, the transformation from the goniometer frame to the laboratory frame is described by  $R(\varphi 90^\circ 90^\circ)$ . The crystal is rotated in small steps from  $0^\circ$  to  $180^\circ$  about mutually orthogonal directions. From fitting of the rotation plots, eigenvalues such as  $\delta_{11}, \delta_{22}, \delta_{33}, V_{11}, V_{22}$ , and  $V_{33}$ , and eigenvectors are obtained. In the case of off-axis rotation experiment ( $\theta \neq 90^\circ$ ), one rotation from  $0^\circ$  to  $360^\circ$  is necessary for the determination of the tensors.

In the case of static experiment, the Euler angles ( $\varphi, \theta, \psi$ ) are zero. Therefore,  $f_{20Q} = f_{20Q'}, f_{2\pm 1Q} = f_{2\pm 1Q'}, f_{2\pm 2Q} = f_{2\pm 2Q'}, f_{20CS} = f_{20CS'}, f_{2\pm 1CS} = f_{2\pm 1CS'}, f_{2\pm 2CS} = f_{2\pm 2CS'}$ . Then, the frequency  $v_m$  becomes as follows;

$$\begin{aligned} v_m = & v_L [1 + \delta_{iso} + \delta \{0.5 (3 \cos^2 \beta - 1) f_{20CS}'' + (\frac{3}{2})^{1/2} \sin \beta \cos \beta (e^{ia} f_{2-1CS}'' - e^{-ia} f_{21CS}'') \\ & + (\frac{3}{8})^{1/2} \sin^2 \beta (e^{i2a} f_{2-2CS}'' + e^{-i2a} f_{22CS}'')\}] - v_Q (m - \frac{1}{2}) 0.5 (3 \cos^2 \beta - 1 + \eta^Q \sin^2 \beta \cos 2\alpha) \\ & + (v_Q^2 / 12 v_L) [\{24 m(m-1) - 4 I(I+1) + 9\} \{(-\eta^{Q2}/6 \cos^2 2\alpha + \eta^Q \cos 2\alpha - 3/2) \cos^4 \beta \\ & + (\eta^{Q2}/3 \cos^2 2\alpha - \eta^Q \cos 2\alpha - \eta^{Q2}/6 + 3/2) \cos^2 \beta + \eta^{Q2}/6 (1 - \cos^2 2\alpha)\} \\ & - \{6 m(m-1) - 2 I(I+1) + 3\} \{(\eta^{Q2}/24 \cos^2 2\alpha - 1/4 \eta^Q \cos 2\alpha + 3/8) \cos^4 \beta \\ & + (-\eta^{Q2}/12 \cos^2 2\alpha + \eta^{Q2}/6 - 3/4) \cos^2 \beta + (\eta^{Q2}/24 \cos^2 2\alpha + 1/4 \eta^Q \cos 2\alpha + 3/8)\}]. \end{aligned}$$

(26)

Eq (26) reduced to equations in the literature by Cheng et al. for  $m=1/2$ .<sup>13</sup>

In the case where the principal axes of the CS and Q tensors be coincident ( $\alpha' = \beta' = \gamma' = 0^\circ$ ), the frequency  $v_m$  is simplified to

$$\begin{aligned}
 v_m = & v_L [1 + \delta_{\text{iso}} + \delta \{0.5 (3 \cos^2 \beta - 1 \eta^Q \sin^2 \beta \cos 2\alpha)\}] \\
 & - v_Q (m - 1/2) \{0.5 (3 \cos^2 \beta - 1 + \eta^Q \sin^2 \beta \cos 2\alpha)\} \\
 & + (v_Q^2 / 12 v_L) [\{24 m(m-1) - 4 I(I+1) + 9\} \{(-\eta^{Q2}/6 \cos^2 2\alpha + \eta^Q \cos 2\alpha - 3/2) \cos^4 \beta \\
 & + (\eta^{Q2}/3 \cos^2 2\alpha - \eta^Q \cos 2\alpha - \eta^{Q2}/6 + 3/2) \cos^2 \beta + \eta^{Q2}/6 (1 - \cos^2 2\alpha)\} \\
 & - \{6 m(m-1) - 2 I(I+1) + 3\} \{(\eta^{Q2}/24 \cos^2 2\alpha - \eta^Q/4 \cos 2\alpha + 3/8) \cos^4 \beta \\
 & + (-\eta^{Q2}/12 \cos^2 2\alpha + \eta^{Q2}/6 - 3/4) \cos^2 \beta + (\eta^{Q2}/24 \cos^2 2\alpha + \eta^Q/4 \cos 2\alpha + 3/8)\}]. 
 \end{aligned} \tag{27}$$

Then, the explicit dependence of the central line  $v_{1/2}^{(2)}$  can be obtained as

$$\begin{aligned}
 v_{1/2}^{(2)} = & -(v_Q^2 / 6 v_L) \{I(I+1) - 3/4\} \{(-3/8 \eta^{Q2} \cos^2 2\alpha + 9/4 \eta^Q \cos 2\alpha - 27/8) \cos^4 \beta \\
 & + (3/4 \eta^{Q2} \cos^2 2\alpha - 2\eta^Q \cos 2\alpha - \eta^{Q2}/2 + 15/4) \cos^2 \beta \\
 & + (-3/8 \eta^{Q2} \cos^2 2\alpha - \eta^Q/4 \cos 2\alpha + \eta^{Q2}/3 - 3/8)\}. 
 \end{aligned} \tag{28}$$

Eq (28) is the same as equations in the literatures.<sup>12,16</sup>

The angle  $\varphi$  may be time dependent ( $\varphi = \omega_r t$ ) in a sample spinning experiment. The time dependence is introduced by rotating the sample by an angle  $\varphi(t)$  about the spinner axis with  $\varphi(t) = \varphi_0 + \omega_r t$ . The angle  $\theta$  is the angle between the static field  $\mathbf{B}_0$  and the spinning axis. For rapid sample spinning, such that  $\omega_r$  is greater than the second-order quadrupole linebreadth of the  $(1/2, -1/2)$  transition, the sidebands are well separated from central line and become vanishingly small with increasing  $\omega_r$ . Time dependent parts,  $e^{\pm i\varphi}$  and  $e^{\pm i2\varphi}$  may be averaged over one cycle to give zero (Strictly speaking, they show cyclic behavior which lead to rotational spin echos.<sup>17</sup>), while other time dependent parts,  $(e^{\pm i\varphi})^2$  and  $(e^{\pm i2\varphi})^2$  become one. In the VAS case where the angle  $\theta$  can vary from  $0^\circ$  to  $90^\circ$  with respect to the magnetic field, the  $v_m^{(2)}$  can be written as

$$\begin{aligned}
v_m = & v_L [1 + \delta_{iso} + \delta \{0.5 (3 \cos^2 \theta - 1)\} \{0.5 (3 \cos^2 \beta - 1) f_{20CS}'' \\
& + (\frac{3}{2})^{1/2} \sin \beta \cos \beta (e^{ia} f_{2-1CS}'' - e^{-ia} f_{21CS}'') + (\frac{3}{8})^{1/2} \sin^2 \beta (e^{i2a} f_{2-2CS}'' + e^{-i2a} f_{22CS}'')\}] \\
& - v_Q (m - 1/2) \{0.5 (3 \cos^2 \theta - 1)\} \{0.5 (3 \cos^2 \beta - 1 + \eta^Q \sin^2 \beta \cos 2\alpha)\} \\
& + (v_Q^2 / 12 v_L) [\{24 m(m-1) - 4 I(I+1) + 9\} [1/2 (1 - \cos^4 \theta) (D^2 + E^2) \\
& + 1/2 (1 - 3 \cos^2 \theta + 4 \cos^4 \theta) (B^2 + C^2) + 3/2 (1 - \cos^2 \theta) \cos^2 \theta A^2] \\
& - \{6 m(m-1) - 2 I(I+1) + 3\} [1/8 (1 + 6 \cos^2 \theta + \cos^4 \theta) (D^2 + E^2) \\
& + 1/2 (1 - 4 \cos^4 \theta) (B^2 + C^2) + 3/8 (1 - \cos^2 \theta)^2 A^2], 
\end{aligned} \tag{29}$$

where

$$\begin{aligned}
A &= 0.5 (3 \cos^2 \beta - 1 + \eta^Q \sin^2 \beta \cos 2\alpha) \\
B &= -(\frac{3}{2})^{1/2} \sin \beta \cos \beta + (\eta^Q / 6^{1/2}) \sin \beta \cos \beta \cos 2\alpha \\
C &= \sin \beta \sin 2\alpha \\
D &= 0.5 \{(\frac{3}{2})^{1/2} \sin^2 \beta + (\eta^Q / 6^{1/2}) (1 + \cos^2 \beta) \cos 2\alpha\} \\
E &= -0.5 \cos \beta \sin 2\alpha.
\end{aligned}$$

Under MAS condition, the angle  $\theta$  is  $54.74^\circ$ , i.e.  $3 \cos^2 \theta - 1 = 0$ . The central line  $v_{1/2}$  can be expressed as

$$\begin{aligned}
v_{1/2} = & v_L [1 + \delta_{iso}] - (v_Q^2 / 6 v_L) \{I(I+1) - 3/4\} \{(7/48 \eta^{Q2} \cos^2 2\alpha - 7/8 \eta^Q \cos 2\alpha \\
& + 21/16) \cos^4 \beta + (-7/24 \eta^{Q2} \cos^2 2\alpha + \eta^Q \cos 2\alpha + \eta^{Q2}/12 - 9/8) \cos^2 \beta \\
& + (7/48 \eta^{Q2} \cos^2 2\alpha - \eta^Q/8 \cos 2\alpha + 5/16)\}.
\end{aligned} \tag{30}$$

Eq (30) is the same as Eq (6) by Müller.<sup>18</sup>

A Fortran program ‘XTALNMR’ is listed at the end of this article, which is based on Eq (1) in order to calculate single-crystal nmr lines at every given rotations with a starting orientation with the Euler angles ( $\alpha, \beta, \gamma$ ) and ( $\alpha', \beta', \gamma'$ ). This program with nonlinear least squares fitting program Kinfit<sup>19</sup> was used to fit single-crystal  $^{133}\text{Cs}$  NMR data of  $\text{Cs}^+(15\text{-crown-5})_2\text{I}$ .<sup>20</sup>

A powder spectrum is computed by summing the spectra from an ensemble of randomly oriented single crystals. A Fortran program ‘VMAS’ is given at the end of this article, which is based on Eq (26) in order to calculate a powder NMR line under static,

magic angle, or variable angle condition. The efficiency of this program is comparable to the program ‘POWDER’<sup>21</sup> (under 10 s on a VAX 8800 computer). This comes from the efficiency in the loop calculation for the polar ( $\cos \beta$ ) and azimuthal angle ( $\alpha$ ) grid point space. In the VMAS, common factors in the equation are factorized out, calculated outside of the loop and stored in memory sites. Therefore, calculation in the loop summation is minimized. Run time for  $360 \times 360$  grid point integration was less than 1 s with Pentium III MMX 500MHz computer. Gauss-point integration method was also comparable to the VMAS in run time.<sup>22</sup> The program VMAS with nonlinear least squares fitting program Kinfit2 was used to fit  $^{85}\text{Rb}$ ,  $^{87}\text{Rb}$  and  $^{39}\text{K}$  NMR powder patterns.<sup>19</sup>

In summary, general expressions for solid state NMR lines are described for transitions under static, magic angle spinning, and variable angle spinning conditions in the case ( $\alpha' \neq 0^\circ$ ,  $\beta' \neq 0^\circ$ ,  $\gamma' \neq 0^\circ$ ) where the principal axis system for the anisotropic chemical shift tensor is noncoincident with that of the quadrupole coupling tensor.

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**PROGRAM XTALNMR**

C

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C

C    THE PROGRAM "XTALNMR.FOR" COMPUTES CHEMICAL SHIFTS CORRESPONDING

C    TO NUCLEAR MAGNETIC RESONANCE TRANSITIONS WHICH RESULT FROM

C    THE SPIN HAMILTONIAN WITH NUCLEAR ZEEMAN, QUADRUPOLE AND CHEM-

C    ICAL SHIFT TERMS. THE QUADRUPOLE AND CHEMICAL SHIFT TERMS ARE

C    TREATED AS PERTURBATIONS ON THE ZEEMAN TERM, AND ARE CORRECT

C    THROUGH SECOND ORDER AND FIRST ORDER, RESPECTIVELY.

C    THE SINGLE CRYSTAL NMR PEAK POSITIONS ARE CALCULATED BY SETTING

C    THE EULERIAN ANGLES, ALPHA, BETA, AND GAMMA, WHICH ARE DEFINED

C    WITH RESPECT TO THE GONIOMETER AXIS SYSTEM. IF BOTH TENSORS ARE

C    AXIALLY SYMMETRIC AND PRINCIPAL AXES ARE COINCIDENT, THE

C    EULERIAN ANGLE ALPHA IS MEANINGLESS. IT CAN BE AN ARBITRARY NUMBER.

C    FOR NONCOINCIDENT SYSTEMS, THE TRANSFORMATION OF THE CSA TENSOR TO

C    THE EFG TENSOR PRINCIPAL AXIS SYSTEM IS PERFORMED WITH THE EULERIAN

C    ANGLES, ALPHA', BETA', AND GAMMA'. THE ANGLE VAN IS DEFINED AS AN

C    ANGLE BETWEEN MAGNETIC FIELD DIRECTION AND THE ROTATION

C    (GONIOMETER Z AXIS) AXIS. THE TRANSFORMATIONS ARE DONE BY ROTATION

C    MATRICES SUCH AS R(A,B,C)=R(C)z'R(B)y'R(A)z DESCRIBED IN

C "ELEMENTARY THEORY OF ANGULAR MOMENTUM" BY M. E. ROSE.

C EACH TRANSITION INTENSITY IS PROPORTIONAL TO  $I(I+1)-M(M-1)$ .

C TOTAL TRANSITIONS TO BE CALCULATED ARE  $2I^*$ TOTAL ANGLE OF

C ROTATION(ANG)/ANGLE OF ROTATION(XI).

C EACH TRANSITION INTENSITY IS PROPORTIONAL TO  $I(I+1)-M(M-1)$ .

C TOTAL TRANSITIONS TO BE CALCULATED ARE  $2I^*$ TOTAL ANGLE OF

C ROTATION(ANG)/ANGLE OF ROTATION(XI).

C

C THE FOLLOWING CONVENTION IS CHOSEN TO LABEL THE PRINCIPAL AXES:

C     ABS(D33-D) > or = ABS(D22-D) > or = ABS(D11-D)

C               WHERE D = (D11+D22+D33)/3 AND

C     ABS(V33) > or = ABS(V22) > or = ABS(V11).

C ASYMMETRY PARAMETERS (ETA=(V11-V22)/V33 AND CETA=(D11-D22)/(D33-D))

C FOR QUADRUPOLAR AND CHEMICAL SHIFT INTERACTIONS, RESPECTIVELY)

C ARE IN THE RANGE OF 0 AND 1.

C

C INPUT

C     TITLE OF RUN	18A4
C     SPIN,TRAN,VL,VAN	FREE FORMAT
C         SPIN              NUCLEAR SPIN I	
C         VL                LARMOR FREQUENCY IN MHZ	
C         TRAN              SELECTS TRANSITIONS TO BE CALCULATED	
C         VAN               POLAR ANGLE BETWEEN FIELD DIRECTION AND GONIOMETER	
C                          Z-AXIS IN DEGREE	
C     QCC,ETA,D11,D22,D33	FREE FORMAT
C         QCC              QUADRUPOLAR COUPLING CONSTANT IN MHZ	
C         ETA              ASYMMETRIC PARAMETER OF ELECTRIC FIELD	
C                          GRADIENT AT NUCLEAR SITE	
C         D11,D22,D33	CHEMICAL SHIFT PARAMETERS IN PPM
C     ALPA,BETA,GAMA,XI,ANG	FREE FORMAT
C                          ALPA,BETA,GAMA:EULERIAN ANGLES IN DEGREE	

C XI ANGLE INCREMENT OF EACH ROTATION IN DEGREE(CONSTANT)  
C ANG MAXIMUM(TOTAL) ANGLE OF ROTATION  
C ACQ,BCQ,GCQ  
C ALPHA',BETA',GAMMA':EULERIAN ANGLES BETWEEN CSA TENSOR PRINCIPAL  
C AXES AND EFG TENSOR PRINCIPAL AXES  
C  
C OUTPUT  
C XL,CHEM,DMAG,FIRS,MULTI,SECO,FNN  
C XL XL\*XI = SUM OF THE ANGLE ROTATED  
C CHEM CHEMICAL SHIFT ANISOTROPY(PPM)  
C DMAG M  
C FIRS FIRST ORDER SATELLITE TRANSITION(PPM)  
C MULTI I(I+1)-M(M-1)  
C SECO SECOND ORDER SHIFT FOR EACH TRANSITION(PPM)  
C PFREQ SUM OF CSA, FIRST ORDER, AND SECOND ORDER  
C QUADRUPOLAR SHIFT(PPM)  
C \*\*\*\*\*  
C  
C EXAMPLES OF DATA FILE  
C Cs1.DAT  
C Cs+ AT GNIOMETER Z-AXIS POLAR ANGLE 90  
C 3.5,3.5,23.61,90  
C .15,.5,60,40,-30  
C 0,90,0,10,180  
C 10,90,0  
C Cs2.DAT  
C Cs+ AT A TILTED GONIOMETER(MAS PROBE) WITH AN ANGLE OF 54.7356  
C 3.5,3.5,23.61,54.7356  
C .15,.5,60,40,-30  
C 0,90,0,10,360  
C 10.90,0

```

C ****
C      IMPLICIT REAL*8 (A-H,O-Z)
C      CHARACTER*20 F20,F21
C      DIMENSION TITLE(18)
C      DIMENSION SGM(400),CGM(400),S2G(400),C2G(400)
C      DIMENSION A(10),B(10),C(10),INT(10),D(10)

C **** OPEN FILES ****
C
7      WRITE(*,20)
      READ(*,40,ERR=7) F20
      OPEN(UNIT=1,FILE=F20,STATUS='OLD')
8      WRITE(*,30)
      READ(*,40,ERR=8) F21
      OPEN(UNIT=2,FILE=F21,STATUS='NEW')

C **** FORMAT ****
C
10     FORMAT(/)
20     FORMAT(' ENTER INPUT FILE NAME')
30     FORMAT(' ENTER OUTPUT FILE NAME')
40     FORMAT(A)
50     FORMAT(9H  ALPHA =,F7.3,1X,8H  BETA =,F7.3,1X,9H  GAMMA =,F7.3,
*    1X,17H VARIABLE ANGLE =,F7.3)
60     FORMAT(8H  SPIN =,F4.1,2X,14HLARMOR FREQ. =,F7.3,
*    2X,21HQUAD. COUPL. CONST. =,F7.3,2X,13HASYMM. PAR. =,F7.4)
70     FORMAT(29H  CHEM. SHIFT PARAMETERS;D11=,F9.3,1X,4HD22=,F9.3,1X
*    4HD33=,F9.3)
80     FORMAT(' ISO. CHEM. SHIFT=',F9.3,1X,'CHEM. SHIFT DELTA=',F9.3,
*    1X,'CHEM. SHIFT ETA=',F5.3)
90     FORMAT(1X,31HDELTA M = 1 TRANSITIONS BETWEEN,F5.1,4H AND,F5.1,

```

```

* 5H ONLY)
100  FORMAT(3X,'NO',5X,'CSA',6X,'M',2X,'SATELLITE',3X,'MULTI',
* 2X,'SECOND',2X,'PEAK PPM')
110  FORMAT(1X,F5.1,1X,F9.3,1X,F4.1,1X,F10.3,1X,I4,1X,F9.3,1X,F10.3)
120  FORMAT(9H ALPHA' =,F7.3,1X,8H BETA' =,F7.3,1X,9H GAMMA' =,F7.3,
* 1X,30H FOR CSA TENSOR TRANSFORMATION)
130  FORMAT(18A4)
140  FORMAT(2X,18A4)
C ****
      READ(1,130)TITLE
      READ(1,*)SPIN,TRAN,VL,VAN
      READ(1,*)QCC,ETA,D11,D22,D33
      READ(1,*)ALPA,BETA,GAMA,XI,ANG
      READ(1,*)ACQ,BCQ,GCQ
C ****
      RAD=3.141592653589793/180.

```

### C TERMS FOR TRANSFORMATION OF NONCOINCIDENT CSA TENSOR

```

TRACQ=2.*RAD*ACQ
STACQ=DSIN(TRACQ)
CTACQ=DCOS(TRACQ)
RBCQ=RAD*BCQ
SBCQ=DSIN(RBCQ)
CBCQ=DCOS(RBCQ)
SQSB=SBCQ*SBCQ
SQCB=CBCQ*CBCQ

```

### C EULERIAN ANGLE ALPHA

```

TAL=2.*RAD*ALPA
STP=DSIN(TAL)
CTP=DCOS(TAL)
RGA=RAD*(GCQ+ALPA)
TRGA=2.*RGA

```

SRGA=DSIN(RGA)  
 CRGA=DCOS(RGA)  
 STRGA=DSIN(TRGA)  
 CTRGA=DCOS(TRGA)

### C EULERIAN ANGLE BETA

BTA=RAD\*BETA  
 SBT=DSIN(BTA)  
 CBT=DCOS(BTA)  
 CSB=CBT\*SBT  
 SSB=SBT\*SBT  
 CCB=CBT\*CBT

### C EULERIAN ANGLE GAMMA

J=ANG/XI+1  
 DO 300 I=1,J  
 YI=I-1.  
 ROT=RAD\*(GAMA+YI\*XI)  
 TROT=2.\*ROT  
 SGM(I)=DSIN(ROT)  
 CGM(I)=DCOS(ROT)  
 S2G(I)=DSIN(TROT)  
 C2G(I)=DCOS(TROT)  
 300 CONTINUE

### C TERMS FOR TRANSFORMATION FROM THE GONIOMETER FRAME TO THE LABORATORY

### C AXIS SYSTEM

THE=RAD\*VAN  
 VSAL=DSIN(THE)  
 VCAL=DCOS(THE)  
 VSVA=VSAL\*VCAL  
 VSSA=VSAL\*VSAL  
 VDCB=1.-2.\*VSSA

```

VCCA=.5*(1.+VCAL*VCAL)
VTCT=.5*(3.*VCAL*VCAL-1.)
SIX=SQRT(6.)

```

### C CSA TERMS

```

CSISO=(D11+D22+D33)/3.
CSDEL=D33-CSISO
IF (CSDEL.EQ.0.) GO TO 350
CSETA=(D11-D22)/CSDEL
GO TO 400
350 CSETA=0.
400 ACS=.5*(3.*SQCB-1.+CSETA*SQSB*CTACQ)
BCS=SBCQ*CBCQ*(3.-CSETA*CTACQ)/SIX
CCS=CSETA*SBCQ*STACQ/SIX
DCS=.5*(3.*SQSB+CSETA*(1.+SQCB)*CTACQ)/SIX
ECS=CSETA*CBCQ*STACQ/SIX
DCNES=DCS*CTRGA-ECS*STRGA
DSEC=DCS*STRGA+ECS*CTRGA
BCCS=BCS*CRGA+CCS*SRGA
XNBSCC=-BCS*SRGA+CCS*CRGA
AO=.5*SIX*SSB*DCNES-SIX*CSB*BCCS+.5*(3.*CCB-1.)*ACS
BR=-CSB*DCNES+(2.*CCB-1.)*BCCS+.5*SIX*CSB*ACS
CI=SBT*DSEC+CBT*XNBSCC
DR=.5*(1.+CCB)*DCNES+CSB*BCCS+.25*SIX*SSB*ACS
EI=CBT*DSEC-SBT*XNBSCC

```

### C QUADRUPOLAR TERMS

```

SS1=SPIN*(SPIN+1.)
IF(SPIN - 0.5) 450,450,500
450 VQQ = 0.0
GO TO 550
500 VQQ=3.*QCC/(2.*SPIN*(2.*SPIN-1.))/VL
550 FACT=VQQ*VQQ/12.

```

```

XMAG=-TRAN+1.
MAG=1
600 CONTINUE
XX1=XMAG*(XMAG-1.)
A(MAG)=24.*XX1-4.*SS1+9.
B(MAG)=6.*XX1-2.*SS1+3.
C(MAG)=-VQQ*(XMAG-.5)
D(MAG)=XMAG
INT(MAG)=SS1-XX1
IF(XMAG.GE.TRAN) GO TO 650
XMAG=XMAG+1.
MAG=MAG+1
GO TO 600
650 CONTINUE
MGMX=MAG
ECP=ETA*CTP
XNECP3=3.-ECP
CON=XNECP3*CSB/SIX
SON=ETA*SBT*STP/SIX
CTW=(3.+ECP-XNECP3*CCB)*.5/SIX
STW=ETA*CBT*STP/SIX
FR=.5*(-1.+ECP+XNECP3*CCB)

```

### C ROTATION ANGLE DEPENDENT TERMS

### C CALCULATE SHIFT CONTRIBUTIONS TO EACH TRANSITION

```

WRITE(2,100)
DO 700 M=1,MGMX
DO 700 L=1,J
XL=XI*(L-1.)
CS2L=CON*SGM(L)-SON*CGM(L)
CS3L=CTW*C2G(L)-STW*S2G(L)
CS4L=CON*CGM(L)+SON*SGM(L)

```

```

CS5L=CTW*C2G(L)+STW*S2G(L)
CS6L=CTW*S2G(L)+STW*C2G(L)
SAT=.5*SIX*VSSA*CS3L-SIX*VSCA*CS4L+VTCT*FR
F1R=-VSCA*CS3L+VDCB*CS4L+.5*SIX*VSCA*FR
F1I=VSAL*CS6L-VCAL*CS2L
F2R=VCCA*CS3L+VSCA*CS4L+.25*SIX*VSSA*FR
F2I=VCAL*CS6L+VSAL*CS2L
CAN=CSDEL*(.5*SIX*VSSA*(DR*C2G(L)-EI*S2G(L))-SIX*VSCA*
*(BR*CGM(L)+CI*SGM(L))+VTCT*AO)
DMAG=D(M)
MULTI=INT(M)
VLCC=CSISO+CAN
FST=C(M)*SAT
SND=FACT*((F1R*F1R+F1I*F1I)*A(M)-(F2R*F2R+F2I*F2I)*B(M))
CHEM=VLCC
FIRS=FST*1.E6
SECO=SND*1.E6
PFREQ=VLCC+(FST+SND)*1.E6
WRITE(2,110) XL,CHEM,DMAG,FIRS,MULTI,SECO,PFREQ
700    CONTINUE
C
C      ***** OUTPUT FOR PARAMETERS *****
C
WRITE(2,10)
WRITE(2,140) TITLE
WRITE(2,10)
WRITE(2,60) SPIN,VL,QCC,ETA
WRITE(2,10)
WRITE(2,70) D11,D22,D33
WRITE(2,10)
WRITE(2,80) CSISO,CSDEL,CSETA

```

```
WRITE(2,10)
WRITE(2,50) ALPA,BETA,GAMA,VAN
WRITE(2,10)
WRITE(2,90) TRAN,-TRAN
WRITE(2,10)
WRITE(2,120) ACQ,BCQ,GCQ
C ****
STOP
END
```

PROGRAM VMAS

```
C
C SOLID STATE NMR POWDER PATTERN(STATIC, MAS, AND VAS) SIMULATION
C PROGRAM FOR NONCOINCIDENT PRINCIPAL AXIS SYSTEMS
C
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C DEPARTMENT OF CHEMISTRY
C GYEONGSANG NATIONAL UNIVERSITY
C CHINJU 660-701 KOREA
C JAMES L. DYE
C DEPARTMENT OF CHEMISTRY
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C
C THIS PROGRAM IS AN EXTENSION OF QCPE 154 AND QCMP 56.
C ***POWPAT PROGRAM***
C H.S. STORY AND D. KLINE DEPARTMENT OF PHYSICS
C STATE UNIVERSITY OF NEW YORK AT ALBANY
C POWPAT PROGRAM COMPUTES AND PLOTS POWDER PATTERNS AND SPECTRA
C CORRESPONDING TO NUCLEAR MAGNETIC RESONANCE TRANSITIONS WHICH
```



C TRAN SELECTS TRANSITIONS TO BE INCLUDED  
 C VL LARMOR FREQUENCY  
 C PS NUMBER OF FREQUENCIES AT WHICH POWDER  
     PATTERN AND SPECTRUM WILL BE COMPUTED.  
 C PPMIN,PPMAX MINIMUM AND MAXIMUM CHEM. SHIFTS IN  
     POWDER PATTERN AND SPECTRUM.  
 C DUMMY NUMBER TO ADD A BACKGROUND SPECTRUM  
 C  
 C QCC NUCLEAR QUADRUPOLE COUPLING CONSTANT  
 C ETA ASYMMETRY PARAMETER OF ELECTRIC FIELD  
     GRADIENT AT NUCLEAR SITE  
 C  
 C D11,D22,D33 CHEMICAL SHIFT PARAMETERS(NOT SHELDING)  
 C  
 C SIGHZ GAUSSIAN BROADENING IN HZ  
 C WEIGHT RELATIVE WEIGHT TO BE ASSIGNED TO SITE  
 C NORM NORMALIZE INTENSITY, IF NORM = 1  
 C  
 C ALP VARIABLE ANGLE  
 C ACQ,BCQ,GCQ EULERIAN ANGLES BETWEEN QUADRUPOLE COUPLING AND  
     CHEMICAL SHIFT TENSORS  
 C ITH NUMBER OF POINTS OF SPHERE DIVISION  
 C  
 C A GRID IS ESTABLISHED IN COSINE THETA, PHI SPACE WITH EACH  
 C COORDINATE HAVING XITH INCREMENTS, SO THAT THE GRID IS XITH\*XITH  
 C IF THE QUADRUPOLE AND CHEMICAL SHIFT TENSORS ARE GENERAL.  
 C IF BOTH ARE AXIAL AND COINCIDENT, THE PHI ROTATION IS NOT  
 C DONE. VL,PMIN, PMAX, AND QCC ARE IN MEGAHERTZ. ETA AND  
 C WEIGHT ARE DIMENSIONLESS. THE PRINCIPAL AXIS SYSTEM OF THE  
 C QUADRUPOLEAR COUPLING TENSOR (Xq,Yq,Zq) IS DEFINED SO THAT  
 C THE ABSOLUTE MAGNITUDE OF Vzz AND Vxx CAN BE THE LARGEST

C AND THE SMALLEST, RESPECTIVELY. ETA IS DEFINED AS: ETA =  
C  $(V_{xx} - V_{yy})/V_{zz}$ . THE ELEMENTS (D11,D22,D33) OF THE CHEMICAL  
C SHIFT TENSOR ARE CHOSEN IN THE SAME PRINCIPAL AXIS SYSTEM AS  
C QUADRUPOLAR COUPLING TENSOR, SPECIFICALLY,  
C  
C \*\*\*\*\*  
C \*\* ABS(D33 - Diso) GE. ABS(D22 - Diso) GE. ABS(D11 - Diso) \*\*  
C \*\*\*\*\*  
C  
C D11, D22, AND D33 ARE IN PPM SCALE. THE TRANSFORMATION OF THE  
C CSA TENSOR TO THE QUADRUPOLAR PRINCIPAL AXIS SYSTEM IS DONE  
C BY THE EULERIAN ANGLES, ACQ, BCQ, AND GCQ.  
C SIG IS THE GAUSSIAN HALFWIDTH EXPRESSED AS A FRACTION OF THE  
C TOTAL RANGE PMAX-PMIN. THE PATTERNS WILL INCLUDE THE INTENSITY  
C FROM ALL ADJACENT LEVEL TRANSITIONS BETWEEN THE TRAN AND-TRAN,  
C I.E. IF TRAN IS SET EQUAL TO .5 FOR HALF ODD-INTEGRAL SPIN, THEN  
C ONLY THE CENTRAL TRANSITION WILL BE DONE. IF TRAN IS SET TO 1.5  
C THEN THE CENTRAL TRANSITION PLUS TWO ADJACENT SATELLITES WILL BE  
C INCLUDED. THE WEIGHTING FACTOR  $I(I+1)-M(M-1)$  FOR THE SATELLITE  
C INTENSITIES IS INCLUDED. I IS THE SPIN AND M IS THE MAGNETIC  
C QUANTUM NUMBER. DIMENSIONS OF PARAMETERS WHICH AGREE WITH  
C EACH OTHER ARE COUPLED AND SHOULD BE RESET ACCORDINGLY.  
C THIS PROGRAM SIMULATES VARIABLE ANGLE SAMPLE SPINNING NMR SPECTRA  
C (NOT SPINNING SIDEBANDS).IF FAST SPINNING CAN AVERAGE FIRST ORDER  
C SATELLITES TO ZERO, THE SIMULATION PROGRAM WHICH USES AVERAGE  
C HAMILTONIAN THEORY WILL BE APPROPRIATE (M. M. MARICQ AND  
C J. S. WAUGH, J. CHEM. PHYS. 70, 3300(1977).  
C STATIC SPECTRA CAN BE SIMULATED BY SETTING THE VARIABLE ANGLE  
C ALP = 0.  
C \* DUMMY = FACTOR TO DETERMINE IF SIMULATION IS TO BE NORMAL(ANY  
C NUMBER WHICH IS NOT EQUAL TO 2) OR IF SIMULATION IS TO ADD IN A

C BACKGROUND SPECTRUM FROM A FILE(SPECTRUM.DAT). THE VL,PPMIN,PPMAX  
 C AND PS MUST BE THE SAME TO ADD. THE RELATIVE RATIOS BETWEEN THE  
 C SITES SHOULD BE USED. NORM SHOULD BE 0.  
 C AN ARRAY P(K) STORES FREQUENCIES.  
 C THE UNIT 2 OUTPUT FILE ' F21 ' KEEPS PARAMETERS.  
 C THE UNIT 1 OUTPUT FILE ' F22 ' HAS CHEMICAL SHIFTS  
 C AND CORRESPONDING FREQUENCIES.  
 C THE UNIT 4 FILE IS A DATA FILE TO EXECUTE THIS PROGRAM(TO ADD A  
 C SPECTRUM).  
 C  
 C VMAS1.DAT VMAS2.DAT  
 C\*TEST1 J.PHYS.CHEM. V94, P591 (1990) \*TEST2 J.PHYS.CHEM. V94, P591 (1990)  
 C\*3.5,3.5,26.247,1000,-1700,1700,0 \*3.5,3.5,26.247,1000,-3500,3500,0  
 C\*0.142,0.11,127,136,-197,200,1,1 \*0.142,0.11,127,136,-197,200,1,0  
 C\*0,0,0,0,360 \*0,0,52,0,360  
 C  
 C VMAS3.DAT VMAS4.DAT  
 C\*TEST3 J.PHYS.CHEM. V94, P591 (1990) \*TEST4 J.PHYS.CHEM. V94, P591 (1990)  
 C\*3.5,0.5,50,1000,-1000,1000,0 \*3.5,3.5,26.247,1000,-3500,3500,2  
 C\*14,0.2,50,150,-200,200,1,1 \*0.373,0.55,-163,-211,59,200,1,0  
 C\*0,0,60,0,360 \*0,90,7,0,360  
 C  
 C VMAS5.DAT VMAS6.DAT  
 C\*TEST5 J.CHEM.PHYS. V77, P4360 (1982)\*TEST6 J.CHEM.PHYS. V77, P4360 (1982)  
 C\*1.5,0.5,95.25,200,-75,25,0 \*1.5,0.5,95.25,200,-85,15,0  
 C\*2.6,0,0,0,0,100,1,1 \*2.6,0.3,0,0,0,100,1,0  
 C\*54.74,0,0,0,360 \*40,0,0,0,360  
 C  
 C 1. THE RESULT OF TEST1 IS A Cs NMR SPECTRUM AT 26.247 MHz WHICH SHOWS 7  
 C TRANSITIONS WITH BOTH QUADRUPOLAR AND CHEMICAL SHIFT ANISOTROPY.  
 C

C 2. THE TEST2 YIELDS A Cs NMR SPECTRUM WITH AN EULERIAN ANGLE BETWEEN  
C QUADRUPOLAR AND CHEMICAL SHIFT TENSORS.  
C  
C 3. THE TEST3 DEMONSTRATES THE EFFECT OF AN EULERIAN ANGLE BETWEEN  
C QUADRUPOLAR AND CHEMICAL SHIFT TENSORS ON A CENTRAL TRANSITION.  
C  
C 4. IF THERE ARE TWO SITES, THE SUM OF THE SPECTRA CAN BE GENERATED BY  
C USING THE VMAS4.DAT AND THE "SPECTRUM.DAT" WHICH IS RENAMED FROM  
C THE OUTPUT OF VMAS2.DAT("WEIGHT" CAN BE VARIED IN ORDER TO CONSIDER  
C THE RELATIVE RATIO BETWEEN TWO SITES.).  
C  
C 5. THE TEST5 SIMULATES A MAS SPECTRUM.  
C  
C 6. THE TEST6 GIVES A VAS POWDER PATTERN WITH A VARIABLE  
C ANGLE 40 DEGREE.  
C  
C IMPLICIT REAL\*8 (A-H,O-Z)  
CHARACTER\*20 F20,F21,F22  
DIMENSION TITLE(18),G(1000),GK(1000),SPTR(1000),PPM(1000)  
DIMENSION ST(720),CT(720),SCT(720),CCT(720),CXCT(720)  
DIMENSION SP(721),CP(721),STP(721),CTP(721),CXP(721)  
DIMENSION HT(1000),BO1(10),BO2(10),BO3(10),BO4(10)  
DIMENSION BO5(10),BO6(10),BO7(10),FINT(10)

C  
C OPEN FILES  
C  
7 WRITE(\*,82)  
READ(\*,84,ERR=7) F20  
OPEN(UNIT=4,FILE=F20,STATUS='OLD')  
8 WRITE(\*,78)

```

READ(*,84,ERR=8) F21
OPEN(UNIT=1,FILE=F21,STATUS='NEW')
9 WRITE(*,80)
READ(*,84,ERR=9) F22
OPEN(UNIT=2,FILE=F22,STATUS='NEW')

C
C ***** FORMAT *****
C
12 FORMAT(/)
14 FORMAT(6HSPIN =,F4.1,2X,13HLARMOR FRQ. =,F7.3,1X,3HMHz,2X,
* 5HQCC =,F7.3,1X,3HMHz,2X,13HASYMM. PAR. =,F7.4)
16 FORMAT(13HSPHERE DIV. =,I4,1X,26HNO. POINTS TO BE PLOTTED =,
* F5.0,/,18HMIN. CHEM. SHIFT =,F8.2,1X,18HMAX. CHEM. SHIFT =,
* F8.2,3Hppm)
18 FORMAT(28H CHEM. SHIFT PARAMETERS;D11=,F7.2,1X,4HD22=,F7.2,1X,
* 4HD33=,F7.2,1X,3Hppm)
20 FORMAT('ISO. CHEM. SHIFT =',F7.2,1X,'CHEM. SHIFT DELTA =',F8.3,
* 1X,'ppm',3X,'CHEM SHIFT ETA=',F7.2)
22 FORMAT(22HGAUSSIAN BROADENNING =,F7.5,1X,20HGAUSSIAN LINEWIDTH
=,
* F10.5,4H kHz)
24 FORMAT(10HDELTA M =,I1,20H TRANSITIONS BETWEEN,F5.1,
* 4H AND,F5.1)
28 FORMAT(18A4)
32 FORMAT(2X,17H VARIABLE ANGLE =,F10.7,7H DEGREE,/,
* 7HALPHA =,F10.6,2X,6HBETA =,F10.6,2X,7HGAMMA =,F10.6,7H DEGREE)
34 FORMAT(2X,'ENTER TITLE')
36 FORMAT(2X,'ENTER SPIN (FOR EXAMPLE I=7/2-->3.5)')
38 FORMAT(8X,'TRANSITIONS TO BE CALCULATED(.5 FOR CENTRAL
* TRANSITION)')
40 FORMAT(8X,'LARMOR FREQUENCY IN MHZ')

```

```
42  FORMAT(8X,'NUMBER OF POINTS FOR PLOT  (<1001)')  
44  FORMAT(8X,'PLOT MINIMUM IN PPM (FOR EXAMPLE -100)')  
46  FORMAT(8X,'PLOT MAXIMUM IN PPM (FOR EXAMPLE 100)')  
48  FORMAT(8X,'DUMMY (0 FOR NORM. SIM., 2 TO ADD A SPECTRUM)')  
50  FORMAT(34H SPIN,TRAN,VL,PS,PPMIN,PPMAX,DUMMY)  
52  FORMAT(2X,'ENTER QUADRUPOLE COUPLING CONSTANT IN MHZ')  
54  FORMAT(8X,'ASYMMETRY PARAMETER OF EFG')  
56  FORMAT(8X,'3 CSA TENSOR PRINCIPAL COMPONENTS IN PPM')  
58  FORMAT(8X,'GAUSSIAN BROADENING IN HZ')  
60  FORMAT(8X,'RELATIVE INTENSITY')  
62  FORMAT(8X,'1 FOR NORMALIZATION, 0 TO ADD A SPECTRUM')  
64  FORMAT(38H QCC,ETA,D11,D22,D33,SIGHZ,WEIGHT,NORM)  
66  FORMAT(8X,'VARIABLE ANGLE (54.7356 FOR MASS, 0 FOR STATIC)')  
68  FORMAT(8X,'EULERIAN ANGLES, ACQ, BCQ, GCQ')  
70  FORMAT(8X,'NUMBER OF POINTS FOR ANGLES (<1001)')  
72  FORMAT(20H ALP,ACQ,BCQ,GCQ,ITH)  
74  FORMAT(2E12.4)  
76  FORMAT('*** THE SPECTRUM.DAT FILE HAS BEEN ADDED.  
*  ADD THE PARAMETER FILE OF THE DATA FILE ***')  
78  FORMAT(' ENTER OUTPUT FILE NAME')  
80  FORMAT(' ENTER OUTPUT PARAMETER FILE NAME')  
82  FORMAT(' ENTER INPUT PARAMETER FILE NAME')  
84  FORMAT(A)  
86  FORMAT('WEIGHT = ',1F10.6)  
88  FORMAT(' PLEASE, DO NOT INTERRUPT.')  
     IO=2  
C  
C      ***** INPUT PARAMETER *****  
C  
C      WRITE(*,34)  
C      READ(*,28)TITLE
```

```

C      WRITE(*,36)
C      WRITE(*,38)
C      WRITE(*,40)
C      WRITE(*,42)
C      WRITE(*,44)
C      WRITE(*,46)
C      WRITE(*,48)
C      WRITE(*,50)
C      READ(*,*)SPIN,TRAN,VL,PS,PPMIN,PPMAX,DUMMY
C      WRITE(*,52)
C      WRITE(*,54)
C      WRITE(*,56)
C      WRITE(*,58)
C      WRITE(*,60)
C      WRITE(*,62)
C      WRITE(*,64)
C      READ(*,*)QCC,ETA,D11,D22,D33,SIGHZ,WEIGHT,NORM
C      WRITE(*,66)
C      WRITE(*,66)
C      WRITE(*,66)
C      WRITE(*,72)
C      READ(*,*)ALP,ACQ,BCQ,GCQ,ITH
      WRITE(*,88)
      READ(4,28)TITLE
      READ(4,*)SPIN,TRAN,VL,PS,PPMIN,PPMAX,DUMMY
      READ(4,*)QCC,ETA,D11,D22,D33,SIGHZ,WEIGHT,NORM
      READ(4,*)ALP,ACQ,BCQ,GCQ,ITH
      ****
      IF(ITH.GT.720.) ITH=720.
      IF(PS.GT.1000.) PS=1000.
      IF(ITH.LT.1.) ITH=90.

```

```
IF(PS.LT.1.) PS=100.  
NS=PS  
XITH=ITH  
PMIN=VL+(PPMIN*VL/1.0E6)  
PMAX=VL+(PPMAX*VL/1.0E6)  
PR=PMAX-PMIN  
GIP=(PS-1.)/PR  
DO 110 K=1,NS  
G(K)=0.  
GK(K)=0.  
SPTR(K)=0.  
XK=K  
PK=PMIN+(XK-1.)*PR/(PS-1.)  
PPM(K)=(PK-VL)*1.0E6/VL  
110 CONTINUE  
SIGKC=SIGHZ/1000.  
SIG=SIGHZ/PR*1.E-6  
CSIGMA=(D11+D22+D33)/3.  
CSDEL=D33-CSIGMA  
IF(CSDEL.EQ.0.) GO TO 120  
CSETA=(D11-D22)/CSDEL  
GO TO 130  
120 CSETA=0.  
130 CISO=CSIGMA/1.E6  
CDEL=CSDEL/1.E6  
RAD=1.745329252E-2  
ANG=RAD*ALP  
CAL=COS(ANG)  
CXC=CAL*CAL  
CTC=CXC*CXC  
THC=.5*(3.*CXC-1.)
```

```

TTC=.5*THC
RTA=2.*RAD*ACQ
SINTA=SIN(RTA)
COSTA=COS(RTA)
RBCQ=RAD*BCQ
SINB=SIN(RBCQ)
COSB=COS(RBCQ)
SQSB=SINB*SINB
SQCB=COSB*COSB
RGCQ=RAD*GCQ
SING=SIN(RGCQ)
COSG=COS(RGCQ)
SINTG=2.*SING*COSG
COSTG=COSG*COSG-SING*SING
VLCL=VL*CDEL*THC
AAB=.5*VLCL*(3.*SQCB-1.+CSETA*SQSB*COSTA)
BAB=.25*VLCL*(3.*SQSB+CSETA*(1.+SQCB)*COSTA)
CAB=.5*VLCL*CSETA*COSB*SINTA
DAB=VLCL*SINB*COSB*(-3.+CSETA*COSTA)
EAB=-VLCL*CSETA*SINB*SINTA
BABS=BAB*SINTG
BABC=BAB*COSTG
CABS=CAB*SINTG
CABC=CAB*COSTG
DABS=DAB*SING
DABC=DAB*COSG
EABS=EAB*SING
EABC=EAB*COSG
BSCC=BABS+CABC
BCNCS=BABC-CABS
DCES=DABC+EABS

```

```
DSEC=DABS+EABC
MPH=ITH+1
IF(ACQ+BCQ+GCQ.EQ.0.) XITH=2.*XITH
DO 100 I=1,ITH
XI=I
CT(I)=1.-2.*(XI-.5)/XITH
CCT(I)=CT(I)*CT(I)
ST(I)=SQRT(1.-CCT(I))
SCT(I)=ST(I)*CT(I)
CXCT(I)=CCT(I)*CCT(I)
100 CONTINUE
IF(ETA+CSETA+ACQ+BCQ+GCQ.EQ.0) MPH=1
DO 200 M=1,MPH
XPI=M
PHI=3.1415927*(XPI-1.)/XITH
TPHI=2.*PHI
SP(M)=SIN(PHI)
CP(M)=COS(PHI)
STP(M)=SIN(TPHI)
CTP(M)=COS(TPHI)
CXP(M)=CTP(M)*CTP(M)
200 CONTINUE
EE=ETA*ETA
EE6=EE/6.
SS1=SPIN*(SPIN+1.)
IF(SPIN-0.5) 140,140,150
140 VQQ=0.0
GO TO 160
150 VQQ=3.*QCC/(2.*SPIN*(2.*SPIN-1.))
160 FACT=VQQ*VQQ/(12.*VL)
VIS=VL*(1.+CISO)-.5*AAB
```

```

FATA=FACT*ETA
XMAG=-TRAN+1.
ITRAN=2*TRAN
MAG=1
170   CONTINUE
XX1=XMAG*(XMAG-1.)
A=24.*XX1-4.*SS1+9.
B=6.*XX1-2.*SS1+3.
C=-VQQ*(XMAG-.5)
VMA=.5*A*(1.-CTC)-.125*B*(1.+6.*CXC+CTC)
VMB=.5*A*(1.-3.*CXC+4.*CTC)-.5*B*(1.-CTC)
VMC=1.5*A*(CXC-CTC)-.375*B*(1.-2.*CXC+CTC)
ETAC=ETA*C
BO1(MAG)=-TTC*C+FACT*(.375*VMA+EE6*VMB+.25*VMC)
BO2(MAG)=BCNCS+TTC*ETAC+FATA*(.25*VMA-.5*VMC)
BO3(MAG)=FACT*EE*(VMA/24.-VMB/6.+.25*VMC)
BO4(MAG)=1.5*AAB+3.*TTC*C+FACT*((EE6-.75)*VMA
* +(1.5-EE6)*VMB-1.5*VMC)
BO5(MAG)=-BCNCS-TTC*ETAC+FATA*(-VMB+2.*VMC)
BO6(MAG)=FACT*(.375*VMA-1.5*VMB+2.25*VMC)
BO7(MAG)=FATA*(-.25*VMA+VMB-1.5*VMC)
FINT(MAG)=SS1-XX1
IF(XMAG.GE.TRAN) GO TO 190
180   XMAG=XMAG+1
MAG=MAG+1
GO TO 170
190   CONTINUE
MGMX=MAG
AMP=WEIGHT/XPI
DO 300 I=1,ITH
DO 300 JP=1.MPH

```

```

DO 300 MAG=1,MGMX
BOC=BO3(MAG)*CXP(JP)
EUN1=VIS+BO1(MAG)+BO2(MAG)*CTP(JP)-BSCC*STP(JP)+BOC
EUN2=BO4(MAG)+BO5(MAG)*CTP(JP)+BSCC*STP(JP)-2.*BOC
EUN3=BO6(MAG)+BO7(MAG)*CTP(JP)+BOC
EUN4=DCES*CP(JP)+DSEC*SP(JP)
FNN=EUN1+EUN2*CCT(I)+EUN3*CXCT(I)+EUN4*SCT(I)-PMIN
NN=FNN*GIP+1.5
IF(NN.LT.2.OR.NN.GT.NS) GO TO 300
GK(NN)=GK(NN)+FINT(MAG)

300 CONTINUE
DO 210 LL=1,NS
GK(LL)=GK(LL)*AMP.

210 CONTINUE
CALL SHIFT(NS,GK)
CALL CONVOL(NS,SIG,GK,0,SPTR)
IF(DUMMY.NE.2) GO TO 230
OPEN(UNIT=3,FILE='SPECTRUM.DAT',STATUS='OLD')
READ(3,74)(PPM(K),HT(K),K=1,NS)
WRITE(*,76)
DO 220 N=1,NS
SPTR(N)=SPTR(N)+HT(N)

220 CONTINUE
230 IF(NORM.NE.1) GO TO 240
CALL UNITY(NS,SPTR)
C **** OUTPUT ****
C
240 WRITE(1,74)(PPM(K),SPTR(K),K=1,NS)
C
C ***** OUTPUT FOR INPUT PARAMETER *****
C

```

```

        WRITE(IO,28)TITLE
        WRITE(IO,12)
        WRITE(IO,16)ITH,PS,PPMIN,PPMAX
        WRITE(IO,12)
        WRITE(IO,14)SPIN,VL,QCC,ETA
        WRITE(IO,24)ITRAN,TRAN,-TRAN
        WRITE(IO,12)
        WRITE(IO,18)D11,D22,D33
        WRITE(IO,20)CSIGMA,CSDEL,CSETA
        WRITE(IO,12)
        WRITE(IO,32)ALP,ACQ,BCQ,GCQ
        WRITE(IO,12)
        WRITE(IO,22)SIG,SIGKC
        WRITE(IO,12)
        WRITE(IO,86)WEIGHT
C ****
C      STOP
C      END
C
C      SUBROUTINE CONVOL(NS,SIGMA,G,NDERIV,SA)
C      CONVOL DOES A CONVOLUTION OF THE NTH GAUSSIAN DERIVATIVE WITH
C      AN ARBITRARY FUNCTION. N IS 0,1,2, OR 3. ENTRY IS NS POINTS,
C      SIGMA THE GAUSSIAN HALF-WIDTH, G CONTAINING THE FUNCTION TO BE
C      CONVOLUTED AND NDERIV THE VALUE OF N. THE OUTPUT IS IN SA.
C
C      IMPLICIT REAL*8 (A-H,O-Z)
C      DIMENSION G(1000),SA(1000),XSTORE(1000)
C      IF(SIGMA.LE.0.) RETURN
C      RX=1./(2.*(SIGMA*NS)**2)
C      SQRX=RX**.5
C      DO 28 L=1,NS

```

```
SA(L)=0.  
XL=L-1  
POW=XL*XL*RX  
IF(POW-30.) 26,27,27  
26  XSTORE(L)=EXP(-POW)  
29  GO TO 28  
27  XSTORE(L)=0.  
28  CONTINUE  
    IF(NDERIV.LT.0.OR.NDERIV.GT.3) NDERIV=1  
    MOVE=NDERIV+1  
    GO TO (20,30,40,50) MOVE  
20  CONTINUE  
    DO 22 I=1,NS  
    DO 22 J=1,NS  
    XXX=I-J  
    IF(XXX) 23,24,24  
23  L=J-I+1  
    GO TO 22  
24  L=I-J+1  
22  SA(J)=SA(J)+G(I)*XSTORE(L)  
    RETURN  
30  CONTINUE  
    DO 32 I=1,NS  
    DO 32 J=1,NS  
    XXX=I-J  
    IF(XXX) 33,34,34  
33  L=J-I+1  
    GO TO 32  
34  L=I-J+1  
32  SA(J)=SA(J)+G(I)*XSTORE(L)*XXX*SQRX  
    RETURN
```

```

40    CONTINUE
      DO 42 I=1,NS
      DO 42 J=1,NS
      XXX=I-J
      IF(XXX) 43,44,44
43    L=J-I+1
      GO TO 42
44    L=I-J+1
42    SA(J)=SA(J)+G(I)*XSTORE(L)*(1.-2.*XXX*XXX*RX)
      RETURN
50    CONTINUE
      DO 52 I=1,NS
      DO 52 J=1,NS
      XXX=I-J
      IF(XXX) 53,54,54
53    L=J-I+1
      GO TO 52
54    L=I-J+1
52    SA(J)=SA(J)-G(I)*XSTORE(L)*(-6.*XXX*SQRX
*   +4.*XXX*XXX*XXX*RX*SQRX)
      RETURN
      END
C
C     SUBROUTINE SHIFT(NS,G)
C     SHIFT SUBTRACTS A FIXED AMOUNT FROM THE POWDER
C     PATTERN TO MINIMIZE THE EFFECTS OF THE CONVOLUTION
C     AT THE ENDS OF THE SIMULATED SPECTRUM.
C
C     IMPLICIT REAL*8 (A-H,O-Z)
C     DIMENSION G(1000)
C     GMIN=G(1)

```

```
DO 1 I=2,NS
IF(GMIN.LE.G(I)) GO TO 1
GMIN=G(I)
1    CONTINUE
DO 2 I=1,NS
G(I)=G(I)-GMIN
2    CONTINUE
RETURN
END

C
SUBROUTINE UNITY(NS,SPTR)
C      UNITY NORMALIZES INTENSITY AFTER CONVOLUTION (MAXIMUM PEAK
C      HEIGHT = 1.).
C      IMPLICIT REAL*8 (A-H,O-Z)
DIMENSION SPTR(1000)
STMAG=SPTR(1)
DO 1 I=1,NS
IF(STMAG.GE.SPTR(I)) GO TO 1
STMAG=SPTR(I)
1    CONTINUE
DO 2 I=1,NS
SPTR(I)=SPTR(I)/STMAG
2    CONTINUE
RETURN
END
```