# A STUDY ON THE RECURRENCE RELATIONS AND <br> VECTORS $X_{\lambda}, S_{\lambda}$ AND $U_{\lambda}$ IN $g-E S X_{n}$ 

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#### Abstract

The manifold $g-E S X_{n}$ is a generalized $n$-dimensional Riemannian manifold on which the differential geometric structure is imposed by the unified field tensor $g_{\lambda \mu}$ through the $E S$-connection which is both Einstein and semi-symmetric. In this paper, we investigate the properties of the vectors $X_{\lambda}, S_{\lambda}$ and $U_{\lambda}$ of $g-E S X_{n}$, with main emphasis on the derivation of several useful generalized identities involving it.


## 1. Introduction

Manifolds with recurrent connections have been studied by many authors, such as Chung, Datta, E.M. Patterson, M.Pravanovitch, Singal, and Takano, etc(refer to [3] and [4]). Examples of such manifolds are those of recurrent curvature, Ricci-recurrent manifolds, and bi-recurrent manifolds.

In this paper, we introduce a new concept of semi-symmetric connection $\Gamma_{\lambda}{ }^{\nu}{ }_{\mu}$ on a generalized $n$-dimensional Riemannian manifold $X_{n}$ and study its recurrence relations in the first. In the second, we investigate the properties of the vectors $X_{\lambda}, S_{\lambda}$ and $U_{\lambda}$ of $g-E S X_{n}$.

The main purpose of the present paper is to obtain several basic identities satisfied by the vectors $X_{\lambda}, S_{\lambda}$ and $U_{\lambda}$ and recurrence relations in $g-E S X_{n}$ which is both semi-symmetric and Einstein .

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## 2. Preliminaries

This section is a brief collection of basic concepts, results, and notations needed in subsequent considerations. They are due to Chung ( [3], 1963), Hwang ([2], 1988), and Mishra([7], 1959) mostly due to [6].
(a) generalized $n$-dimensional Riemannian manifold $X_{n}$

Let $X_{n}$ be a generalized $n$-dimensional Riemannian manifold referred to a real coordinate system $x^{\nu}$, which obeys the coordinate transformations $x^{\nu} \rightarrow x^{\nu^{\prime}}$ for which

$$
\begin{equation*}
\operatorname{det}\left(\frac{\partial x^{\prime}}{\partial x}\right) \neq 0 \tag{2.1}
\end{equation*}
$$

In $n-g-U F T$ the manifold $X_{n}$ is endowed with a real nonsymmetric tensor $g_{\lambda \mu}$, which may be decomposed into its symmetric part $h_{\lambda \mu}$ and skew-symmetric part $k_{\lambda \mu}$ :

$$
\begin{array}{rlrl}
g_{\lambda \mu} & =h_{\lambda \mu}+k_{\lambda \mu}, & \quad \text { where } & \\
\mathfrak{g} & =\operatorname{det}\left(g_{\lambda \mu}\right) \neq 0, \quad \mathfrak{h}=\operatorname{det}\left(h_{\lambda \mu}\right) \neq 0, \quad \mathfrak{k}=\operatorname{det}\left(k_{\lambda \mu}\right) \tag{2.2b}
\end{array}
$$

In virtue of $(2.2 b)$ we may define a unique tensor $h^{\lambda \nu}$ by

$$
\begin{equation*}
h_{\lambda \mu} h^{\lambda \nu}=\delta_{\mu}^{\nu} \tag{2.3}
\end{equation*}
$$

which together with $h_{\lambda \mu}$ will serve for raising and/or lowering indices of tensors in $X_{n}$ in the usual manner. There exists a unique tensor $* g^{\lambda \nu}$ satisfying

$$
\begin{equation*}
g_{\lambda \mu}{ }^{*} g^{\lambda \nu}=g_{\mu \lambda}{ }^{*} g^{\nu \lambda}=\delta_{\mu}^{\nu} \tag{2.4}
\end{equation*}
$$

It may be also decomposed into its symmetric part ${ }^{*} h_{\lambda \mu}$ and skewsymmetric part ${ }^{*} k_{\lambda \mu}$ :

$$
\begin{equation*}
{ }^{*} g^{\lambda \nu}={ }^{*} h^{\lambda \nu}+{ }^{*} k^{\lambda \nu} \tag{2.5}
\end{equation*}
$$

The manifold $X_{n}$ is connected by a general real connection $\Gamma_{\lambda}{ }^{\nu}{ }_{\mu}$ with the following transformation rule:

$$
\begin{equation*}
\Gamma_{\lambda^{\prime}}^{\nu^{\prime}}{ }_{\mu^{\prime}}=\frac{\partial x^{\nu^{\prime}}}{\partial x^{\alpha}}\left(\frac{\partial x^{\beta}}{\partial x^{\lambda^{\prime}}} \frac{\partial x^{\gamma}}{\partial x^{\mu^{\prime}}} \Gamma_{\beta}{ }^{\alpha}{ }_{\gamma}+\frac{\partial^{2} x^{\alpha}}{\partial x^{\lambda^{\prime}} \partial x^{\mu^{\prime}}}\right) \tag{2.6}
\end{equation*}
$$

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It may also be decomposed into its symmetric part $\Lambda_{\lambda}{ }^{\nu}{ }_{\mu}$ and its skewsymmetric part $S_{\lambda \nu}{ }^{\nu}$, called the torsion of $\Gamma_{\lambda}{ }^{\nu}{ }_{\mu}$ :

$$
\begin{equation*}
\Gamma_{\lambda}{ }^{\nu}{ }_{\mu}=\Lambda_{\lambda}{ }^{\nu}{ }_{\mu}+S_{\lambda \mu}{ }^{\nu} ; \quad \Lambda_{\lambda}{ }^{\nu}{ }_{\mu}=\Gamma_{\left(\lambda^{\nu}{ }_{\mu)} ; \quad S_{\lambda \mu}{ }^{\nu}=\Gamma_{[\lambda}{ }^{\nu}{ }_{\mu]}, ~\right.}^{0} \tag{2.7}
\end{equation*}
$$

A connection $\Gamma_{\lambda}{ }^{\nu}{ }_{\mu}$ is said to be Einstein if it satisfies the following system of Einstein's equations:

$$
\begin{array}{rlr}
\partial_{\omega} g_{\lambda \mu} & -\Gamma_{\lambda}{ }^{\alpha}{ }_{\omega} g_{\alpha \mu}-\Gamma_{\omega}{ }^{\alpha}{ }_{\mu} g_{\lambda \alpha}=0, & \text { or equivalently } \\
D_{\omega} g_{\lambda \mu} & =2 S_{\omega \mu}{ }^{\alpha} g_{\lambda \alpha} & \tag{2.8b}
\end{array}
$$

where $D_{\omega}$ is the symbolic vector of the covariant derivative with respect to $\Gamma_{\lambda}{ }^{\nu}{ }_{\mu}$. In order to obtain $g_{\lambda \mu}$ involved in the solution for $\Gamma_{\lambda}{ }^{\nu}{ }_{\mu}$ in (2.8), certain conditions are imposed. These conditions may be condensed to

$$
\begin{equation*}
S_{\lambda}=S_{\lambda \alpha}^{\alpha}=0, \quad R_{[\mu \lambda]}=\partial_{[\mu} Y_{\lambda]}, \quad R_{(\mu \lambda)}=0 \tag{2.9}
\end{equation*}
$$

where $Y_{\lambda}$ is an arbitrary vector, and

$$
\begin{equation*}
\left.R_{\omega \mu \lambda}{ }^{\nu}=2\left(\partial_{[\mu} \Gamma_{|\lambda|}{ }^{\nu} \omega\right]+\Gamma_{\alpha}{ }^{\nu}\left[\mu \Gamma_{|\lambda|}{ }^{\alpha} \omega\right]\right), \quad R_{\mu \lambda}=R_{\alpha \mu \lambda}{ }^{\alpha} \tag{2.10}
\end{equation*}
$$

If the system (2.8) admits a solution $\Gamma_{\lambda}{ }^{\nu}{ }_{\mu}$, it must be of the form (Hlavatý, 1957)

$$
\Gamma_{\lambda}^{\nu}{ }_{\mu}=\left\{\begin{array}{l}
\nu  \tag{2.11}\\
\lambda \mu
\end{array}\right\}+S_{\lambda \mu}{ }^{\nu}+U^{\nu}{ }_{\lambda \mu}
$$

where $U^{\nu}{ }_{\lambda \mu}=2 h^{\nu \alpha} S_{\alpha(\lambda}{ }^{\beta} k_{\mu) \beta}$ and $\left\{\begin{array}{l}\nu \\ \lambda \mu\end{array}\right\}$ are Christoffel symbols defined by $h_{\lambda \mu}$
(b) Some notations and results The following quantities are frequently used in our further considerations:

$$
\begin{gather*}
g=\frac{\mathfrak{g}}{\mathfrak{h}}, \quad k=\frac{\mathfrak{k}}{\mathfrak{h}}  \tag{2.12}\\
K_{p}=k_{\left[\alpha_{1}\right.}{ }^{\alpha_{1}}{k_{\alpha_{2}}}^{\alpha_{2}} \cdots k_{\left.\alpha_{p}\right]}{ }^{\alpha^{p}}, \quad(p=0,1,2, \cdots)  \tag{2.13}\\
{ }^{(0)} k_{\lambda}{ }^{\nu}=\delta_{\lambda}^{\nu}, \quad{ }^{(p)} k_{\lambda}{ }^{\nu}=k_{\lambda}{ }^{\alpha}{ }^{(p-1)} k_{\alpha}{ }^{\nu} \quad(p=1,2, \cdots) \tag{2.14}
\end{gather*}
$$

In $X_{n}$ it was proved in [3] that
$K_{0}=1, \quad K_{n}=k \quad$ if $n$ is even, and $K_{p}=0 \quad$ if $p$ is odd
(2.16) $\mathfrak{g}=\mathfrak{h}\left(1+K_{1}+K_{2}+\cdots+K_{n}\right) \quad$ or $\quad g=1+K_{1}+K_{2}+\cdots+K_{n}$

$$
\begin{equation*}
\sum_{s=0}^{n-\sigma} K_{s}^{(n-s+p)} k_{\lambda}^{\nu}=0 \quad(p=01,2, \cdots) \tag{2.17}
\end{equation*}
$$

We also use the following useful abbreviations for an arbitrary vector $Y$, for $p=1,2,3, \cdots$ :

$$
\begin{align*}
& { }^{(p)} Y_{\lambda}={ }^{(p-1)} k_{\lambda}{ }^{\alpha} Y_{\alpha}  \tag{2.18}\\
& { }^{(p)} Y^{\nu}={ }^{(p-1)} k^{\nu}{ }_{\alpha} Y^{\alpha} \tag{2.19}
\end{align*}
$$

(c) n-dimensional $E S$ manifold $E S X_{n}$

In this subsection, we display an useful representation of the $E S$ connection in $n-g$-UFT.

Definition 2.1. A connection $\Gamma_{\lambda}{ }^{\nu}{ }_{\mu}$ is said to be semi-symmetric if its torsion tensor $S_{\lambda \mu}{ }^{\nu}$ is of the form

$$
\begin{equation*}
S_{\lambda \mu}{ }^{\nu}=2 \delta_{[\lambda}^{\nu} X_{\mu]} \tag{2.20}
\end{equation*}
$$

for an arbitrary non-null vector $X_{\mu}$. A connection which is both semisymmetric and Einstein is called a ES-connection. An n-dimensional generalized Riemannian manifold $X_{n}$, on which the differential geometric structure is imposed by $g_{\lambda \mu}$ by means of a $E S$-connection, is called an $n$-dimensional $E S$-manifold. We denote this manifold by $g-E S X_{n}$ in our further considerations.

Theorem 2.2. Under the condition (2.20), the system of equations (2.8) is equivalent to

$$
\Gamma_{\lambda}^{\nu}{ }_{\mu}=\left\{\begin{array}{l}
\nu  \tag{2.21}\\
\lambda \mu
\end{array}\right\}+2 k_{(\lambda}^{\nu} X_{\mu)}+2 \delta_{[\lambda}^{\nu} X_{\mu]}
$$

Proof. Substituting (2.20) for $S_{\lambda \mu}{ }^{\nu}$ into (2.11), we have the representation (2.21).

## 3. Properties of the vectors $X_{\lambda}, S_{\lambda}$ and $U_{\lambda}$

This section is concerned with identities satisfied by the vectors $X_{\lambda}$, given by (2.19) and the vectors $S_{\lambda}$ in (2.7) and

$$
\begin{equation*}
U_{\lambda}=U^{\alpha}{ }_{\lambda \alpha} \tag{3.1}
\end{equation*}
$$

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Theorem 3.1. In $g-E S X_{n}$ under present conditions, the following recurrence relation hold:

$$
\begin{equation*}
\sum_{s=0}^{n-\sigma} K_{s}^{(n-s+p)} k_{\lambda}^{\nu}=0 \quad(p=0,1,2, \cdots) \tag{3.2}
\end{equation*}
$$

where

$$
\sigma= \begin{cases}0 & \text { if } n \text { is even } \\ 1 & \text { if } n \text { is odd }\end{cases}
$$

Proof. The relation (3.2) is a direct consequence of (2.17) and (2.18).

Theorem 3.2. In $g-E S X_{n}$, the vectors $S_{\lambda}$ and $U_{\lambda}$ are given by

$$
\begin{gather*}
S_{\lambda}=(1-n) X_{\lambda}  \tag{3.3}\\
U_{\lambda}=\frac{1}{2} \partial_{\lambda} \ln \mathfrak{g} \tag{3.4}
\end{gather*}
$$

Proof. Putting $\mu=\nu$ in (2.20), we have (3.3). In order to prove (3.4), consider the following Einstein's equations (2.8a). Multiplying * $g^{\lambda \mu}$ to both sides of (2.8a) and making use of (2.4), we have

$$
\begin{equation*}
\partial_{\omega} \ln \mathfrak{g}-\Gamma_{\alpha}{ }^{\alpha}{ }_{\omega}-\Gamma_{\omega}{ }^{\alpha}{ }_{\alpha}=0 \tag{3.5}
\end{equation*}
$$

or equivalently

$$
\begin{equation*}
\partial_{\omega} \ln \mathfrak{g}+2 S_{\omega}-2 \Gamma_{\omega}{ }^{\alpha}{ }_{\alpha}=0 \tag{3.6}
\end{equation*}
$$

On the other hand, in virtue of classical result

$$
\left\{\begin{array}{l}
\alpha  \tag{3.7}\\
\omega \alpha
\end{array}\right\}=\frac{1}{2} \partial_{\omega} \ln \mathfrak{h}
$$

the result (2.2) gives

$$
\begin{equation*}
\Gamma_{\omega}{ }^{\alpha}{ }_{\alpha}=\frac{1}{2} \partial_{\omega} \ln \mathfrak{h}+S_{\omega}+U_{\omega} \tag{3.8}
\end{equation*}
$$

The relation (3.4) immediately follows from (3.6) and (3.8).
Theorem 3.3. In $g-E S X_{n}$, the following relations hold for $p, q=$ $1,2,3, \cdots$ :

$$
\begin{gather*}
{ }^{(p+1)} S_{\lambda}=(1-n)^{(p)} U_{\lambda}  \tag{3.9}\\
{ }^{(p)} U_{\alpha}^{(q)} X^{\alpha}=0 \quad \text { if } \quad p+q-1 \quad \text { is } \quad \text { odd } \tag{3.10}
\end{gather*}
$$

Proof. The relation (3.9) is a direct consequence of (3.3), (3.4), and (2.18). Making use of (3.9), the relation

$$
\begin{equation*}
{ }^{(p)} U_{\alpha}^{(q)} X^{\alpha}={ }^{(p+1)} X_{\alpha}^{(q)} X^{\alpha}=(-1)^{q(p+q-1)} k_{\alpha \beta} X^{\alpha} X^{\beta} \tag{3.11}
\end{equation*}
$$

follows. The statement (3.10) may be proved from (3.11), since ${ }^{(p+q-1)} k_{\alpha \beta}$ is skew-symmetric if $p+q-1$ is odd.

Theorem 3.4. In $g-E S X_{n}$, the following relations hold:

$$
\begin{gather*}
D_{\lambda} X_{\mu}=\nabla_{\lambda} X_{\mu}  \tag{3.12}\\
D_{[\lambda} X_{\mu]}=\nabla_{[\lambda} X_{\mu]}=\partial_{[\lambda} X_{\mu]}  \tag{3.13}\\
\nabla_{[\lambda} U_{\mu]}=0, \quad D_{[\lambda} U_{\mu]}=2 U_{[\lambda} X_{\mu]}=2^{(2)} X_{[\lambda} X_{\mu]} \tag{3.14}
\end{gather*}
$$

where $\nabla_{\omega}$ is the symbolic vector of the covariant derivative with respect to the Christoffel symbols defined by $h_{\lambda \mu}$.

Proof. In virtue of (2.18) and Theorem 2.2, the relation (3.11) follows as in the following way:

$$
\begin{aligned}
D_{\lambda} X_{\mu} & =\nabla_{\lambda} X_{\mu}-X_{\alpha} S_{\mu \lambda}{ }^{\alpha}-X_{\alpha} U^{\alpha}{ }_{\mu \lambda} \\
& =\nabla_{\lambda} X_{\mu}-2 X_{[\mu} X_{\lambda]}+h_{\mu \lambda}\left(k_{\alpha \beta} X^{\alpha} X^{\beta}\right) \\
& =\nabla_{\alpha} X_{\mu}
\end{aligned}
$$

The relation (3.13) are direct consequences of (3.12). Since $\partial_{[\lambda} U_{\mu]}=0$ in virtue of the (3.4), we have the first relation of (3.14). Similarly, the second relation of (3.14) may be proved in virtue of (3.4).

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