# STABILITY OF THE MILSTEIN METHOD FOR STOCHASTIC DIFFERENTIAL EQUATIONS WITH JUMPS ${ }^{\dagger}$ 

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#### Abstract

In this paper the Milstein method is proposed to approximate the solution of a linear stochastic differential equation with Poisson-driven jumps. The strong Milstein method and the weak Milstein method are shown to capture the mean square stability of the system. Furthermore using some technique, our result shows that these two kinds of Milstein methods can well reproduce the stochastically asymptotical stability of the system for all sufficiently small time-steps. Some numerical experiments are given to demonstrate the conclusions.


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## 1. Introduction

Stochastic differential equations (SDEs) with jumps have many important applications in the areas such as biology, finance, economics and so on. In general, SDEs with jumps have no explicit solutions. Thus, it is necessary to develop numerical methods and to study the properties of these numerical schemes. There is a vast literature on the stability of numerical methods for SDEs, see, for example, Higham [5], Milstein [12], Higham et al.[7] and for stochastic differential delay equations (SDDEs), for example, see [17, 9]. There are some results concerned with the weak approximate schemes for SDEs, see, for example, Saito and Mistui [14], Burrage et al. [2], Cao and Liu [4]. Recently, there has been some work done about the numerical schemes for SDEs with jumps [6, 3] and SDDEs with jumps [10, 15].

[^0]However, the numerical methods in $[5,7,6,3,10,15]$ converge with strong order of one half. In this paper, we will focus on the Milstein method which has strong convergence rate of one, see, for example, [11]. There is an extensive literature on the numerical analysis of the Milstein method (e.g., [16, 13]). Our aim in this work is to investigate the mean square stability and the asymptotical stability of the Milstein method for SDEs with jumps.

In this paper, we consider the linear, scalar SDEs with jumps:

$$
\begin{equation*}
d x(t)=a x\left(t^{-}\right) d t+b x\left(t^{-}\right) d W(t)+c x\left(t^{-}\right) d N(t) \tag{1}
\end{equation*}
$$

with $x(0)=x_{0}$. Here $t>0, x_{0} \neq 0$ with probability one, $x\left(t^{-}\right)$denotes $\lim _{s \rightarrow t^{-}} x(s)$. Here, $W(t)$ is a scalar Brownian motion and $N(t)$ is a scalar Poisson process with intensity $\lambda(\lambda>0)$, both defined on an appropriate complete probability space $\left(\Omega, \mathcal{F},\left\{\mathcal{F}_{t}\right\}_{t \geqslant 0}, \mathbb{P}\right)$, with a filtration $\left\{\mathcal{F}_{t}\right\}_{t \geqslant 0}$ satisfying the usual conditions(i.e. it is increasing and right-continuous while $\mathcal{F}_{0}$ contains all $\mathbb{P}$-null sets). $W(t)$ is independent of $N(t)$. The coefficients $a, b, c \in \mathbb{R}$ and $c \neq 0$ (if $c=0$, the model reduces to a SDEs without jumps).

The main results is organized as follows. Sections 2 and 3 deal with the mean square stabilities of the strong Milstein method and the weak Milstein method respectively. In Sections 4 and 5, we apply our asymptotic stability analysis to the strong Milstein method and the weak Milstein method. In addition, some numerical results are presented to confirm the theoretical results in Section 6.

## 2. mean square stability of the strong Milstein method

Given a stepsize $h>0$, the strong Milstein scheme [11] is proposed for (1) as follows.

$$
\begin{align*}
& Y_{n+1}=\left[1+\left(a-\frac{1}{2} b^{2}\right) h+b \Delta W_{n}+\frac{1}{2} b^{2}\left(\Delta W_{n}\right)^{2}+\frac{1}{2}\left(2 c-c^{2}\right) \Delta N_{n}\right.  \tag{2}\\
& \left.+b c \Delta W_{n} \Delta N_{n}+\frac{1}{2} c^{2}\left(\Delta N_{n}\right)^{2}\right] Y_{n}
\end{align*}
$$

with $Y(0)=x_{0}$. Here $Y_{n}$ is an approximation to $x\left(t_{n}\right)$ with $t_{n}=n h, \Delta W_{n}=$ $W\left(t_{n+1}\right)-W\left(t_{n}\right), \Delta N_{n}=N\left(t_{n+1}\right)-N\left(t_{n}\right)$ and $\Delta W_{n}$ is independent of $\Delta N_{n}$.

To begin with, we cite here a linear mean square stability result of underlying the system (1)

Lemma 2.1 ([6]). The analytical solution of the system (1) is mean square stable, that is to say, $\lim _{t \rightarrow \infty} \mathbb{E}(x(t))^{2}=0$, if and only if

$$
\begin{equation*}
2 a+b^{2}+\lambda c(2+c)<0 \tag{3}
\end{equation*}
$$

Given parameters $a, b, c, \lambda$ and stepsize $h$, we say the Milstein method is mean-square stable if $\lim _{n \rightarrow \infty} \mathbb{E}\left(Y_{n}\right)^{2}=0$ for any $Y_{0}$. We are now in a position to state our result on the mean-square stability of the strong Milstein method.

Theorem 2.1. Assume the condition (3) is satisfied. Then the strong Milstein method (2) is mean square stable.

Proof. It follows from (2) that

$$
\begin{aligned}
& \mathbb{E} Y_{n+1}^{2} \\
= & \left\{1+\left(a-\frac{1}{2} b^{2}\right)^{2} h^{2}+b^{2} \mathbb{E}\left(\Delta W_{n}\right)^{2}+\frac{1}{4} b^{4} \mathbb{E}\left(\Delta W_{n}\right)^{4}+\frac{1}{4}\left(2 c-c^{2}\right)^{2} \mathbb{E}\left(\Delta N_{n}\right)^{2}\right. \\
+ & b^{2} c^{2} \mathbb{E}\left[\left(\Delta W_{n}\right)^{2}\left(\Delta N_{n}\right)^{2}\right]+\frac{1}{4} c^{4} \mathbb{E}\left(\Delta N_{n}\right)^{4}+2\left(a-\frac{1}{2} b^{2}\right) h+2 b \mathbb{E} \Delta W_{n} \\
+ & b^{2} \mathbb{E}\left(\Delta W_{n}\right)^{2}+\left(2 c-c^{2}\right) \mathbb{E} \Delta N_{n}+2 b c \mathbb{E}\left[\Delta W_{n} \Delta N_{n}\right]+c^{2} \mathbb{E}\left(\Delta N_{n}\right)^{2} \\
+ & 2\left(a-\frac{1}{2} b^{2}\right) h b \mathbb{E} \Delta W_{n}+\left(a-\frac{1}{2} b^{2}\right) h b^{2} \mathbb{E}\left(\Delta W_{n}\right)^{2}+\left(a-\frac{1}{2} b^{2}\right)\left(2 c-c^{2}\right) h \mathbb{E} \Delta N_{n} \\
+ & 2\left(a-\frac{1}{2} b^{2}\right) b c h \mathbb{E}\left[\Delta W_{n} \Delta N_{n}\right]+\left(a-\frac{1}{2} b^{2}\right) c^{2} h \mathbb{E}\left(\Delta N_{n}\right)^{2}+b^{3} \mathbb{E}\left(\Delta W_{n}\right)^{3} \\
+ & \left(2 c-c^{2}\right) b \mathbb{E}\left[\Delta W_{n} \Delta N_{n}\right]+2 b^{2} c \mathbb{E}\left[\left(\Delta W_{n}\right)^{2} \Delta N_{n}\right]+c^{2} b \mathbb{E}\left[\left(\Delta N_{n}\right)^{2} \Delta W_{n}\right] \\
+ & \frac{1}{2} b^{2}\left(2 c-c^{2}\right) \mathbb{E}\left[\left(\Delta W_{n}\right)^{2} \Delta N_{n}\right]+b^{3} c \mathbb{E}\left[\left(\Delta W_{n}\right)^{3} \Delta N_{n}\right]+\frac{1}{2} b^{2} c^{2} \mathbb{E}\left[\left(\Delta W_{n}\right)^{2}\left(\Delta N_{n}\right)^{2}\right] \\
+ & \left.\left(2 c-c^{2}\right) b c \mathbb{E}\left[\Delta W_{n}\left(\Delta N_{n}\right)^{2}\right]+\frac{1}{2} c^{2}\left(2 c-c^{2}\right) \mathbb{E}\left(\Delta N_{n}\right)^{3}+b c^{3} \mathbb{E}\left[\Delta W_{n}\left(\Delta N_{n}\right)^{3}\right]\right\} \mathbb{E} Y_{n}^{2} \\
= & {\left[1+R_{1}(a, b, c, \lambda) h+R_{2}(a, b, c, \lambda) h^{2}+R_{3}(a, b, c, \lambda) h^{3}+R_{4}(a, b, c, \lambda) h^{4}\right] \mathbb{E} Y_{n}^{2}, }
\end{aligned}
$$

where $R_{1}(a, b, c, \lambda)=2 a+b^{2}+\lambda c(c+2), R_{3}(a, b, c, \lambda)=b^{2} c^{2} \lambda^{2}+c^{4} \lambda^{3}+a c^{2} \lambda^{2}+$ $c^{3} \lambda^{3}, R_{2}(a, b, c, \lambda)=\frac{1}{2} b^{4}+a^{2}+b^{2} c^{2} \lambda+\frac{1}{2} \lambda^{2} c^{4}+2 \lambda^{2} c^{2}+2 b^{2} c \lambda+2 a c \lambda+2 c^{3} \lambda^{2}$, $R_{4}(a, b, c, \lambda)=\frac{1}{4} c^{4} \lambda^{4}$. Here the fact has been used that $\mathbb{E}\left(\Delta W_{n}\right)=\mathbb{E}\left(\Delta W_{n}\right)^{3}=$ $0, \mathbb{E}\left(\Delta W_{n}\right)^{2}=h, \mathbb{E}\left(\Delta W_{n}\right)^{4}=3 h^{2}$, and the Poisson increments satisfy $\mathbb{E}\left(\Delta N_{n}\right)=$ $\lambda h, \mathbb{E}\left(\Delta N_{n}\right)^{2}=\lambda h(1+\lambda h), \mathbb{E}\left(\Delta N_{n}\right)^{3}=\lambda h+3 \lambda^{2} h^{2}+\lambda^{3} h^{3}, \mathbb{E}\left(\Delta N_{n}\right)^{4}=$ $\lambda h+7 \lambda^{2} h^{2}+6 \lambda^{3} h^{3}+\lambda^{4} h^{4}$. Thus the numerical solution is mean square stable if and only if

$$
R_{1}(a, b, c, \lambda) h+R_{2}(a, b, c, \lambda) h^{2}+R_{3}(a, b, c, \lambda) h^{3}+R_{4}(a, b, c, \lambda) h^{4}<0
$$

That is to say

$$
R_{1}(a, b, c, \lambda)+R_{2}(a, b, c, \lambda) h+R_{3}(a, b, c, \lambda) h^{2}+R_{4}(a, b, c, \lambda) h^{3}<0
$$

Letting $\psi(h)=R_{1}(a, b, c, \lambda)+R_{2}(a, b, c, \lambda) h+R_{3}(a, b, c, \lambda) h^{2}+R_{4}(a, b, c, \lambda) h^{3}$, by (3), we easily obtain $\psi(0)<0$. Because of the continuity of $\psi$ with respect to $h$, there must exist a $h_{0}(a, b, c, \lambda) \in(0, \mu), \mu$ is small enough, such that for any $h \in\left(0, h_{0}(a, b, c, \lambda)\right)$

$$
\psi(h)<0 .
$$

The proof is completed.

## 3. Mean square stability of the weak Milstein method

Applying the weak Milstein method, which is equipped with two-point random variables for the driving process, to system (1) leads to

$$
\begin{align*}
& Y_{n+1}=\left[1+\left(a-\frac{1}{2} b^{2}\right) h+b \widehat{\Delta W_{n}}+\frac{1}{2} b^{2}\left(\widehat{\Delta W_{n}}\right)^{2}+\frac{1}{2}\left(2 c-c^{2}\right) \widehat{\Delta N_{n}}\right.  \tag{4}\\
& \left.+b c \widehat{\Delta W_{n}} \widehat{\Delta N_{n}}+\frac{1}{2} c^{2}\left(\widehat{\Delta N_{n}}\right)^{2}\right] Y_{n}
\end{align*}
$$

with $Y(0)=x_{0}$. Here $\mathbb{P}\left(\widehat{\Delta W_{n}}=\sqrt{h}\right)=\mathbb{P}\left(\widehat{\Delta W_{n}}=-\sqrt{h}\right)=1 / 2$ and $\mathbb{P}\left(\widehat{\Delta N_{n}}=\right.$ $0)=1-\lambda h, \mathbb{P}\left(\widehat{\Delta N_{n}}=1\right)=\lambda h$.

Now we give the main theorem in this section.
Theorem 3.1. Assume the condition (3) is satisfied. Then the weak Milstein method (4) is mean square stable.

Proof. It follows from (4) that

$$
\begin{align*}
& \mathbb{E} Y_{n+1}^{2} \\
= & {\left[1+\widehat{R}_{1}(a, b, c, \lambda) h+\widehat{R}_{2}(a, b, c, \lambda) h^{2}\right] \mathbb{E} Y_{n}^{2}, } \tag{5}
\end{align*}
$$

where $\widehat{R}_{1}(a, b, c, \lambda)=2\left(a-\frac{1}{2} b^{2}\right)+\lambda c(2+c), \widehat{R}_{2}(a, b, c, \lambda)=a^{2}+b^{2} c^{2} \lambda+2 a c \lambda+$ $2 b^{2} c \lambda$. Here the fact has been used that $\mathbb{E}\left(\widehat{\Delta W_{n}}\right)=\mathbb{E}\left(\widehat{\Delta W_{n}}\right)^{3}=0, \mathbb{E}\left(\widehat{\Delta W_{n}}\right)^{2}=$ $h, \mathbb{E}\left(\widehat{\Delta W_{n}}\right)^{4}=h^{2}, \mathbb{E}\left(\widehat{\Delta N_{n}}\right)^{i}=\lambda h(i=1, \cdots, 4)$. Hence the weak Milstein method is mean square stable if and only if

$$
\begin{equation*}
\widehat{R}_{1}(a, b, c, \lambda)+\widehat{R}_{2}(a, b, c, \lambda) h<0 \tag{6}
\end{equation*}
$$

If $\widehat{R}_{2}(a, b, c, \lambda) \leq 0$, by (3), then (6) holds for all $h \in(0,1 / \lambda)$. If $\widehat{R}_{2}(a, b, c, \lambda)>0$, by (3), then (6) holds for all $h \in\left(0,-\frac{\widehat{R}_{1}(a, b, c, \lambda)}{\widehat{R}_{2}(a, b, c, \lambda)}\right)$. Thus, there exists $\widehat{h}_{0}(a, b, c, \lambda)=$ $\min \left\{\frac{1}{\lambda},-\frac{\widehat{R}_{1}(a, b, c, \lambda)}{\widehat{R}_{2}(a, b, c, \lambda)}\right\}$, for any $h \in\left(0, \widehat{h}_{0}(a, b, c, \lambda)\right)$, the inequality (6) holds. This completes the proof.

Theorem 2.1 and Theorem 3.1 show that the strong Milstein method (2) and the weak Milstein method (4) can preserve the mean square stability of the system (1).

## 4. Asymptotic stability of the strong Milstein method

In Sections 2 and 3, we consider the mean square stabilities of the two kinds of Milstein methods. This section is concerned with the asymptotic stability. We begin by giving the necessary and sufficient condition for the stochastically asymptotical stability in the large (hereafter, asymptotical stability) of the system (1).

Lemma 4.1 ([3]). The system (1) is stochastically asymptotically stable in the large (hereafter, asymptotically stable) if and only if

$$
\begin{equation*}
a-\frac{1}{2} b^{2}+\lambda \ln |1+c|<0, \tag{7}
\end{equation*}
$$

where $\ln |1+c|=-\infty$, as $c=-1$. Thus when $c=-1$, under the condition (7), the system (1) is asymptotically stable for any $a, b \in \mathbb{R}$.

We say the Milstein method (2) are asymptotically stable for a particular choice of $a, b, c, \lambda$ and $h$ if $\lim _{n \rightarrow \infty}\left|Y_{n}\right|=0$, with probability one, for any $Y_{0}$.

We now establish Lemma 4.2, which will play a key role in the proof of Theorem 4.1. In the similar way as the Theorem 3.2 in [1] we can obtain Lemma 4.2.

Lemma 4.2. Let $\xi$ be a standard Normal random variable. Then, for the integrable function $\varphi: \mathbb{R} \rightarrow \mathbb{R}$, we have

$$
\begin{align*}
& \mathbb{E}\left[\varphi\left(1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}\right)\right] \\
= & \varphi(1)+\varphi^{\prime}(1)\left(\left(a-\frac{1}{2} b^{2}\right) h+\frac{1}{2} b^{2} h\right)+\frac{\varphi^{\prime \prime}(1)}{2} b^{2} h+o(h), \tag{8}
\end{align*}
$$

where $a, b \in \mathbb{R}, h \rightarrow 0$.

Proof. Consider the function $\varphi$, such that there exists $\sigma>0$ and the integral function $\widetilde{\varphi}: \mathbb{R} \rightarrow \mathbb{R}$ satisfying
(1) $\widetilde{\varphi} \equiv \varphi$ on $A=[1-\sigma, 1+\sigma]$,
(2) $\widetilde{\varphi} \in C^{3}(\mathbb{R})$ and $\left|\widetilde{\varphi}^{\prime \prime \prime}(x)\right| \leq H$ for some $H$ and all $x \in \mathbb{R}$,
(3) $\int_{\mathbb{R}}|\varphi-\widetilde{\varphi}| d x \leq K<+\infty$ for some $K$.

The following proof is divided into two parts:
Part 1: We proves formula (8) for $\widetilde{\varphi}\left(1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}\right)$. Using Taylor expansion gives

$$
\widetilde{\varphi}(1+x)=\widetilde{\varphi}(1)+\widetilde{\varphi}^{\prime}(1) x+\frac{\widetilde{\varphi}^{\prime \prime}(1)}{2} x^{2}+\frac{\widetilde{\varphi}^{\prime \prime \prime}(\eta)}{6} x^{3}
$$

where $\eta$ lying between 1 and $1+x$.
Substituting $x=\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}$ and taking expectation lead to

$$
\begin{aligned}
& \mathbb{E}[\widetilde{\varphi}(1+x)] \\
= & \widetilde{\varphi}(1)+\widetilde{\varphi}^{\prime}(1)\left(\left(a-\frac{1}{2} b^{2}\right) h+\frac{1}{2} b^{2} h\right)+\frac{\widetilde{\varphi}^{\prime \prime}(1)}{2}\left(b^{2} h+o(h)\right)+\frac{1}{6} \mathbb{E}\left(\widetilde{\varphi}^{\prime \prime \prime}(\eta) x^{3}\right) .
\end{aligned}
$$

Here the fact has been used that $\mathbb{E} \xi=0, \mathbb{E}|\xi|=2 / \sqrt{2 \pi}, \mathbb{E}|\xi|^{2}=1$. Noticing that $\widetilde{\varphi} \in C^{3}(\mathbb{R})$ and $\left|\widetilde{\varphi}^{\prime \prime \prime}(x)\right| \leq H$ for all $x \in \mathbb{R}$, we have

$$
\left|\frac{1}{6} \mathbb{E}\left(\widetilde{\varphi}^{\prime \prime \prime}(\eta) x^{3}\right)\right| \leq \frac{H}{6} \mathbb{E}\left|x^{3}\right|=o(h),
$$

which implies

$$
\mathbb{E}[\widetilde{\varphi}(1+x)]=\widetilde{\varphi}(1)+\widetilde{\varphi}^{\prime}(1)\left(\left(a-\frac{1}{2} b^{2}\right) h+\frac{1}{2} b^{2} h\right)+\frac{\widetilde{\varphi}^{\prime \prime}(1)}{2} b^{2} h+o(h) .
$$

This proves formula (8) for $\widetilde{\varphi}\left(1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}\right)$.
Part 2: Letting $u_{1}=\frac{1}{2}+\left(a-\frac{1}{2} b^{2}\right) h, u_{2}=b \sqrt{h}$ and, now we estimate the error term $e=\mathbb{E}\left[\varphi\left(u_{1}+\frac{1}{2}\left(u_{2} \xi+1\right)^{2}\right)-\widetilde{\varphi}\left(u_{1}+\frac{1}{2}\left(u_{2} \xi+1\right)^{2}\right)\right]$ as following

$$
|e|=\left|\int_{-\infty}^{+\infty}\left[\varphi\left(u_{1}+\frac{1}{2}\left(u_{2} z+1\right)^{2}\right)-\widetilde{\varphi}\left(u_{1}+\frac{1}{2}\left(u_{2} z+1\right)^{2}\right)\right] p(z) d z\right|,
$$

where $p(z)$ is the density of $\xi$.
In the following, we give the two cases to discuss:
Case 1: $u_{2} z+1 \neq 0$. In this case, we suppose $\sigma \in(0,1 / 2)$. We set $v=$ $u_{1}+\frac{1}{2}\left(u_{2} z+1\right)^{2}$, since $u_{2} z+1 \neq 0$, which means $v-u_{1} \neq 0$. Here without loss of generality, we can assume $b \neq 0$ which implies $u_{2}=b \sqrt{h} \neq 0$. Noting that $A$
is the integration range on which $\varphi(v)-\widetilde{\varphi}(v)=0$, we compute

$$
\begin{align*}
& |e| \leq \int_{\mathbb{R} \backslash A}|\varphi(v)-\widetilde{\varphi}(v)| p\left(\frac{-1+\sqrt{2\left(v-u_{1}\right)}}{u_{2}}\right) \frac{d v}{\sqrt{2\left(v-u_{1}\right)}\left|u_{2}\right|} \\
+ & \int_{\mathbb{R} \backslash A}|\varphi(v)-\widetilde{\varphi}(v)| p\left(\frac{-1-\sqrt{2\left(v-u_{1}\right)}}{u_{2}}\right) \frac{d v}{\sqrt{2\left(v-u_{1}\right)}\left|u_{2}\right|} \\
\leq & \sup _{v \notin A}\left\{p\left(\frac{-1+\sqrt{2\left(v-u_{1}\right)}}{u_{2}}\right) \frac{1}{\sqrt{2\left(v-u_{1}\right)}\left|u_{2}\right|}\right\} \int_{\mathbb{R}}|\varphi(v)-\widetilde{\varphi}(v)| d v \\
+ & \sup _{v \notin A}\left\{p\left(\frac{-1-\sqrt{2\left(v-u_{1}\right)}}{u_{2}}\right) \frac{1}{\sqrt{2\left(v-u_{1}\right)}\left|u_{2}\right|}\right\} \int_{\mathbb{R}}|\varphi(v)-\widetilde{\varphi}(v)| d v  \tag{9}\\
\leq & K\left|u_{2}\right|^{2} \sup _{v \notin A}\left\{p\left(\frac{-1+\sqrt{2\left(v-u_{1}\right)}}{u_{2}}\right) \frac{1}{\sqrt{2\left(v-u_{1}\right)}\left|u_{2}\right|^{3}}\right\} \\
+ & K\left|u_{2}\right|^{2} \sup _{v \notin A}\left\{p\left(\frac{-1-\sqrt{2\left(v-u_{1}\right)}}{u_{2}}\right) \frac{1}{\sqrt{2\left(v-u_{1}\right)}\left|u_{2}\right|^{3}}\right\} \\
= & K b^{2} h \sup _{v \notin A}\left\{p\left(y_{1}\right)\left|y_{1}\right|^{3} \frac{1}{\left|-1+\sqrt{2\left(v-u_{1}\right)}\right|^{3}} \cdot \frac{1}{\sqrt{2\left(v-u_{1}\right)}}\right\} \\
+ & K b^{2} h \sup _{v \notin A}\left\{p\left(y_{2}\right)\left|y_{2}\right|^{3} \frac{1}{\left|-1-\sqrt{2\left(v-u_{1}\right)}\right|^{3}} \cdot \frac{1}{\sqrt{2\left(v-u_{1}\right)}}\right\},
\end{align*}
$$

where $y_{1}=\frac{-1+\sqrt{2\left(v-u_{1}\right)}}{u_{2}}, y_{2}=\frac{-1-\sqrt{2\left(v-u_{1}\right)}}{u_{2}}$.
Noticing that when $v \notin A$, that is to say, $v>1+\sigma$, we have

$$
\begin{aligned}
& -1+\sqrt{2\left(v-u_{1}\right)} \\
= & -1+\sqrt{2\left(v-\frac{1}{2}-\left(a-\frac{1}{2} b^{2}\right) h\right)} \\
> & -1+\sqrt{1+2 \sigma}=\frac{2 \sigma}{1+\sqrt{1+2 \sigma}}, h \rightarrow 0,
\end{aligned}
$$

or $v<1-\sigma$, thus

$$
\begin{equation*}
-1+\sqrt{2\left(v-u_{1}\right)}<-1+\sqrt{1-2 \sigma}=-\frac{2 \sigma}{1+\sqrt{1-2 \sigma}}, h \rightarrow 0 . \tag{11}
\end{equation*}
$$

Noting that $y_{1}=\frac{-1+\sqrt{2\left(v-\frac{1}{2}-\left(a-\frac{1}{2} b^{2}\right) h\right)}}{b \sqrt{h}}$, thus using (10) and (11) gives $y_{1} \rightarrow \infty$ uniformly on $v \notin A$, as $h \rightarrow 0$. By (10) and (11), we find that

$$
\begin{aligned}
& \left(-1+\sqrt{2\left(v-\frac{1}{2}\right)}\right)^{3} \cdot \sqrt{2\left(v-u_{1}\right)} \rightarrow\left(-1+\sqrt{2\left(v-\frac{1}{2}\right)}\right)^{3} \cdot \sqrt{2\left(v-\frac{1}{2}\right)} \\
> & \frac{8 \sigma^{3}}{(1+\sqrt{1+2 \sigma})^{3}} \sqrt{1+2 \sigma}, h \rightarrow 0,
\end{aligned}
$$

or

$$
\left(-1+\sqrt{\left.2\left(v-\frac{1}{2}\right)\right)^{3}} \cdot \sqrt{2\left(v-u_{1}\right)}<\frac{-8 \sigma^{3}}{(1+\sqrt{1-2 \sigma})^{3}} \sqrt{1-2 \sigma}, h \rightarrow 0\right.
$$

which implies $\left|-1+\sqrt{2\left(v-\frac{1}{2}\right)}\right|^{3} \cdot \sqrt{2\left(v-\frac{1}{2}\right)}$ is bounded away from zero. With the properties of $p(z)$, we can prove $p\left(y_{1}\right)\left|y_{1}\right|^{3} \rightarrow 0$, as $y_{1} \rightarrow \infty$, thus

$$
\begin{equation*}
\sup _{v \notin A}\left\{p\left(y_{1}\right)\left|y_{1}\right|^{3} \frac{1}{\left|-1+\sqrt{2\left(v-u_{1}\right)}\right|^{3}} \cdot \frac{1}{\sqrt{2\left(v-u_{1}\right)}}\right\}=o(1), h \rightarrow 0 . \tag{12}
\end{equation*}
$$

Similarly

$$
\begin{equation*}
\sup _{v \notin A}\left\{p\left(y_{2}\right)\left|y_{2}\right|^{3} \frac{1}{\left|-1-\sqrt{2\left(v-u_{1}\right)}\right|^{3}} \cdot \frac{1}{\sqrt{2\left(v-u_{1}\right)}}\right\}=o(1), h \rightarrow 0 . \tag{13}
\end{equation*}
$$

Inserting (12) and (13) into (9) yields

$$
|e|=o(h) .
$$

Case 2: $u_{2} z+1=0$. In this case, we suppose $\sigma \in(1 / 2,1)$. Noticing that $u_{1}=\frac{1}{2}+\left(a-\frac{1}{2} b^{2}\right) h \in A$, when $h \rightarrow 0$, we have

$$
\begin{aligned}
& |e|=\left|\int_{-\infty}^{+\infty}\left[\varphi\left(u_{1}+\frac{1}{2}\left(u_{2} z+1\right)^{2}\right)-\widetilde{\varphi}\left(u_{1}+\frac{1}{2}\left(u_{2} z+1\right)^{2}\right)\right] p(z) d z\right| \\
& =\left|\int_{-\infty}^{+\infty}\left[\varphi\left(u_{1}\right)-\widetilde{\varphi}\left(u_{1}\right)\right] p(z) d z\right|=0 .
\end{aligned}
$$

Combining Part 1 and Part 2 leads the assertion (8).
Based on Lemma 4.1, we can now begin to establish the main theorem in this section.

Theorem 4.1. Given $a, b, c, \lambda$, the system (1) is asymptotically stable if and only if there exists a $h_{1}^{*}>0$, the strong Milstein method (2) is asymptotically stable for all $0<h<h_{1}^{*}$.

Proof. It follows from (2) that

$$
\begin{align*}
& Y_{n+1}=\left[1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}+\frac{1}{2}\left(2 c-c^{2}\right) \Delta N_{n}\right.  \tag{14}\\
& \left.+b c \sqrt{h} \xi \Delta N_{n}+\frac{1}{2} c^{2}\left(\Delta N_{n}\right)^{2}\right] Y_{n}
\end{align*}
$$

where $\xi$ is standard Normal random variable. By the Lemma 5.1 of [5], we deduce immediately that the asymptotic stability of the strong Milstein method (14) is equivalent to

$$
\begin{align*}
& \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}+\frac{1}{2}\left(2 c-c^{2}\right) \Delta N_{n}\right.\right.  \tag{15}\\
+ & \left.\left.b c \sqrt{h} \xi \Delta N_{n}+\frac{1}{2} c^{2}\left(\Delta N_{n}\right)^{2} \right\rvert\,\right]<0
\end{align*}
$$

Multiplying the expected value in (15) by $e^{\lambda h}$ yields

$$
\begin{align*}
& e^{\lambda h} \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}+\frac{1}{2}\left(2 c-c^{2}\right) \Delta N_{n}\right.\right. \\
+ & \left.\left.b c \sqrt{h} \xi \Delta N_{n}+\frac{1}{2} c^{2}\left(\Delta N_{n}\right)^{2} \right\rvert\,\right] \\
= & \mathbb{E}\left[\ln \left|1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}\right|\right]+\lambda h \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h\right.\right. \\
+ & \left.\left.b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}+\frac{1}{2}\left(2 c-c^{2}\right)+b c \sqrt{h} \xi+\frac{1}{2} c^{2} \right\rvert\,\right]  \tag{16}\\
+ & \sum_{k=1}^{\infty} \frac{(\lambda h)^{k}}{k!} \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}+\frac{1}{2}\left(2 c-c^{2}\right) k\right.\right. \\
+ & \left.\left.b c \sqrt{h} \xi k+\frac{1}{2} c^{2} k^{2} \right\rvert\,\right] .
\end{align*}
$$

Using Lemma 4.2 with $\varphi=\ln |\cdot|$ yields

$$
\begin{align*}
& \mathbb{E}\left[\ln \left|1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}\right|\right]  \tag{17}\\
= & \left(a-\frac{1}{2} b^{2}\right) h+\frac{1}{2} b^{2} h-\frac{1}{2} b^{2} h+o(h)=\left(a-\frac{1}{2} b^{2}\right) h+o(h) .
\end{align*}
$$

To discuss the second term at the right side of (16), we consider the following two different cases:
Case 1: $c \neq-1$ then

$$
\begin{aligned}
& \lambda h \mathbb{E}\left[\ln \left|1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}+\frac{1}{2}\left(2 c-c^{2}\right)+b c \sqrt{h} \xi+\frac{1}{2} c^{2}\right|\right] \\
= & \lambda h\left(\ln |1+c|+\mathbb{E}\left[\ln \left|1+\frac{a-\frac{1}{2} b^{2}}{1+c} h+b \sqrt{h} \xi+\frac{b^{2} h \xi^{2}}{2(1+c)}\right|\right]\right) .
\end{aligned}
$$

Recalling the fundamental inequality

$$
\begin{equation*}
\ln |u| \leq|u-1|, u \geq 1 \tag{18}
\end{equation*}
$$

Thus

$$
\begin{aligned}
& \left|\mathbb{E}\left[\ln \left|1+\frac{a-\frac{1}{2} b^{2}}{1+c} h+b \sqrt{h} \xi+\frac{b^{2} h \xi^{2}}{2(1+c)}\right|\right]\right| \\
\leq & \mathbb{E}\left[\ln \left(1+\left|\frac{a-\frac{1}{2} b^{2}}{1+c} h+b \sqrt{h} \xi+\frac{b^{2} h \xi^{2}}{2(1+c)}\right|\right)\right] \\
\leq & \left|\frac{a-\frac{1}{2} b^{2}}{1+c}\right| h+|b| \sqrt{h} \mathbb{E}|\xi|+\frac{b^{2} h}{2|1+c|} \mathbb{E} \xi^{2}=O(\sqrt{h}),
\end{aligned}
$$

which implies

$$
\begin{align*}
& \lambda h \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}\right.\right. \\
+ & \left.\left.\frac{1}{2}\left(2 c-c^{2}\right)+b c \sqrt{h} \xi+\frac{1}{2} c^{2} \right\rvert\,\right]  \tag{19}\\
= & \lambda h \ln |1+c|+o(h) .
\end{align*}
$$

Without loss of generality, we suppose $0<h \leq h_{1}<1$. By using the inequality (18), the third term in the expansion of (16) becomes

$$
\begin{align*}
& \left\lvert\, \sum_{k=2}^{\infty} \frac{(\lambda h)^{k}}{k!} \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}+\frac{1}{2}\left(2 c-c^{2}\right) k\right.\right.\right. \\
&+\left.\left.b c \sqrt{h} \xi k+\frac{1}{2} c^{2} k^{2} \right\rvert\,\right] \mid \\
&<(\lambda h)^{2}\left[\left(\left|a-\frac{1}{2} b^{2}\right| h+\frac{2|b|}{\sqrt{2 \pi}} \sqrt{h}+\frac{1}{2} b^{2} h\right) \sum_{k=2}^{\infty} \frac{\lambda^{k-2}}{k!}\right. \\
&+\left(\frac{1}{2}\left|2 c-c^{2}\right|+\frac{2}{\sqrt{2 \pi}}|b c| \sqrt{h}\right) \sum_{k=2}^{\infty} \lambda^{k-2}(k-1)!  \tag{20}\\
&<\left.\frac{1}{2} c^{2} \sum_{k=2}^{\infty} \frac{\lambda^{k-2} k}{(k-1)!}\right] \\
&<(\lambda h)^{2}\left[\left(\left|a-\frac{1}{2} b^{2}\right| h+\frac{2|b|}{\sqrt{2 \pi}} \sqrt{h}+\frac{1}{2} b^{2} h\right) M_{1}+\left(\frac{1}{2}\left|2 c-c^{2}\right|\right.\right. \\
&+\left.\left.\frac{2}{\sqrt{2 \pi}}|b c| \sqrt{h}\right) M_{2}+\frac{1}{2} c^{2} M_{3}\right]=o(h) .
\end{align*}
$$

where $\sum_{k=2}^{\infty} \frac{\lambda^{k-2}}{k!} \leq M_{1}, \sum_{k=2}^{\infty} \frac{\lambda^{k-2}}{(k-1)!} \leq M_{2}, \sum_{k=2}^{\infty} \frac{\lambda^{k-2} k}{(k-1)!} \leq M_{3}$. Here the fact has been used that $\mathbb{E}|\xi|=2 / \sqrt{2 \pi}, \mathbb{E}|\xi|^{2}=1$. Combining (17), (19), (20) and (16) yields

$$
\begin{align*}
& e^{\lambda h} \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}+\frac{1}{2}\left(2 c-c^{2}\right) \Delta N_{n}\right.\right. \\
+ & \left.\left.b c \sqrt{h} \xi \Delta N_{n}+\frac{1}{2} c^{2}\left(\Delta N_{n}\right)^{2} \right\rvert\,\right]  \tag{21}\\
= & \left(a-\frac{1}{2} b^{2}+\lambda \ln |1+c|\right) h+o(h) .
\end{align*}
$$

From (21), the sign of $\mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}+\frac{1}{2}\left(2 c-c^{2}\right) \Delta N_{n}+\right.\right.$ $\left.\left.b c \sqrt{h} \xi \Delta N_{n}+\frac{1}{2} c^{2}\left(\Delta N_{n}\right)^{2} \right\rvert\,\right]$ is same to the sign of $a-\frac{1}{2} b^{2}+\lambda \ln |1+c|$ when $h$ is sufficiently small, and hence the assertion follows.
Case 2: $c=-1$. In this case, we will show that the Milstein method (2) is also asymptotically stable with sufficiently small stepsize $h$ for all values of $a, b$ and $\lambda>0$. It follows from (16) that

$$
\begin{align*}
& e^{\lambda h} \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}+\frac{1}{2}\left(2 c-c^{2}\right) \Delta N_{n}\right.\right. \\
+ & \left.\left.b c \sqrt{h} \xi \Delta N_{n}+\frac{1}{2} c^{2}\left(\Delta N_{n}\right)^{2} \right\rvert\,\right] \\
= & \mathbb{E}\left[\ln \left|1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}\right|\right]+\lambda h \mathbb{E}\left[\ln \left\lvert\,\left(a-\frac{1}{2} b^{2}\right) h\right.\right.  \tag{22}\\
+ & \left.\left.\frac{1}{2} b^{2} h \xi^{2} \right\rvert\,\right]+\sum_{k=2}^{\infty} \frac{(\lambda h)^{k}}{k!} \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}\right.\right. \\
- & \left.\left.\frac{3 k}{2}-b \sqrt{h} \xi k+\frac{1}{2} k^{2} \right\rvert\,\right] .
\end{align*}
$$

By (18)

$$
\begin{align*}
& \lambda h \mathbb{E}\left[\ln \left|\left(a-\frac{1}{2} b^{2}\right) h+\frac{1}{2} b^{2} h \xi^{2}\right|\right] \\
= & \lambda h \mathbb{E}\left[\ln h+\ln \left|\left(a-\frac{1}{2} b^{2}\right)+\frac{1}{2} b^{2} \xi^{2}\right|\right] \\
\leq & \lambda h\left[\ln h+\mathbb{E}\left(\left|a-\frac{1}{2} b^{2}\right|+\frac{1}{2} b^{2} \xi^{2}\right)\right]  \tag{23}\\
= & \lambda h\left[\ln h+\left|a-\frac{1}{2} b^{2}\right|+\frac{1}{2} b^{2}\right] .
\end{align*}
$$

In a similar way as in deriving (20), we have

$$
\begin{align*}
& \left\lvert\, \sum_{k=2}^{\infty} \frac{(\lambda h)^{k}}{k!} \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}\right.\right.\right.  \tag{24}\\
& \left.\left.-\frac{3 k}{2}-b \sqrt{h \xi} k+\frac{1}{2} k^{2} \right\rvert\,\right] \mid=o(h) .
\end{align*}
$$

Combining (17), (23), (24) with (22) gives

$$
\begin{aligned}
& \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}+\frac{1}{2}\left(2 c-c^{2}\right) \Delta N_{n}\right.\right. \\
+ & \left.\left.b c \sqrt{h} \xi \Delta N_{n}+\frac{1}{2} c^{2}\left(\Delta N_{n}\right)^{2} \right\rvert\,\right] \\
\leq & e^{-\lambda h} \mathbb{E}\left[\ln \left|1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}\right|\right]+\lambda h\left(\ln h+\left|a-\frac{1}{2} b^{2}\right|+\frac{1}{2} b^{2}\right) \\
+ & \sum_{k=2}^{\infty} \frac{(\lambda h)^{k}}{k!} \mathbb{E}\left[\ln \left|1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}-\frac{3 k}{2}-b \sqrt{h} \xi k+\frac{1}{2} k^{2}\right|\right] \\
= & e^{-\lambda h}\left[\left(a-\frac{1}{2} b^{2}\right) h+\lambda h\left(\ln h+\left|a-\frac{1}{2} b^{2}\right|+\frac{1}{2} b^{2}\right)+o(h)\right] .
\end{aligned}
$$

For all sufficiently small $h$, we immediately obtain

$$
\begin{aligned}
& \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \xi+\frac{1}{2} b^{2} h \xi^{2}+\frac{1}{2}\left(2 c-c^{2}\right) \Delta N_{n}\right.\right. \\
+ & \left.\left.b c \sqrt{h} \xi \Delta N_{n}+\frac{1}{2} c^{2}\left(\Delta N_{n}\right)^{2} \right\rvert\,\right]<0
\end{aligned}
$$

Combining these different cases together leads the desired results. The proof is completed.

## 5. Asymptotical stability of the weak Milstein method

In section 3, we have discussed the mean square stability of the weak Milstein method (4). In this section, we will investigate whether the weak Milstein method (4) can share the asymptotical stability of the system (1).

Theorem 5.1. Given $a, b, c, \lambda$, the system (1) is asymptotically stable if and only if there exists a $\widehat{h}_{1}^{*}>0$, the weak Milstein method (4) is asymptotically stable for all $0<h<\widehat{h}_{1}^{*}$.
Proof. In view of Lemma 5.1 of [5], the weak Milstein method (4) is asymptotically stable if and only if

$$
\begin{align*}
& \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \widehat{\xi}+\frac{1}{2} b^{2} h \widehat{\xi}^{2}+\frac{1}{2}\left(2 c-c^{2}\right) \widehat{\Delta N_{n}}\right.\right. \\
+ & \left.\left.b c \sqrt{h} \widehat{\xi} \widehat{\Delta N_{n}}+\frac{1}{2} c^{2}\left(\widehat{\Delta N_{n}}\right)^{2} \right\rvert\,\right]<0 \tag{25}
\end{align*}
$$

where $\mathbb{P}(\widehat{\xi}=-1)=\mathbb{P}(\widehat{\xi}=1)=1 / 2$. Noticing $\widehat{\Delta N_{n}}$ comes from a two point distribution: $\mathbb{P}\left(\widehat{\Delta N_{n}}=0\right)=1-\lambda h, \mathbb{P}\left(\widehat{\Delta N_{n}}=1\right)=\lambda h$, we have

$$
\begin{align*}
& \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \widehat{\xi}+\frac{1}{2} b^{2} h \widehat{\xi}^{2}+\frac{1}{2}\left(2 c-c^{2}\right) \widehat{\Delta N_{n}}\right.\right. \\
&+ b c \sqrt{h} \widehat{\xi} \Delta N_{n}  \tag{26}\\
&=\left.\left.\frac{1}{2} c^{2}\left(\widehat{\Delta N_{n}}\right)^{2} \right\rvert\,\right] \\
&(1-\lambda h) \mathbb{E}\left[\ln \left|1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \widehat{\xi}+\frac{1}{2} b^{2} h \widehat{\xi}^{2}\right|\right]+(\lambda h) \mathbb{E}[\ln \mid 1+ \\
&\left.\left.\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \widehat{\xi}+\frac{1}{2} b^{2} h \widehat{\xi}^{2}+\frac{1}{2}\left(2 c-c^{2}\right)+b c \sqrt{h} \widehat{\xi}+\frac{1}{2} c^{2} \right\rvert\,\right]
\end{align*}
$$

Note that for any $h \in\left(0, \widehat{h}_{1}\right)$ with $\widehat{h}_{1}=\min \left\{1 /\left(\left|a-\frac{1}{2} b^{2}\right|+|b|+\frac{1}{2} b^{2}\right)^{2}, 1\right\}$,

$$
\left|\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \widehat{\xi}+\frac{1}{2} b^{2} h \widehat{\xi}^{2}\right|<\left(\left|a-\frac{1}{2} b^{2}\right|+|b|+\frac{1}{2} b^{2}\right) \sqrt{h}<1 .
$$

Here the fact has been used that $\mathbb{E} \widehat{\xi}=0, \mathbb{E}|\widehat{\xi}|=1, \mathbb{E}|\widehat{\xi}|^{2}=1$. Thus using Taylor expansion gives

$$
\begin{align*}
& \mathbb{E}\left[\ln \left|1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \widehat{\xi}+\frac{1}{2} b^{2} h \widehat{\xi}^{2}\right|\right] \\
& =\mathbb{E}\left(\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \widehat{\xi}+\frac{1}{2} b^{2} h \widehat{\xi}^{2}\right)-\frac{1}{2} \mathbb{E}\left(\left(a-\frac{1}{2} b^{2}\right) h\right.  \tag{27}\\
& \left.+b \sqrt{h} \widehat{\xi}+\frac{1}{2} b^{2} h \widehat{\xi}^{2}\right)^{2}+\frac{1}{3} \mathbb{E}\left[\left(\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \widehat{\xi}+\frac{1}{2} b^{2} h \widehat{\xi}^{2}\right) \varsigma\right]^{3} \\
& =\left(a-\frac{1}{2} b^{2}\right) h+o(h),
\end{align*}
$$

where $0<\varsigma<1$.
To discuss the second term in the expansion of (26), let us discuss the following two possible cases:
Case 1: $c \neq-1$. In this case, similarly to (19), we have

$$
\begin{align*}
& (\lambda h) \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \widehat{\xi}+\frac{1}{2} b^{2} h \widehat{\xi}^{2}+\frac{1}{2}\left(2 c-c^{2}\right)\right.\right. \\
+ & \left.\left.b c \sqrt{h} \widehat{\xi}+\frac{1}{2} c^{2} \right\rvert\,\right] \\
= & \lambda h \ln |1+c|+\lambda h \mathbb{E}\left[\ln \left|1+\frac{\left(a-\frac{1}{2} b^{2}\right) h}{1+c}+b \sqrt{h} \widehat{\xi}+\frac{b^{2} h \widehat{\xi}^{2}}{2(1+c)}\right|\right]  \tag{28}\\
= & \lambda h \ln |1+c|+o(h) .
\end{align*}
$$

Combing (27), (28) and (26) yields

$$
\begin{aligned}
& \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \widehat{\xi}+\frac{1}{2} b^{2} h \widehat{\xi}^{2}+\frac{1}{2}\left(2 c-c^{2}\right) \widehat{\Delta N_{n}}\right.\right. \\
+ & \left.\left.b c \sqrt{h} \widehat{\xi} \widehat{\Delta N_{n}}+\frac{1}{2} c^{2}\left(\widehat{\Delta N_{n}}\right)^{2} \right\rvert\,\right] \\
= & {\left[a-\frac{1}{2} b^{2}+\lambda \ln |1+c|\right] h+o(h) . }
\end{aligned}
$$

With sufficiently small stepsize $h$, we find that (25) holds if and only if

$$
a-\frac{1}{2} b^{2}+\lambda \ln |1+c|<0
$$

Case 2: $\mathrm{c}=-1$. In this case, in a similar way as in deriving (23), we have

$$
\begin{align*}
& \lambda h \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \widehat{\xi}+\frac{1}{2} b^{2} h \widehat{\xi}^{2}+\frac{1}{2}\left(2 c-c^{2}\right)\right.\right. \\
+ & \left.\left.b c \sqrt{h} \widehat{\xi}+\frac{1}{2} c^{2} \right\rvert\,\right] \leq \lambda h\left[\ln h+\left|a-\frac{1}{2} b^{2}\right|+\frac{1}{2} b^{2}\right] . \tag{29}
\end{align*}
$$

Substituting (27), (29) into (26), for sufficiently small stepsize $h$, we have

$$
\begin{aligned}
& \mathbb{E}\left[\ln \left\lvert\, 1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \widehat{\xi}+\frac{1}{2} b^{2} h \widehat{\xi}^{2}+\frac{1}{2}\left(2 c-c^{2}\right) \widehat{\Delta N_{n}}\right.\right. \\
+ & \left.\left.b c \sqrt{h} \widehat{\xi \Delta N_{n}}+\frac{1}{2} c^{2}\left(\widehat{\Delta N_{n}}\right)^{2} \right\rvert\,\right] \\
\leq & (1-\lambda h) \mathbb{E}\left[\ln \left|1+\left(a-\frac{1}{2} b^{2}\right) h+b \sqrt{h} \widehat{\xi}+\frac{1}{2} b^{2} h \widehat{\xi}^{2}\right|\right] \\
+ & \lambda h\left[\ln h+\left|a-\frac{1}{2} b^{2}\right|+\frac{1}{2} b^{2}\right] \\
= & \left(a-\frac{1}{2} b^{2}\right) h+\lambda h\left[\ln h+\left|a-\frac{1}{2} b^{2}\right|+\frac{1}{2} b^{2}\right]+o(h)<0 .
\end{aligned}
$$

Thus, combining these two different cases completes the proof.
Theorem 4.1 and Theorem 5.1 show that the strong Milstein method (2) and the weak Milstein method (4) can well reproduce the asymptotical stability of the system (1).


Figure 1. The mean square stability of the strong Milstein method.

## 6. Numerical experiments

Consider the linear equation with real and scalar coefficients

$$
\begin{equation*}
d x(t)=a x\left(t^{-}\right) d t+b x\left(t^{-}\right) d W(t)+c x\left(t^{-}\right) d N(t), t>0 \tag{30}
\end{equation*}
$$

with $x(0)=1$.
It is know that (30) has the solution[8]

$$
\begin{equation*}
x(t)=(1+c)^{N(t)} \exp \left[\left(a-\frac{1}{2} b^{2}\right) t+b W(t)\right] . \tag{31}
\end{equation*}
$$

Now we test the mean square stability behavior of the Milstein method. In the experiments examining the mean square stability, the expectations are estimated by averaging 2000 different discretized Brownian paths. In the figures, the blue broken lines represent the numerical solutions produced by the Milstein method and we plot the graphs with the vertical axis scaled logarithmically.

In Figures 1, 2, we consider the equation (30) with the parameters $a=-5, b=$ $1, c=-1.1, \lambda=2$ which satisfy the condition (3), hence the system is mean square stable.

In Figure 1, we see that the strong Milstein method (2) is mean square stable on $h=0.2,0.25$. With stepsize increasing to 0.3 , the solution produced by the strong Milstein method grows rapidly to the scale of $10^{12}$ and instability is displayed without doubt. It is shown that the strong Milstein method (2) is unstable if $h$ is large enough and mean square stable if $h \leq 0.25$.


Figure 2. The mean square stability of the weak Milstein method.


Figure 3. The asymptotical stabilities of the stable explicit solution and the strong Milstein method. left: $h=0.01$; middle : $h=0.02$; right: $h=0.2$.

In Figure 2, we draw the numerical solutions produced by the weak Milstein method (4). From Theorem 3.1, we compute $\widehat{h}_{0}(a, b, c, \lambda)=0.3195$. Figure 2 shows that the weak Milstein method (4) is mean square stable on $h=0.2,0.3<$ 0.3195 and unstable on $h=0.4>0.3195$. Hence we conclude that the weak Milstein method (4) is mean square stable when $h<0.3195$.

In order to show the asymptotical stability of the Milstein method, we plot the numerical solutions and the explicit solution (31) by one sample trajectory in all the figures. we use the red real lines and the blue broken lines to represent


Figure 4. The asymptotical stabilities of the unstable explicit solution and the strong Milstein method. left: $h=0.01$; middle : $h=0.02$; right: $h=0.2$.


Figure 5. The asymptotical stabilities of the stable explicit solution and the weak Milstein method. left: $h=0.02$; middle : $h=0.04$; right: $h=0.1$.


Figure 6. The asymptotical stabilities of the unstable explicit solution and the weak Milstein method. left: $h=0.02$; middle : $h=0.04 ;$ right $: h=0.1$.
the explicit solution and the numerical solutions respectively in all the pictures. All the the graphs are drawn with the vertical axis scaled logarithmically.

To verify our result concerning asymptotical stability for the strong Milstein method, we take the coefficients $a=0.5, b=0.1, c=-0.2, \lambda=4$. Note that these parameters satisfy the condition (7), thus the system is asymptotically stable. In Figure 3, we plot the explicit solution (31) and the numerical solutions produced by the strong Milstein method (2) with different stepsize $h$. We replace the parameter $c=-0.2$ by $c=-0.1$, then the the coefficients $a=0.5, b=$ $0.1, c=-0.1, \lambda=4$ don't satisfy the condition (7), so the system becomes
unstable. In this case, we draw the explicit solution (31) and the numerical solutions of the strong Milstein method (2) changing the stepsize $h$ in Figure 4.. Figures 3 and 4 suggest that the two curves match well on $h=0.01,0.02$, whereas, when for some large stepsize $h=0.2$, the two curves can't match well. We conclude that the strong Milstein method (2) can well behave the same stability of the explicit solution (31) with sufficiently small stepsize $h \leq 0.02$.

To examine the asymptotical stability of the weak Milstein method (4), we choose $a=1, b=0.4, c=-0.8, \lambda=4$ which satisfy the condition (7), so the system is asymptotically stable. We change the parameter $a=1$ with $a=20$ which leads to the unstability of the system. In Figure 6, we simulate the explicit solution (31) and the numerical solutions of the weak Milstein method by the different stepsize $h$. In Figures 5 and 6, we find that the two curves match well on $h=0.02,0.04$, whereas, when for some large stepsize $h=0.1$, the two curves can't match well. This implies that the weak Milstein method (4) can well reproduce the same stability of the explicit solution (31) with sufficiently small stepsize $h \leq 0.04$.

## 7. Conclusion

In this paper, we have investigated the mean square stability and the asymptotical stability of the strong Milstein method and the weak Milstein method for the stochastic differential equations with jumps. The forgoing stability results show that both kinds of the Milstein methods can reproduce the mean square stability. Moreover, it is shown that the asymptotical stability for stochastic jump-diffusion differential equations is inherited by the two kinds of Milstein methods with sufficiently small stepsizes $h$.

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