

THE STUDY ON THE EINSTEIN'S CONNECTION IN 5-DIMENSIONAL *ES*-MANIFOLD FOR THE SECOND CLASS

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ABSTRACT. The manifold $*g-ESX_n$ is a generalized n -dimensional Riemannian manifold on which the differential geometric structure is imposed by the unified field tensor $*g^{\lambda\nu}$ through the *ES*-connection which is both Einstein and semi-symmetric. The purpose of the present paper is to derive a new set of powerful recurrence relations and to prove a necessary and sufficient condition for a unique Einstein's connection to exist in 5-dimensional $*g-ESX_5$ and to display a surveyable tensorial representation of 5-dimensional Einstein's connection in terms of the unified field tensor, employing the powerful recurrence relations in the second class.

1. Preliminaries

This section is a brief collection of basic concepts, notations, and results needed in subsequent considerations. They are due to [1],[4] and [5].

(a) n -dimensional $*g$ -unified field theory

Let X_n be an n -dimensional generalized Riemannian manifold referred to a real coordinate system x^ν , which obeys the coordinate transformations $x^\nu \rightarrow x^{\nu'}$ for which

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$$(1.1) \quad \det\left(\frac{\partial x'}{\partial x}\right) \neq 0$$

In $n - g - UFT$ the manifold X_n is endowed with a real nonsymmetric tensor $g_{\lambda\mu}$, which may be decomposed into its symmetric part $h_{\lambda\mu}$ and skew-symmetric part $k_{\lambda\mu}$:

$$(1.2) \quad g_{\lambda\mu} = h_{\lambda\mu} + k_{\lambda\mu}$$

where

$$(1.3) \quad \mathfrak{g} = \det(g_{\lambda\mu}) \neq 0, \quad \mathfrak{h} = \det(h_{\lambda\mu}) \neq 0, \quad \mathfrak{k} = \det(k_{\lambda\mu})$$

In $n - {}^*g - UFT$ the algebraic structure on X_n is imposed by the basic real tensor ${}^*g^{\lambda\nu}$ defined by

$$(1.4) \quad g_{\lambda\mu} {}^*g^{\lambda\nu} = g_{\mu\lambda} {}^*g^{\nu\lambda} = \delta_\mu^\nu$$

It may be also decomposed into its symmetric part ${}^*h^{\lambda\nu}$ and skew-symmetric part ${}^*k^{\lambda\nu}$:

$$(1.5) \quad {}^*g^{\lambda\nu} = {}^*h^{\lambda\nu} + {}^*k^{\lambda\nu}$$

Since $\det({}^*h^{\lambda\nu}) \neq 0$, we may define a unique tensor ${}^*h_{\lambda\mu}$ by

$$(1.6) \quad {}^*h_{\lambda\mu} {}^*h^{\lambda\nu} = \delta_\mu^\nu$$

In $n - {}^*g - UFT$ we use both ${}^*h^{\lambda\nu}$ and ${}^*h_{\lambda\mu}$ as tensors for raising and/or lowering indices of all tensors in X_n in the usual manner. We then have

$$(1.7) \quad {}^*k_{\lambda\mu} = {}^*k^{\rho\sigma} {}^*h_{\lambda\rho} {}^*h_{\mu\sigma}, \quad {}^*g_{\lambda\mu} = {}^*g^{\rho\sigma} {}^*h_{\lambda\rho} {}^*h_{\mu\sigma}$$

so that

$$(1.8) \quad {}^*g_{\lambda\mu} = {}^*h_{\lambda\mu} + {}^*k_{\lambda\mu}$$

The differential geometric structure on X_n is imposed by the tensor ${}^*g^{\lambda\nu}$ by means of a connection $\Gamma_{\lambda}^{\nu}{}_{\mu}$ defined by a system of equations

$$(1.9) \quad D_\omega {}^*g^{\lambda\nu} = -2S_{\omega\alpha}{}^\nu {}^*g^{\lambda\alpha}$$

where D_ω denotes the symbol of the covariant derivative with respect to $\Gamma_{\lambda}^{\nu}{}_{\mu}$ and $S_{\lambda\mu}{}^\nu$ is the torsion tensor of $\Gamma_{\lambda}^{\nu}{}_{\mu}$. Under certain conditions the system (1.9) admits a unique solutions $\Gamma_{\lambda}^{\nu}{}_{\mu}$.

It has been shown in [4] that if the system (1.9) admits $\Gamma_{\lambda}^{\nu}{}_{\mu}$, it must be of the form

$$(1.10) \quad \Gamma_{\lambda}^{\nu}{}_{\mu} = * \left\{ \begin{array}{c} \nu \\ \lambda\mu \end{array} \right\} + U^{\nu}{}_{\lambda\mu} + S_{\lambda\mu}{}^{\nu},$$

where

$$(1.11) \quad U_{\nu\lambda\mu} = S_{(\lambda\mu)\nu}^{100} + 2 S_{\nu(\lambda\mu)}^{(10)0}.$$

(b) Some notations and results

The following quantities are frequently used in our further considerations:

$$(1.12) \quad *g = \det(*g_{\lambda\mu}), \quad *h = \det(*h_{\lambda\mu}), \quad *k = \det(*k_{\lambda\mu})$$

$$(1.13) \quad *g = \frac{*g}{*h}, \quad *k = \frac{*k}{*h}.$$

$$(1.14) \quad K_p = *k_{[\alpha_1}{}^{\alpha_1} *k_{\alpha_2}{}^{\alpha_2} \dots *k_{\alpha_p]}{}^{\alpha_p}, \quad (p = 0, 1, 2, \dots).$$

$$(1.15) \quad {}^{(0)}*k_{\lambda}{}^{\nu} = \delta_{\lambda}^{\nu}, \quad {}^{(p)}*k_{\lambda}{}^{\nu} = *k_{\lambda}{}^{\alpha} {}^{(p-1)}*k_{\alpha}{}^{\nu} \quad (p = 1, 2, \dots).$$

$$(1.16) \quad K_{\omega\mu\nu} = \nabla_{\nu} *k_{\omega\mu} + \nabla_{\omega} *k_{\nu\mu} + \nabla_{\mu} *k_{\omega\nu}$$

where ∇_{ω} is the symbolic vector of the covariant derivative with respect to the christoffel symbols $* \left\{ \begin{array}{c} \nu \\ \lambda\mu \end{array} \right\}$ defined by $*h_{\lambda\mu}$ in the usual way.

In X_n it was proved in [4] that

$$(1.17) \quad K_0 = 1, \quad K_n = *k \text{ if } n \text{ is even, and } K_n = 0 \text{ if } n \text{ is odd.}$$

$$(1.18) \quad *g = 1 + K_2 + \dots + K_{n-\sigma}.$$

$$(1.19) \quad \sum_{s=0}^{n-\sigma} K_s {}^{(n-s)}*k_{\lambda}{}^{\nu} = 0 \quad (p = 0, 1, 2, \dots).$$

We also use the following useful abbreviations, denoting an arbitrary tensor $T_{\omega\mu\nu}$ skew-symmetric in the first two indices by T :

$$(1.20) \quad \overset{pqr}{T} = \overset{pqr}{T}_{\omega\mu\lambda} = T_{\alpha\beta\gamma} {}^{(p)}*k_{\omega}{}^{\alpha} {}^{(q)}*k_{\mu}{}^{\beta} {}^{(r)}*k_{\lambda}{}^{\gamma}$$

and for an arbitrary tensor T_{\dots} for $p = 1, 2, 3, \dots$:

$$(1.21) \quad {}^{(p)}T_{\dots}^{\nu\dots} = {}^{(p-1)}*k_{\alpha}{}^{\nu} T_{\dots}^{\alpha\dots}.$$

DEFINITION 1.1. The tensors $*g_{\lambda\mu}$ is said to be

(1) of the first class, if $K_{n-\sigma} \neq 0$

(2) of the second class with j th category ($j \geq 1$), if

$$(1.22) \quad K_{2j} \neq 0, \quad K_{2j+2} = K_{2j+4} = \cdots = K_{n-\sigma} = 0$$

(3) of the third class, if $K_2 = K_4 = \cdots = K_{n-\sigma} = 0$

In 5 - $*g$ - UFT , there are three classes: namely the first class when $K_4 \neq 0$, the second class when $K_4 = 0, K_2 \neq 0$ and the third class when $K_2 = K_4 = 0$.

It is well known that the basic scalars M are solutions of the characteristic equation

$$(1.23) M^\sigma (M^{n-\sigma} + K_2 M^{n-2-\sigma} + \cdots + K_{n-2-\sigma} M^2 + K_{n-\sigma}) = 0$$

On the other hand, it has shown in [5] that the tensor $S_{\lambda\mu}{}^\nu$ satisfies

$$(1.24) \quad S = B - 3 \overset{(110)}{S}$$

where

$$(1.25) \quad 2B_{\omega\mu\nu} = K_{\omega\mu\nu} + 3K_{\alpha[\mu\beta} * k_\omega]^\alpha * k_\nu^\beta$$

In our subsequent chapter, we start with the relation (1.24) to solve the system (1.9). Furthermore, for the second class, the nonholonomic solution of (1.24) is given by

$$(1.26) \quad M S_{xyz} = B_{xyz}$$

or equivalently

$$(1.27) \quad 4M S_{xyz} = (2 + \underset{z}{M} \underset{x}{M} + \underset{z}{M} \underset{y}{M}) K_{xyz} + \underset{z}{M} (\underset{x}{M} + \underset{z}{M}) K_{zxy} + \underset{z}{M} (\underset{y}{M} + \underset{z}{M}) K_{yzx}$$

where

$$(1.28) \quad M = 1 + \underset{xyz}{M} \underset{x}{M} \underset{y}{M} + \underset{y}{M} \underset{z}{M} \underset{z}{M} \underset{x}{M}$$

Therefore, in virtue of (1.26), we see that a necessary and sufficient condition for the system (1.9) to have a unique solution in the second class is

$$(1.29) \quad \underset{xyz}{M} \neq 0 \text{ for all } x, y, z.$$

DEFINITION 1.2. A connection Γ_{λ}^{ν} is said to be *semi-symmetric* if its torsion tensor $S_{\lambda\mu}^{\nu}$ is of the form

$$(1.30) \quad S_{\lambda\mu}^{\nu} = 2\delta_{[\lambda}^{\nu} X_{\mu]}.$$

for an arbitrary non-null vector X_{μ} .

A connection which is both semi-symmetric and Einstein is called an ES connection. An n -dimensional generalized Riemannian manifold X_n , on which the differential geometric structure is imposed by $*g^{\lambda\nu}$ by means of an ES connection, is called an n -dimensional $*g-ES$ manifold. We denote this manifold by $*g-ESX_n$ in our further considerations.

THEOREM 1.3. *In the second class, the basic scalars in $*g-ESX_5$ are given by*

$$(1.31) \quad M_1 = -M_2 = \sqrt{-K_2} \neq 0, \quad M_3 = M_4 = M_5 = 0$$

Proof. For the second class of $5-*g-UFT$, the characteristic equation (1.23) is reduced to

$$(1.32) \quad M(M^4 + K_2M^2) = 0$$

from which our assertion follows. \square

THEOREM 1.4. *The main recurrence relation in the second class is*

$$(1.33) \quad {}^{(p+3)*}k_{\lambda}^{\nu} = -K_2 {}^{(p+1)*}k_{\lambda}^{\nu} \quad (p = 0, 1, 2, \dots)$$

Proof. When $*g_{\lambda\mu}$ belongs to the second class, the characteristic equation (1.23) is reduced to

$$(1.34) \quad \sum_{f=0}^2 K_f M^{n-f} = M^{n-2} \sum_{f=0}^2 K_f M^{2-f}$$

If M is a root of (1.34), it satisfies

$$(1.35) \quad 0 = M_x \sum_{f=0}^2 K_f M_x^{2-f} = \sum_{f=0}^2 K_f M_x^{2-f+1}$$

Multiplying δ_x^i to both sides of (1.35) and making use of (1.15), we have (1.33). \square

The following theorem is a simple consequence of (1.31).

THEOREM 1.5. *In the second class, the basic scalars M_x satisfy*

$$(1.36) \quad M_1 + M_2 = M_x + M_y = 0 \quad (x, y = 3, 4, 5)$$

$$(1.37) \quad M_1 M_2 = K_2, \quad M_1 M_x = M_2 M_x = M_x M_y = 0 \quad (x, y = 3, 4, 5)$$

In virtue of the above theorem, we have

THEOREM 1.6. *In the second class, the following identities hold for all values of x and y when $x \neq y$*

$$(1.38) \quad M_x^2 M_y^2 = K_2 M_x M_y$$

$$(1.39) \quad M_x^{(2)} M_y^{(1)} = 0$$

THEOREM 1.7. *If $T_{\omega\mu\nu}$ is a tensor skew-symmetric in the first two indices, then the following recurrence relations hold in the second class of $*g - ESX_5$:*

$$(1.40) \quad T^{(21)r} = 0$$

$$(1.41) \quad T^{22r} = K_2 T^{11r}$$

Proof. The relations (1.40) and (1.41) follow from (1.38) and (1.39). For example, the relation (1.41) is obtained as in the following way:

$$\begin{aligned} T^{22r} &= \sum_{x,y,z} T_{xyz} M_x^2 M_y^2 M_z^r A_\omega^x A_\mu^y A_\nu^z \\ &= \sum_{x,y,z} T_{xyz} (K_2 M_x M_y M_z^r) A_\omega^x A_\mu^y A_\nu^z \\ &= K_2 T^{11r} \end{aligned}$$

$$(1.42)$$

□

2. Einstein's connection $\Gamma_{\lambda}^{\nu}{}_{\mu}$ in the second class

In this section, we shall derive surveyable representations of $\Gamma_{\lambda}^{\nu}{}_{\mu}$ in terms of $*g^{\lambda\nu}$, employing the recurrence relations.

In the following theorem, we shall prove two relations in X_n . These relations will be used in our subsequent theorem when we are concerned with the solution of (1.9).

THEOREM 2.1. We have

$$(2.1) \quad B^{(pq)r} = S^{(pq)r} + S^{(p'q')r} + S^{(p'q)r'} + S^{(pq')r'}$$

$$(2.2) \quad \begin{aligned} 2 B_{\omega\mu\nu}^{(pq)r} &= K_{\omega\mu\nu}^{(pq)r} + K_{\nu[\omega\mu]}^{r''(pq)} + \\ &+ \frac{1}{2} (K_{\omega\mu\nu}^{(pq')r'} + K_{\omega\mu\nu}^{(p'q)r'} + K_{\nu[\omega\mu]}^{r'p'q} + K_{\nu[\omega\mu]}^{r'q'p}) \end{aligned}$$

where

$$(2.3) \quad p' = p + 1, \quad q' = q + 1, \quad r' = r + 1, \quad r'' = r + 2$$

Proof. In virtue of (1.24) and (1.20), the first relation (2.1) is obtained as in the following way:

$$(2.4) \quad \begin{aligned} B^{(pq)r} &= B_{\omega\mu\nu}^{(pq)r} = \frac{1}{2} B_{\omega\beta\gamma} ((p)^* k_{\omega}^{\alpha} (q)^* k_{\mu}^{\beta} + (q)^* k_{\omega}^{\alpha} (p)^* k_{\mu}^{\beta}) (r)^* k_{\nu}^{\gamma} \\ &= \frac{1}{2} (S_{\alpha\beta\gamma} + S_{\epsilon\eta\gamma} {}^* k_{\alpha}^{\epsilon} {}^* k_{\beta}^{\eta} + S_{\epsilon\beta\eta} {}^* k_{\alpha}^{\epsilon} {}^* k_{\gamma}^{\eta} + S_{\alpha\epsilon\eta} {}^* k_{\beta}^{\epsilon} {}^* k_{\gamma}^{\eta}) \times \\ &\quad \times ((p)^* k_{\omega}^{\alpha} (q)^* k_{\mu}^{\beta} + (q)^* k_{\omega}^{\alpha} (p)^* k_{\mu}^{\beta}) (r)^* k_{\nu}^{\gamma} \end{aligned}$$

After a lengthy calculation, we note that the right-hand side of the above equation is equal to that of (2.1). Similarly, we verify (2.2) using (1.20) and (1.25). \square

THEOREM 2.2. A necessary and sufficient condition for the system (1.9) to admit a unique solution $\Gamma_{\lambda}^{\nu}{}_{\mu}$ in the second class is that

$$(2.5) \quad 1 - (K_2)^2 \neq 0$$

Proof. In virtue of (1.31), the symmetric scalars M_{xyz} defined by (1.28) takes values as in the following 3 cases:

If two of the indices x, y, z are 1, 2 or 3, 4, then

$$(2.6) \quad M_{xyz} = 1 + K_2, \quad 1$$

If at least one of x, y, z is 5 and no two take the values 1, 2 nor 3, 4, then

$$(2.7) \quad M_{xyz} = 1$$

In the remaining cases,

$$(2.8) \quad M_{xyz} = 1 - K_2, \quad 1$$

It is easily verified that the product of two factors in the right of (2.6) is $1 + K_2$, that of one factor in the right of (2.7) is 1, and that of two factors in the right of (2.8) is $1 - K_2$. Hence we have proved our assertion (2.5) in virtue of (1.29). \square

THEOREM 2.3. *The system of equations (1.24) in the second class is reduced to the following 5 equations:*

$$(2.9) \quad \begin{cases} B = S + 2 \begin{smallmatrix} (10)1 \\ S \end{smallmatrix} + \begin{smallmatrix} 110 \\ S \end{smallmatrix} \\ B = \begin{smallmatrix} (10)1 \\ S \end{smallmatrix} + \begin{smallmatrix} (10)1 \\ S \end{smallmatrix} + \begin{smallmatrix} (20)2 \\ S \end{smallmatrix} + \begin{smallmatrix} 112 \\ S \end{smallmatrix} \\ B = (K_2)^2 \begin{smallmatrix} (20)2 \\ S \end{smallmatrix} + \begin{smallmatrix} (10)1 \\ S \end{smallmatrix} + \begin{smallmatrix} (20)2 \\ S \end{smallmatrix} - K_2 \begin{smallmatrix} 112 \\ S \end{smallmatrix} \\ B = (1 + K_2) \begin{smallmatrix} 110 \\ S \end{smallmatrix} \\ B = (1 + K_2) \begin{smallmatrix} 112 \\ S \end{smallmatrix} \end{cases}$$

Proof. This assertion follows from (2.1) using (1.40), (1.41) and (1.33). \square

THEOREM 2.4. *If the conditions (2.5) is satisfied, the unique solution of (1.24) is given by*

$$(2.10) (1 - K_2^2)(S - B) = -2 \begin{smallmatrix} (10)1 \\ B \end{smallmatrix} + (K_2 - 1) \begin{smallmatrix} 110 \\ B \end{smallmatrix} + 2 \begin{smallmatrix} (20)2 \\ B \end{smallmatrix} + 2 \begin{smallmatrix} 112 \\ B \end{smallmatrix}$$

Proof. (2.10) is a solution of (2.9). \square

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