# WAVELET CHARACTERIZATIONS OF VARIABLE HARDY-LORENTZ SPACES 

Yao He


#### Abstract

In this paper, let $q \in(0,1]$. We establish the boundedness of intrinsic $g$-functions from the Hardy-Lorentz spaces with variable exponent $H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ into Lorentz spaces with variable exponent $L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$. Then, for any $q \in(0,1]$, via some estimates on a discrete LittlewoodPaley $g$-function and a Peetre-type maximal function, we obtain several equivalent characterizations of $H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ in terms of wavelets.


## 1. Introduction

As a generalization of $L^{p}\left(\mathbb{R}^{n}\right)$, the variable Lebesgue space $L^{p(\cdot)}\left(\mathbb{R}^{n}\right)$ was introduced by Orlicz [22] in 1930's. Lorentz spaces on $\mathbb{R}^{n}$ were studied by Lorentz in the early 1950 's. Lorentz spaces, as generalizations of $L^{p}\left(\mathbb{R}^{n}\right)$, are known to be the intermediate spaces of Lebesgue spaces in the real interpolation method; see $[1,18]$. Over the past couple of years, the study of Hardy-Lorentz spaces has always been an interesting topic. For example, the real interpolation of the Hardy-Lorentz space $H^{p, q}\left(\mathbb{R}^{n}\right)$ was investigated by Fefferman, Riviére, and Sagher [4]; the space $H^{1, \infty}\left(\mathbb{R}^{n}\right)$ was considered by Fefferman and Soria [5].

Nowadays, due to the development of variable Lebesgue spaces, there has been a lot of research on the study of Hardy spaces with variable exponents in harmonic analysis. A major breakthrough on Lebesgue spaces with variable exponent is that under some regularity assumptions on $p(\cdot)$, the Hardy-Littlewood maximal operator is bounded on $L^{p(\cdot)}\left(\mathbb{R}^{n}\right)$ [3]. Moreover, Nakai and Sawano [21] made a lot of progress on variable Hardy spaces $H^{p(\cdot)}\left(\mathbb{R}^{n}\right)$. They established the atomic decompositions and the dual spaces of $H^{p(\cdot)}\left(\mathbb{R}^{n}\right)$ in [3]. Later, Sawano [23] extended the atomic characterization of $H^{p(\cdot)}\left(\mathbb{R}^{n}\right)$ and improved the corresponding results in [21]. Recently, Jiao et al. [14] established some realvariable characterizations of variable Hardy-Lorentz spaces. As applications of

[^0]the atomic decompositions, they developed a theory of real interpolation and formulated the dual space of the variable Hardy-Lorentz space with $0<p_{-} \leq$ $p_{+} \leq 1$ and $0<q<\infty$.

In the 1990s, the wavelet theory was established involving different Hardytype spaces. Precisely, several equivalent wavelet characterizations of $H^{1}\left(\mathbb{R}^{n}\right)$ were established by Meyer [20]; some equivalent wavelet characterizations of the weak Hardy space $H^{1, \infty}\left(\mathbb{R}^{n}\right)$ were studied by Liu [19]; a wavelet area integral characterization of the weighted Hardy space $H_{\omega}^{p}\left(\mathbb{R}^{n}\right)$ for any $p \in(0,1]$ was established by Wu [25]; and independently, via the vector-valued CalderónZygmund theory, a characterization of $H_{\omega}^{p}\left(\mathbb{R}^{n}\right)$ for $p \in(0,1]$ in terms of wavelets without compact supports was established by García-Cuerva and Martell [10]. Later, the wavelet inequalities of Lebesgue spaces with variable exponents were introduced by Kopaliani [16] and Izuki [12] independently. In addition, the wavelet characterization for weighted Lebesgue spaces with variable exponents was established by Izuki, Nakai, and Sawano [13].

Recently, via wavelets, several equivalent characterizations of the MusielakOrlicz Hardy space $H^{\varphi}\left(\mathbb{R}^{n}\right)$ were established by Fu and Yang [8]. Later, via wavelets, several equivalent characterizations of $H^{p(\cdot)}\left(\mathbb{R}^{n}\right)$ were established by $\mathrm{Fu}[7]$, which extends the wavelet characterizations of the classical Hardy space in $[20$, Theorems 5.1, 6.4]. In addition, when $(\mathcal{X}, d, \mu)$ is a metric measure space of homogeneous type in the sense of R. R. Coifman and G. Weiss and $H_{a t}^{1}(\mathcal{X})$ is the atomic Hardy space, Fu and Yang [9] established several equivalent characterizations of $H_{a t}^{1}(\mathcal{X})$ in terms of wavelets.

Motivated by the above results, especially by [8,14], we establish several equivalent characterizations of $H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ in terms of wavelets where $q \in(0,1]$.

We describe how we organize this paper. In Section 2, we first recall some known notions and notation. Then, recall the atomic characterizations of $H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ from [14, Theorem 5.4] (see Lemma 2.12 below). In Section 3, for any $q \in(0,1]$, we establish the boundedness of intrinsic $g$-functions from the Hardy-Lorentz spaces with variable exponent $H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ into Lorentz spaces with variable exponent $L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ (see Theorem 3.3 below), and get some estimates on a discrete Littlewood-Paley $g$-function and a Peetre-type maximal function (see Propositions 3.5 and 3.6 , respectively, below). In Section 4, we prove Theorem 4.1. Via the estimate on the Peetre-type maximal function, the wavelet characterizations of Lebesgue spaces from [20] and some standard arguments on the wavelet characterizations of the classical Hardy spaces, we complete the proof of Theorem 4.1.
Notation. In this paper, we denote by $C$ a positive constant which is independent of the main parameters, but it may vary from line to line. We also use $C_{(\alpha, \beta, \ldots)}$ to denote a positive constant depending on the parameters $\alpha, \beta, \ldots$. The symbol $f \lesssim g$ means $f \leq C g$ for a positive constant $C$, and $f \sim g$ amounts to $f \gtrsim g \gtrsim f$. For any $a \in \mathbb{R}$, the symbol $\lfloor a\rfloor$ denotes the largest integer $m$ such that $m \leq a$. Let $\mathbb{N}:=\{1,2, \ldots\}, \mathbb{Z}_{+}:=\mathbb{N} \cup\{0\}$ and $\mathbb{Z}_{+}^{n}:=\left(\mathbb{Z}_{+}\right)^{n}$. For any
$p \in[1, \infty], p^{\prime}$ denotes its conjugate number, namely, $1 / p+1 / p^{\prime}=1$. For any subset $E$ of $\mathbb{R}^{n}$, we use $\chi_{E}$ to denote its characteristic function. Moreover, $\langle\cdot, \cdot \cdot\rangle$ and $(\cdot, \cdot)_{L^{2}\left(\mathbb{R}^{n}\right)}$ represent the duality relation and the $L^{2}\left(\mathbb{R}^{n}\right)$ inner product, respectively.

## 2. Preliminaries

In this section, we first recall some notions and notation. For any $x \in \mathbb{R}^{n}$ and $r \in(0, \infty)$, let

$$
B(x, r):=\left\{y \in \mathbb{R}^{n}:|y-x|<r\right\}
$$

denote the open ball. Let $\mathcal{A}\left(\mathbb{R}^{n}\right)$ be the set of all Lebesgue measurable functions on $\mathbb{R}^{n}$.

A measurable function $p(\cdot): \mathbb{R}^{n} \rightarrow(0, \infty)$ is called a variable exponent. Denote by $\mathcal{P}\left(\mathbb{R}^{n}\right)$ the collection of all variable exponents $p(\cdot)$ satisfying

$$
0<p_{-}:=\operatorname{ess} \inf _{x \in \mathbb{R}^{n}} p(x) \leq \operatorname{ess} \sup _{x \in \mathbb{R}^{n}} p(x)=: p_{+}<\infty
$$

In the following, let

$$
\underline{p}=\min \left\{p_{-}, 1\right\} .
$$

Definition $2.1\left(\left[2\right.\right.$, Definition 2.16]). Let $p(\cdot) \in \mathcal{P}\left(\mathbb{R}^{n}\right)$. Then the Lebesgue space with variable exponent is defined by setting

$$
L^{p(\cdot)}\left(\mathbb{R}^{n}\right):=\left\{f \in \mathcal{A}\left(\mathbb{R}^{n}\right):\|f\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}<\infty\right\}
$$

where
$\|f\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}:=\inf \left\{\lambda \in(0, \infty): \rho_{p(\cdot)}\left(\frac{|f|}{\lambda}\right) \leqslant 1\right\}, \rho_{p(\cdot)}(f):=\int_{\mathbb{R}^{n}}|f(x)|^{p(x)} d x$.
Definition 2.2 ([15, Definition 2.2]). Let $p(\cdot) \in \mathcal{P}\left(\mathbb{R}^{n}\right)$ and let $0<q \leq \infty$. Then the Lorentz space with variable exponent is defined by setting

$$
L^{p(\cdot), q}\left(\mathbb{R}^{n}\right):=\left\{f \in \mathcal{A}\left(\mathbb{R}^{n}\right):\|f\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}<\infty\right\}
$$

where

$$
\|f\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}:= \begin{cases}\left(\int_{0}^{\infty} \lambda^{q}\left\|\chi_{\left\{x \in \mathbb{R}^{n}:|f(x)|>\lambda\right\}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}^{q} \frac{d \lambda}{\lambda}\right)^{1 / q}, & \text { if } 0<q<\infty \\ \sup _{\lambda>0} \lambda\left\|\chi_{\left\{x \in \mathbb{R}^{n}:|f(x)|>\lambda\right\}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}, & \text { if } q=\infty .\end{cases}
$$

Lemma 2.3 ([14, Lemma 2.8]). Let $p(\cdot) \in \mathcal{P}\left(\mathbb{R}^{n}\right)$ and let $0<q \leq \infty$. Then, for all $f \in L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ and $s \in(0, \infty)$, it holds true that

$$
\left\||f|^{s}\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}=\|f\|_{L^{s p(\cdot), s q}\left(\mathbb{R}^{n}\right)}^{s}
$$

A function $p(\cdot) \in \mathcal{P}\left(\mathbb{R}^{n}\right)$ is said to satisfy the globally log-Hölder continuous condition, denoted by $p(\cdot) \in C^{\log }\left(\mathbb{R}^{n}\right)$, if there exist positive constants $C_{p(\cdot)}$, $C_{\infty}$ and $p_{\infty}$ such that, for all $x, y \in \mathbb{R}^{n}$,

$$
\begin{equation*}
|p(x)-p(y)| \leq \frac{C_{p(\cdot)}}{\log (e+1 /|x-y|)} \tag{2.1}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|p(x)-p_{\infty}\right| \leq \frac{C_{\infty}}{\log (e+|x|)} \tag{2.2}
\end{equation*}
$$

For $r \in(0, \infty)$, we denote $L_{\text {loc }}^{r}\left(\mathbb{R}^{n}\right)$ to be the set of all $r$-locally integrable functions on $\mathbb{R}^{n}$. Recall that the Hardy-Littlewood maximal operator $\mathcal{M}$ is defined by setting, for all $f \in L_{\text {loc }}^{1}\left(\mathbb{R}^{n}\right)$,

$$
\begin{equation*}
\mathcal{M}(f)(x):=\sup _{x \in B} \frac{1}{|B|} \int_{B}|f(y)| d y, \quad x \in \mathbb{R}^{n} \tag{2.3}
\end{equation*}
$$

where the supremum is taken over all balls $B$ of $\mathbb{R}^{n}$ containing $x$.
We denote $S\left(\mathbb{R}^{n}\right)$ to be the space of all Schwartz functions and $S^{\prime}\left(\mathbb{R}^{n}\right)$ to be its topological dual space equipped with the weak-* topology. For $N \in \mathbb{N}$, let

$$
\mathcal{F}_{N}\left(\mathbb{R}^{n}\right):=\left\{\psi \in \mathcal{S}\left(\mathbb{R}^{n}\right): \sum_{\beta \in \mathbb{Z}_{+}^{n},|\beta| \leq N} \sup _{x \in \mathbb{R}^{n}}\left[(1+|x|)^{N}\left|D^{\beta} \psi(x)\right|\right] \leq 1\right\}
$$

where, for any $\beta:=\left(\beta_{1}, \ldots, \beta_{n}\right) \in \mathbb{Z}_{+}^{n},|\beta|=\beta_{1}+\cdots+\beta_{n}$ and $D^{\beta}:=$ $\left(\frac{\partial}{\partial x_{1}}\right)^{\beta_{1}} \cdots\left(\frac{\partial}{\partial x_{n}}\right)^{\beta_{n}}$. Then for all $f \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$, the radial grand maximal function $f_{N,+}^{*}$ of $f$ is defined by setting, for all $x \in \mathbb{R}^{n}$,

$$
f_{N,+}^{*}(x):=\sup \left\{\left|f * \psi_{t}(x)\right|: t \in(0, \infty) \text { and } \psi \in \mathcal{F}_{N}\left(\mathbb{R}^{n}\right)\right\},
$$

where, for all $t \in(0, \infty)$ and $\xi \in \mathbb{R}^{n}, \psi_{t}(\xi):=t^{-n} \psi(\xi / t)$. We simply use $f^{*}$ to denote $f_{N,+}^{*}$.
Definition 2.4 ([14, Definition 2.14]). Let $p(\cdot) \in \mathcal{P}\left(\mathbb{R}^{n}\right)$, let $0<q \leq \infty$ and let $N \in\left(\frac{n}{\underline{p}}+n+1, \infty\right)$ be a positive integer. The variable Hardy-Lorentz space $H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ is defined by setting

$$
H^{p(\cdot), q}\left(\mathbb{R}^{n}\right):=\left\{f \in S^{\prime}\left(\mathbb{R}^{n}\right):\left\|f^{*}\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}<\infty\right\}
$$

equipped with the quasi-norm

$$
\|f\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}:=\left\|f^{*}\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} .
$$

We denote $\mathbb{P}^{s}\left(\mathbb{R}^{n}\right)$ to be the set of all polynomials having degree at most $s$. For a locally integrable function $f$, a ball $B$ and a nonnegative integer $s$, there exists a unique polynomial $P$ such that for any polynomial $R \in \mathbb{P}^{s}\left(\mathbb{R}^{n}\right)$,

$$
\int_{B}(f(x)-P(x)) R(x) d x=0
$$

Denote this unique polynomial $P$ by $P_{B}^{s} f$.
Definition $2.5([21])$. Let $p(\cdot) \in \mathcal{P}\left(\mathbb{R}^{n}\right)$ and let $1 \leq r<\infty$. The BMO space $B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)$ is defined by setting

$$
B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right):=\left\{f \in L_{\mathrm{loc}}^{r}\left(\mathbb{R}^{n}\right):\|f\|_{B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)}<\infty\right\}
$$

where

$$
\begin{aligned}
\|f\|_{B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)}:= & \sup _{B \in \mathcal{B}} \frac{|B|}{\left\|\chi_{B}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}}\left(\frac{1}{|B|} \int_{B}\left|f(x)-P_{B}^{s} f(x)\right|^{r} d x\right)^{1 / r}, \\
& (1 \leq r<\infty)
\end{aligned}
$$

where $\mathcal{B}$ is the set of all balls.
Remark 2.6. Let $p(\cdot) \in C^{\log }\left(\mathbb{R}^{n}\right)$, let $d$ be an integer satisfying

$$
\begin{equation*}
d \geqslant d_{p(\cdot)}:=\min \left\{d \in \mathbb{Z}_{+}: p_{-}(n+1+d)>n\right\}, \tag{2.4}
\end{equation*}
$$

and let $r \in\left(\max \left\{p_{+}, 1\right\}, \infty\right]$ and $q \in(0,1]$. By [14, Theorems 5.4 and 7.2], we find that
(i) $\left(H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)\right)^{*}=B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)$;
(ii) $f \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ if and only if $f \in\left(B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)\right)^{*}$ and $f^{*} \in L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$.

Indeed, for any $q \in(0,1]$, if $f \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$, then, by (i), we deduce that

$$
f \in\left(H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)\right)^{* *}=\left(B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)\right)^{*}
$$

It follows from Definition 2.4 that $f^{*} \in L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$. On the other hand, for any $q \in(0,1]$, if $f \in\left(B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)\right)^{*}$ and $f^{*} \in L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$, then, by [26, Lemma 2.8], we deduce that $S\left(\mathbb{R}^{n}\right) \subset B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)$ and hence $f \in S^{\prime}\left(\mathbb{R}^{n}\right)$, which, together with Definition 2.4, implies that $f \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$.

In the following, we recall the following notion of the multiresolution analysis on $\mathbb{R}$ (see $[20,24]$ for more details).
Definition 2.7. An increasing sequence $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$ of closed subspaces in $L^{2}(\mathbb{R})$ is called a multiresolution analysis (MRA) on $\mathbb{R}$ if
(i) $\bigcup_{j \in \mathbb{Z}} V_{j}$ is dense in $L^{2}(\mathbb{R})$ and $\bigcap_{j \in \mathbb{Z}} V_{j}=\{\theta\}$, where $\theta$ denotes the zero function;
(ii) for any $j \in \mathbb{Z}$ and $f \in L^{2}(\mathbb{R}), f \in V_{j}$ if and only if $f\left(2^{-j}\right) \in V_{0}$;
(iii) for any $k \in \mathbb{Z}$ and $f \in L^{2}(\mathbb{R}), f \in V_{0}$ if and only if $f(\cdot-k) \in V_{0}$;
(iv) there exists a function $\phi \in L^{2}(\mathbb{R})$ (called father wavelet) such that $\left\{\phi_{k}(\cdot)\right\}_{k \in \mathbb{Z}}:=\{\phi(\cdot-k)\}_{k \in \mathbb{Z}}$ forms an orthonormal basis of $V_{0}$.
For any fixed $s \in \mathbb{Z}_{+}$, according to [24, Theorem 1.61(ii)], we choose the father and the mother wavelets $\phi, \psi \in C_{c}^{s+1}(\mathbb{R})$, the set of all functions with compact supports having continuous derivatives up to order $s+1$, such that $\widehat{\phi}(0)=(2 \pi)^{-1 / 2}$ and, for any $l \in\{0, \ldots, s+1\}, \int_{\mathbb{R}} x^{l} \psi(x) d x=0$, where $\widehat{\phi}$ denotes the Fourier transform of $\phi$; namely, for any $\xi \in \mathbb{R}$,

$$
\widehat{\phi}(\xi):=\frac{1}{\sqrt{2 \pi}} \int_{\mathbb{R}} \phi(y) e^{-i \xi y} d y
$$

In the following, we assume that

$$
\begin{equation*}
\operatorname{supp} \phi, \operatorname{supp} \psi \subset 1 / 2+m(-1 / 2,1 / 2) \tag{2.5}
\end{equation*}
$$

namely, $x \in[(1-m) / 2,(1+m) / 2]$ if and only if $\left|x-\frac{1}{2}\right| \leqslant m / 2$. Here $m \in[1, \infty)$ is a positive constant independent of the main parameters involved in the whole paper.

By the standard procedure of tensor products, we can extend the above considerations from 1-dimension to $n$-dimension. More precisely, let

$$
\vec{\theta}_{n}:=(\overbrace{0, \ldots, 0}^{n \text { times }}) \text { and } E:=\{0,1\}^{n} \backslash\left\{\vec{\theta}_{n}\right\} .
$$

Assume that $\mathcal{D}$ is the set of all dyadic cubes in $\mathbb{R}^{n}$, i.e., for any $Q \in \mathcal{D}$, there exist $j \in \mathbb{Z}_{+}$and $k:=\left\{k_{1}, \ldots, k_{n}\right\} \in \mathbb{Z}^{n}$ such that

$$
\begin{equation*}
Q=Q_{j, k}:=\left\{x \in \mathbb{R}^{n}: k_{i} \leq 2^{j} x_{i}<k_{i}+1 \text { for any } i \in\{1, \ldots, n\}\right\} . \tag{2.6}
\end{equation*}
$$

Let $m Q$ be the $m$ dilation of $Q$ with the same center as $Q$ and $m$ as in (2.5). According to the tensor product in [20, p. 108], for any $\lambda:=\left(\lambda_{1}, \ldots, \lambda_{n}\right) \in E$, $Q:=Q_{j, k}$ with $k:=\left(k_{1}, \ldots, k_{n}\right) \in \mathbb{Z}^{n}, j \in \mathbb{Z}$, and $x=\left(x_{1}, \ldots, x_{n}\right)$, let

$$
\begin{aligned}
& \psi_{Q}^{\lambda}(x):=2^{j n / 2} \psi^{\lambda_{1}}\left(2^{j} x_{1}-k_{1}\right) \cdots \psi^{\lambda_{n}}\left(2^{j} x_{n}-k_{n}\right), \\
& \phi_{Q}(x):=2^{j n / 2} \phi\left(2^{j} x_{1}-k_{1}\right) \cdots \phi\left(2^{j} x_{n}-k_{n}\right)
\end{aligned}
$$

where $\psi^{0}:=\phi$ and $\psi^{1}:=\psi$.
A family $\left\{\psi_{Q}^{\lambda}\right\}_{Q \in \mathcal{D}, \lambda \in E} \subset C^{s+1}\left(\mathbb{R}^{n}\right)$ (the set of all functions having continuous derivatives up to order $s+1$ ) is called an $s$-order wavelet system (see $\left[8\right.$, p. 6]) if $\left\{\psi_{Q}^{\lambda}\right\}_{Q \in \mathcal{D}, \lambda \in E}$ satisfy
(i) $\left\{\psi_{Q}^{\lambda}\right\}_{Q \in \mathcal{D}, \lambda \in E}$ forms an orthonormal basis of $L^{2}\left(\mathbb{R}^{n}\right)$;
(ii) $\psi_{Q}^{\lambda}$ are compactly supported, namely,

$$
\operatorname{supp} \psi_{Q}^{\lambda} \subset m Q ;
$$

(iii) there exists a positive constant $C$, depending on $s$, such that, for any $\beta:=\left(\beta_{1}, \ldots, \beta_{n}\right) \in \mathbb{Z}_{+}^{n}$ with $|\beta|:=\beta_{1}+\cdots+\beta_{n} \leq s+1$,

$$
\begin{equation*}
\left|\partial^{\beta} \psi_{Q}^{\lambda}(x)\right| \leq C 2^{j|\beta|} 2^{j n / 2}, \quad \forall x \in \mathbb{R}^{n} ; \tag{2.7}
\end{equation*}
$$

(iv) for any $\beta:=\left(\beta_{1}, \ldots, \beta_{n}\right) \in \mathbb{Z}_{+}^{n}$ with $|\beta| \leq s, \int_{\mathbb{R}^{n}} x^{\beta} \psi_{Q}^{\lambda}(x) d x=0$, here and hereafter, for any $x:=\left(x_{1}, \ldots, x_{n}\right) \in \mathbb{R}^{n}, x^{\beta}:=x_{1}^{\beta_{1}} \cdots x_{n}^{\beta_{n}}$.
See more details in [20, p. 108].
Hence, for any $f \in L^{2}\left(\mathbb{R}^{n}\right)$, we find that

$$
f=\sum_{\lambda \in E} \sum_{Q \in \mathcal{D}}\left(f, \psi_{Q}^{\lambda}\right) \psi_{Q}^{\lambda}=\sum_{\lambda \in E} \sum_{(j, k) \in \mathbb{Z} \times \mathbb{Z}^{n}}\left(f, \psi_{j, k}^{\lambda}\right) \psi_{j, k}^{\lambda} \text { in } L^{2}\left(\mathbb{R}^{n}\right),
$$

and for any $k \in \mathbb{Z}^{n}, j \in \mathbb{Z}_{+}$with $Q=Q_{j, k} \in \mathcal{D}$ as in (2.6) and $\lambda \in E$,

$$
\psi_{j, k}^{\lambda}:=\psi_{Q}^{\lambda} .
$$

By [20, p. 142], for any $\lambda \in E$, we assume that there exists some set $W^{\lambda} \subset$ $[0,1)^{n}$ such that $0<\lambda \leq\left|W^{\lambda}\right|$ and $c_{0} \chi_{W^{\lambda}} \leq\left|\psi^{\lambda}\right|$ for some fixed positive constants $\gamma$ and $c_{0}$, where

$$
\begin{equation*}
\psi^{\lambda}:=\psi_{[0,1)^{n}}^{\lambda} \tag{2.8}
\end{equation*}
$$

For every $j \in \mathbb{Z}, \lambda \in E, k \in \mathbb{Z}^{n}$, and $Q:=Q_{j, k}$, let

$$
\begin{equation*}
W_{j, k}^{\lambda}:=\left\{x \in \mathbb{R}^{n}: 2^{j} x-k \in W^{\lambda}\right\}=: W_{Q}^{\lambda} \tag{2.9}
\end{equation*}
$$

Then, for each $\lambda \in E, j \in \mathbb{Z}$ and $k \in \mathbb{Z}^{n}$, we obtain

$$
\begin{equation*}
W_{j, k}^{\lambda} \subset Q_{j, k}, \quad\left|W_{j, k}^{\lambda}\right| \geq \gamma\left|Q_{j, k}\right| \tag{2.10}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|\psi_{j, k}^{\lambda}\right| \geq c_{0} \frac{\chi_{W_{j, k}^{\lambda}}}{\left|Q_{j, k}\right|} \tag{2.11}
\end{equation*}
$$

In the following, let

$$
\begin{equation*}
\Lambda:=\left\{(\lambda, j, k): \lambda \in E,(j, k) \in \mathbb{Z} \times \mathbb{Z}^{n}\right\} \tag{2.12}
\end{equation*}
$$

Further, for any $j \in \mathbb{Z}$, let $V_{j}$ be the closed subspace of $L^{2}\left(\mathbb{R}^{n}\right)$ spanned by $\left\{\phi_{Q}\right\}_{|Q|=2^{-j n}}$. It is known that $\left\{V_{j}\right\}_{j \in \mathbb{Z}}$ is an MRA on $\mathbb{R}^{n}$, whose definition extends MRA on $\mathbb{R}$ in Definition 2.7 (see [20, Chapter 2] for more details).

Next, we show that the wavelets belong to Campanato spaces with variable exponent.

For any $\alpha \in(0,1], s \in \mathbb{Z}_{+}$, and $\epsilon \in(0, \infty)$. Denote by $\mathcal{C}_{(\alpha, \epsilon), s}\left(\mathbb{R}^{n}\right)$ the class of all functions $\eta \in C^{s}\left(\mathbb{R}^{n}\right)$, the set of all functions having continuous derivatives up to order $s$, such that, for any $\nu \in \mathbb{Z}_{+}^{n}$, with $|\nu| \leq s$, and for any $x \in \mathbb{R}^{n}$,

$$
\begin{equation*}
\left|\partial^{\nu} \eta(x)\right| \leq(1+|x|)^{-n-\epsilon} \tag{2.13}
\end{equation*}
$$

and, for any $\nu \in \mathbb{Z}_{+}^{n}$, with $|\nu|=s$, and for any $x_{1}, x_{2} \in \mathbb{R}^{n}$,

$$
\begin{equation*}
\left|\partial^{\nu} \eta\left(x_{1}\right)-\partial^{\nu} \eta\left(x_{2}\right)\right| \leq\left|x_{1}-x_{2}\right|^{\alpha}\left[\left(1+\left|x_{1}\right|\right)^{-n-\epsilon}+\left(1+\left|x_{2}\right|\right)^{-n-\epsilon}\right] \tag{2.14}
\end{equation*}
$$

In the following, let $\mathcal{C}_{\epsilon, s}\left(\mathbb{R}^{n}\right):=\mathcal{C}_{(1, \epsilon), s}\left(\mathbb{R}^{n}\right)$.
The following results are [7, Proposition 1 and Corollary 2], respectively.
Proposition 2.8. Let $p(\cdot) \in C^{\log }\left(\mathbb{R}^{n}\right)$ and $s \in \mathbb{Z}_{+}$.
(i) if $\alpha \in(0,1], \varepsilon \in(\alpha+s, \infty)$ and $p_{-} \in(n /(n+\alpha+s), 1]$, then

$$
\mathcal{C}_{(\alpha, \epsilon), s}\left(\mathbb{R}^{n}\right) \subset B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)
$$

(ii) if $\varepsilon \in(1+s, \infty)$ and $p_{-} \in(n /(n+1+s), 1]$, then

$$
\mathcal{C}_{\epsilon, s}\left(\mathbb{R}^{n}\right) \subset B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)
$$

Corollary 2.9. Let $p(\cdot) \in C^{\log }\left(\mathbb{R}^{n}\right)$, $s \in \mathbb{Z}_{+}$and $p_{-} \in(n /(n+1+s), 1]$.
Then, for any $(\lambda, j, k) \in \Lambda$ with $\Lambda$ as in (2.12), $\psi_{j, k}^{\lambda} \in B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)$.
We now recall the definition of $(p(\cdot), r, s)$-atom introduced in [21].

Definition 2.10. Let $p(\cdot) \in \mathcal{P}\left(\mathbb{R}^{n}\right)$ and let $1<r \leq \infty$. Fix an integer

$$
\begin{equation*}
d \in\left(\frac{n}{p_{-}}-n-1, \infty\right) \cap \mathbb{Z}_{+} \tag{2.15}
\end{equation*}
$$

A measurable function $a$ on $\mathbb{R}^{n}$ is called a $(p(\cdot), r, d)$-atom if there exists a ball $B$ such that
(i) $\operatorname{supp} a \subset B$;
(ii) $\|a\|_{L^{r}\left(\mathbb{R}^{n}\right)} \leq \frac{|B|^{1 / r}}{\left\|\chi_{B}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}}$;
(iii) $\int_{\mathbb{R}^{n}} a(x) x^{\alpha} d x=0$ for all $\alpha \in \mathbb{Z}_{+}^{n}$ with $|\alpha| \leq d$.

We recall the notion of the variable atomic Hardy-Lorentz space, which is taken from [14, Definition 5.2].

Definition 2.11. Let $p(\cdot) \in \mathcal{P}\left(\mathbb{R}^{n}\right)$, let $0<q \leq \infty, 1<r \leq \infty$ and let $d$ be as (2.15). The variable atomic Hardy-Lorentz space $H_{\mathrm{atom}, r, d}^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ is defined as the space of all $f \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ which can be decomposed as

$$
\begin{equation*}
f=\sum_{i \in \mathbb{Z}} \sum_{j \in \mathbb{N}} \lambda_{i, j} a_{i, j} \quad \text { in } \quad S^{\prime}\left(\mathbb{R}^{n}\right) \tag{2.16}
\end{equation*}
$$

where $\left\{a_{i, j}\right\}_{i \in \mathbb{Z}, j \in \mathbb{N}}$ is a sequence of $(p(\cdot), r, d)$-atoms, associated with balls $\left\{B_{i, j}\right\}_{i \in \mathbb{Z}, j \in \mathbb{N}}$, satisfying that, for all $x \in \mathbb{R}^{n}$ and $i \in \mathbb{Z}, \sum_{j \in \mathbb{N}} \chi_{B_{i, j}}(x) \leq A$ with $A$ being a positive constant independent of $x$ and $i$; and for all $i \in \mathbb{Z}$ and $j \in \mathbb{N}, \lambda_{i, j}:=\tilde{A} 2^{i}\left\|\chi_{B_{i, j}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}$ with $\tilde{A}$ being a positive constant independent of $i$ and $j$. Moreover, for $f \in H_{\mathrm{atom}, r, d}^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$, we define

$$
\|f\|_{H_{\mathrm{atom}, r, d}^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}:=\inf \left(\sum_{i \in \mathbb{Z}}\left\|\left(\sum_{j \in \mathbb{N}}\left(\frac{\lambda_{i, j} \chi_{B_{i, j}}}{\left\|\chi_{B_{i, j}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}}\right)^{\underline{\underline{p}}}\right)^{\frac{1}{\underline{p}}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}^{q}\right)^{\frac{1}{q}},
$$

where the infimum is taken over all decompositions of $f$ as (2.16).
Then, we give the atomic characterization of $H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ from $[14$, Theorem 5.4].

Lemma 2.12. Let $p(\cdot) \in C^{\log }\left(\mathbb{R}^{n}\right)$, let $0<q \leq \infty$, let $r \in\left(\max \left\{p_{+}, 1\right\}, \infty\right]$ and let $d$ be as in (2.15). Then $H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)=H_{\mathrm{atom}, r, d}^{p(\cdot),}\left(\mathbb{R}^{n}\right)$ with equivalent quasi-norms.

In the following, we introduce some notation, let

$$
\mathcal{U}_{\psi,(\mathrm{i})} f:=\left[\sum_{(\lambda, j, k) \in \Lambda}\left|\left\langle f, \psi_{j, k}^{\lambda}\right\rangle\right|^{2}\left|\psi_{j, k}^{\lambda}\right|^{2}\right]^{1 / 2},
$$

$$
\begin{aligned}
& \mathcal{U}_{\psi,(\mathrm{ii)}} f:=\left[\sum_{(\lambda, j, k) \in \Lambda}\left|\left\langle f, \psi_{j, k}^{\lambda}\right\rangle\right|^{2} \frac{\chi_{Q_{j, k}}}{\left|Q_{j, k}\right|}\right]^{1 / 2}, \\
& \mathcal{U}_{\psi,(\mathrm{iii})} f:=\left[\sum_{(\lambda, j, k) \in \Lambda}\left|\left\langle f, \psi_{j, k}^{\lambda}\right\rangle\right|^{2} \frac{\chi_{W_{j, k}^{\lambda}}}{\left|Q_{j, k}\right|}\right]^{1 / 2},
\end{aligned}
$$

where $\Lambda$ is as in (2.12). By Corollary 2.9, we know that $\mathcal{U}_{\psi,(\mathrm{i})} f, \mathcal{U}_{\psi, \text { (ii) }} f$, and $\mathcal{U}_{\psi, \text { (iii) }} f$ are well defined.

Now, we recall the following definition of atoms introduced in [10, Definition 4.17].

Definition 2.13. Let $p(\cdot) \in \mathcal{P}\left(\mathbb{R}^{n}\right)$ with $p_{+}<r$ and $r \in(1, \infty)$, and let $\psi$ be the mother wavelet. A function $a \in L^{2}\left(\mathbb{R}^{n}\right)$ is called a $(p(\cdot), r, \psi)$-atom if there exists a dyadic cube $R$ such that

$$
a=\sum_{\lambda \in E} \sum_{Q \subset R, Q \in \mathcal{D}}\left(a, \psi_{Q}^{\lambda}\right) \psi_{Q}^{\lambda}
$$

supported in $m R$ with $m$ as in (2.5) and

$$
\begin{aligned}
\left\|\mathcal{U}_{\psi,(\mathrm{ii})} a\right\|_{L^{r}\left(\mathbb{R}^{n}\right)} & =\left\|\left[\sum_{\lambda \in E} \sum_{Q \subset R, Q \in \mathcal{D}}\left|\left(a, \psi_{Q}^{\lambda}\right)\right|^{2} \frac{\chi_{Q}}{|Q|}\right]^{1 / 2}\right\|_{L^{r}\left(\mathbb{R}^{n}\right)} \\
& \leqslant \frac{|m R|^{1 / r}}{\left\|\chi_{m R}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}} .
\end{aligned}
$$

Remark 2.14. By [8, Remark 2.15], we know that Definition 2.13 is well defined.
The following lemma is just [7, Lemma 1].
Lemma 2.15. Let $p(\cdot) \in \mathcal{P}\left(\mathbb{R}^{n}\right)$ with $p_{+}<r$ and $r \in(1, \infty)$, and let $\psi$ be the mother wavelet. If $a$ is a $(p(\cdot), r, \psi)$-atom related to a cube $R$, then there exists a positive harmless constant $c$, independent of $a$, such that $a / c$ is a $(p(\cdot), r, d)$ atom.

Lemma 2.16 ([21, Lemma 2.4]). Let $1<u<\infty$. Suppose $p(\cdot) \in C^{\log }\left(\mathbb{R}^{n}\right)$ satisfy $p_{-}>1$. Then there exists a positive constant $C$ such that, for any sequence of measurable functions $\left\{f_{j}\right\}_{j=1}^{\infty}$,

$$
\left\|\left(\sum_{j=1}^{\infty}\left[\mathcal{M} f_{j}\right]^{u}\right)^{1 / u}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)} \leqslant C\left\|\left(\sum_{j=1}^{\infty}\left|f_{j}\right|^{u}\right)^{1 / u}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)},
$$

where $\mathcal{M}$ denotes the Hardy-Littlewood maximal operator as in (2.3).

Lemma 2.17 ([14, Theorem 3.4]). Let $1<u<\infty$ and let $q \in(0, \infty]$. Suppose $p(\cdot) \in C^{\log }\left(\mathbb{R}^{n}\right)$ satisfy $p_{-}>1$. Then there exists a positive constant $C$ such that, for any sequence of measurable functions $\left\{f_{j}\right\}_{j=1}^{\infty}$,

$$
\left\|\left(\sum_{j=1}^{\infty}\left[\mathcal{M} f_{j}\right]^{u}\right)^{1 / u}\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \leqslant C\left\|\left(\sum_{j=1}^{\infty}\left|f_{j}\right|^{u}\right)^{1 / u}\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}
$$

where $\mathcal{M}$ denotes the Hardy-Littlewood maximal operator as in (2.3).
3. Intrinsic $\boldsymbol{g}$-function characterization of $\boldsymbol{H}^{p(\cdot), q}\left(\mathbb{R}^{\boldsymbol{n}}\right)$

In this section, firstly, we recall the definition of intrinsic $g$-functions from [17]. For any $\alpha \in(0,1]$ and $s \in \mathbb{Z}_{+}$, let $\mathcal{T}_{\alpha, s}\left(\mathbb{R}^{n}\right)$ be the class of all functions $\eta \in C^{s}\left(\mathbb{R}^{n}\right)$ such that $\operatorname{supp} \eta \subset\left\{x \in \mathbb{R}^{n}:|x| \leq 1\right\}$,

$$
\int_{\mathbb{R}^{n}} \eta(x) x^{\gamma} d x=0 \quad \text { for any } \gamma \in \mathbb{Z}_{+}^{n} \text { with }|\gamma| \leq s
$$

and there exists a positive constant $C$, depending on $s$, such that, for any $\nu \in \mathbb{Z}_{+}^{n}$, with $|\nu|=s$, and any $x_{1}, x_{2} \in \mathbb{R}^{n}$,

$$
\left|\partial^{\nu} \eta\left(x_{1}\right)-\partial^{\nu} \eta\left(x_{2}\right)\right| \leq C\left|x_{1}-x_{2}\right|^{\alpha}
$$

For any $f \in L_{\mathrm{loc}}^{1}\left(\mathbb{R}^{n}\right)$ and $(y, t) \in \mathbb{R}_{+}^{n+1}:=\mathbb{R}^{n} \times(0, \infty)$, let

$$
A_{\alpha, s}(f)(y, t):=\sup _{\eta \in \mathcal{T}_{\alpha, s}\left(\mathbb{R}^{n}\right)}\left|f * \eta_{t}(y)\right|,
$$

where $\eta_{t}(\cdot):=t^{-n} \eta(\dot{\bar{t}})$ for any $t \in(0, \infty)$. Then the intrinsic $g$-function from [17] is defined by setting, for any $x \in \mathbb{R}^{n}$,

$$
g_{\alpha, s}(f)(x):=\left\{\int_{0}^{\infty}\left[A_{\alpha, s}(f)(x, t)\right]^{2} \frac{d t}{t}\right\}^{1 / 2}
$$

Recall that, for all $f \in S^{\prime}\left(\mathbb{R}^{n}\right)$, the Littlewood-Paley $g$-function is defined by setting, for all $x \in \mathbb{R}^{n}$,

$$
g(f)(x):=\left(\int_{0}^{\infty}\left|f * \phi_{t}(x)\right|^{2} \frac{d t}{t}\right)^{1 / 2}
$$

where, $\phi \in S\left(\mathbb{R}^{n}\right)$ is a radial function satisfying [14, (8.1), (8.2) and (8.3)] and, for any $t \in(0, \infty), \phi_{t}(\cdot):=\frac{1}{t^{n}} \phi(\dot{\bar{t}})$.

Recall that $f \in S^{\prime}\left(\mathbb{R}^{n}\right)$ is said to vanish weakly at infinity if, for every $\psi \in S\left(\mathbb{R}^{n}\right), f * \psi_{t} \rightarrow 0$ in $S^{\prime}\left(\mathbb{R}^{n}\right)$ as $t \rightarrow \infty($ see $[6$, p. 50]).

The following result follows from [14, Theorem 8.2].
Lemma 3.1. Let $p(\cdot) \in C^{\log }\left(\mathbb{R}^{n}\right)$ and let $0<q \leq \infty$. Then $f \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ if and only if $f \in S^{\prime}\left(\mathbb{R}^{n}\right)$, $f$ vanishes weakly at infinity and $g(f) \in L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$. Moreover, for all $f \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$,

$$
C^{-1}\|g(f)\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \leq\|f\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \leq C\|g(f)\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}
$$

where $C$ is a positive constant independent of $f$.
The following lemma is a special case of [17, Proposition 3.2].
Lemma 3.2. Let $\alpha \in(0,1], s \in \mathbb{Z}_{+}$and $r \in(1, \infty)$. Then there exists a positive constant $C$ such that, for any $f \in \mathcal{A}\left(\mathbb{R}^{n}\right)$,

$$
\int_{\mathbb{R}^{n}}\left[g_{\alpha, s}(f)(x)\right]^{r} d x \leq C \int_{\mathbb{R}^{n}}|f(x)|^{r} d x .
$$

Now we are ready to state and prove the main result of this section.
Theorem 3.3. Let $p(\cdot) \in C^{\log }\left(\mathbb{R}^{n}\right)$, $p_{+} \in(0,1]$ and $q \in(0,1]$. Suppose that $\alpha \in(0,1]$, $s \in \mathbb{Z}_{+}$and $p_{-} \in(n / n+\alpha+s, 1]$. Then $f \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ if and only if $f \in\left(B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)\right)^{*}$, the dual space of $B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)$, $f$ vanishes weakly at infinity and $g_{\alpha, s}(f) \in L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$; moreover, it holds true that

$$
\frac{1}{C}\left\|g_{\alpha, s}(f)\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \leq\|f\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \leq C\left\|g_{\alpha, s}(f)\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}
$$

with $C$ being a positive constant independent of $f$.
Proof. Let $q \in(0,1], f \in\left(B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)\right)^{*}$ vanish weakly at infinity and $g_{\alpha, s}(f) \in L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$. Then, by Proposition 2.8 , we find that $f \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$. Notice that, for all $x \in \mathbb{R}^{n}, g(f)(x) \lesssim g_{\alpha, s}(f)(x)$, it follows that $g(f) \in$ $L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$. From this and Lemma 3.1, we find that there exists a distribution $\tilde{f} \in \mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ such that $\tilde{f}=f$ in $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right), \tilde{f} \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ and $\|\widetilde{f}\|_{H^{p(\cdot), q\left(\mathbb{R}^{n}\right)}} \lesssim$ $\|g(f)\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}$, which, together with [14, Lemma 8.4] and the fact that $f$ vanishes weakly at infinity, implies that $f=\tilde{f}$ in $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ and hence

$$
\|f\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \sim\|\widetilde{f}\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \lesssim\|g(f)\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \lesssim\left\|g_{\alpha, s}(f)\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} .
$$

This finishes the proof of the sufficiency of Theorem 3.3.
It remains to prove the necessity. For $q \in(0,1]$, let $f \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$. Then, by [14, Lemma 8.4], we see that $f$ vanishes weakly at infinity and, by Lemma 2.12 and [14, Theorem 7.2], we have $f \in\left(B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)\right)^{*}$. By Lemma 2.12, there exist sequences of $(p(\cdot), \infty, d)$-atoms $\left\{a_{i, j}\right\}_{i \in \mathbb{Z}, j \in \mathbb{N}}$ and nonnegative numbers $\left\{\lambda_{i, j}\right\}_{i \in \mathbb{Z}, j \in \mathbb{N}}$ such that the series $\sum_{i \in \mathbb{Z}} \sum_{j \in \mathbb{N}} \lambda_{i, j} a_{i, j}$ converges to $f$ in $\mathcal{S}^{\prime}\left(\mathbb{R}^{n}\right)$ and $\lambda_{i, j} \approx 2^{i}\left\|\chi_{B_{i, j}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}$. For $i_{0} \in \mathbb{Z}$, let

$$
f=\sum_{i=-\infty}^{i_{0}-1} \sum_{j \in \mathbb{N}} \lambda_{i, j} a_{i, j}+\sum_{i=i_{0}}^{\infty} \sum_{j \in \mathbb{N}} \lambda_{i, j} a_{i, j}=: f_{1}+f_{2} .
$$

Hence, we get

$$
\begin{aligned}
& \left\|\chi_{\left\{x \in \mathbb{R}^{n}: g_{\alpha, s}(f)(x)>2^{i_{0}}\right\}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)} \\
\lesssim & \left\|\chi_{\left\{x \in \mathbb{R}^{n}: g_{\alpha, s}\left(f_{1}\right)(x)>2^{i_{0}-1}\right\}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}+\left\|\chi_{\left\{x \in A_{i_{0}}: g_{\alpha, s}\left(f_{2}\right)(x)>2^{i_{0}-1}\right\}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)} \\
& +\left\|\chi_{\left\{x \in A_{i_{0}}^{\mathrm{C}}: g_{\alpha, s}\left(f_{2}\right)(x)>2^{i_{0}-1}\right\}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}
\end{aligned}
$$

$$
=: J_{1}+J_{2}+J_{3},
$$

where $A_{i_{0}}:=\bigcup_{i=i_{0}}^{\infty} \bigcup_{j \in \mathbb{N}}\left(4 B_{i, j}\right)$ and $\left\{B_{i, j}\right\}_{i \in \mathbb{Z}, j \in \mathbb{N}}$ are the balls as in [14, Theorem 5.4].

Let $\delta_{1}, \delta_{2}, \delta_{3}$ and $\delta$ be the same as in [14, Theorem 5.4]. We first estimate $J_{1}$. It is obvious that

$$
\begin{aligned}
J_{1} \lesssim & \left.\| \chi_{\left\{x \in \mathbb{R}^{n}:\right.} \sum_{i=-\infty}^{i_{0}-1} \sum_{j \in \mathbb{N}} \lambda_{i, j} g_{\alpha, s}\left(a_{i, j}\right)(x) \chi_{4 B_{i, j}>2^{i} 0-2}\right\} \|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)} \\
& +\left\|\chi_{\left\{x \in \mathbb{R}^{n}: \sum_{i=-\infty}^{i_{0}-1} \sum_{j \in \mathbb{N}} \lambda_{i, j} g_{\alpha, s}\left(a_{i, j}\right)(x) \chi_{\left(4 B_{i, j}\right)^{\mathrm{C}}}>2^{i_{0}-2}\right\}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)} \\
= & J_{1,1}+J_{1,2} .
\end{aligned}
$$

We now estimate $J_{1,1}$. By Lemma 3.2 and the proof of [14, (5.8)], we obtain

$$
J_{1,1} \lesssim 2^{-i_{0} \delta_{1}}\left(\sum_{i=-\infty}^{i_{0}-1} 2^{i q \delta_{1}}\left\|\sum_{j \in \mathbb{N}} \chi_{B_{i, j}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}^{q}\right)^{\frac{1}{q}}
$$

and

$$
\begin{equation*}
\sum_{i_{0}=-\infty}^{\infty} 2^{i_{0} q}\left\|\chi_{\left\{x \in \mathbb{R}^{n}: g_{\alpha, s}\left(f_{1}\right) \chi_{4 B_{i, j}}(x)>2^{i_{0}}\right\}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}^{q} \lesssim\|f\|_{H_{\mathrm{atom}, r, d}^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \tag{3.1}
\end{equation*}
$$

Next, we deal with $J_{1,2}$. By the similar argument that used in [26, p. 1564], for all $i \in \mathbb{Z}, j \in \mathbb{N}$ and $x \in\left(4 \chi_{B_{i}, j}\right)^{\complement}$, we obtain

$$
\begin{equation*}
\left|g_{\alpha, s}\left(a_{i, j}\right)(x)\right| \lesssim\left\|\chi_{B_{i, j}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}^{-1}\left(\mathcal{M}\left(\chi_{B_{i, j}}\right)(x)\right)^{\frac{n+s+\alpha}{n}} \tag{3.2}
\end{equation*}
$$

Then by the Hölder inequality and a similar argument that used in the proof of $[14,(5.10)]$, we get

$$
J_{1,2} \lesssim 2^{-i_{0} \delta_{2}}\left(\sum_{i=-\infty}^{i_{0}-1} 2^{i q \delta_{2}}\left\|\sum_{j \in \mathbb{N}} \chi_{B_{i, j}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}^{q}\right)^{\frac{1}{q}}
$$

and
(3.3) $\left.\sum_{i_{0}=-\infty}^{\infty} 2^{i_{0} q} \| \chi_{\left\{x \in \mathbb{R}^{n}: g_{\alpha, s}\left(f_{1}\right) \chi_{\left(4 B_{i, j}\right)^{\mathrm{C}}}(x)>2^{i}\right\}}\right\}\left\|_{L^{p(\cdot)\left(\mathbb{R}^{n}\right)}}^{q} \lesssim\right\| f \|_{H_{\mathrm{atom}, r, d}^{p(\cdot),}\left(\mathbb{R}^{n}\right)}^{q}$.

For $I_{2}$, by an argument similar to that used in the proof of $[14,(5.11)]$, we get

$$
J_{2} \leq\left\|\chi_{A_{i_{0}}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)} \lesssim 2^{-i_{0} \delta_{3}}\left(\sum_{i=i_{0}}^{\infty} 2^{i \delta_{3} q}\left\|\sum_{j \in \mathbb{N}} \chi_{B_{i, j}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}^{q}\right)^{\frac{1}{q}}
$$

and

$$
\begin{equation*}
\sum_{i_{0}=-\infty}^{\infty} 2^{i_{0} q}\left\|\chi_{\left\{x \in A_{i_{0}}: g_{\alpha, s}\left(f_{2}\right)(x)>2^{i_{0}}\right\}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}^{q} \lesssim\|f\|_{H_{\mathrm{atom}, r, d}^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}^{q} . \tag{3.4}
\end{equation*}
$$

For $J_{3}$, by (3.2), Lemma 2.16 and an argument similar to that used in the proof of $[14,(5.12)]$, we find that

$$
J_{3} \lesssim 2^{-i_{0} \delta}\left(\sum_{i=i_{0}}^{\infty} 2^{i \delta q}\left\|\sum_{j \in \mathbb{N}} \chi_{B_{i, j}}\right\|_{L^{p(\cdot)}}^{q}\right)^{\frac{1}{q}}
$$

and

$$
\begin{equation*}
\sum_{i_{0}=-\infty}^{\infty} 2^{i_{0} q}\left\|\chi_{\left\{x \in\left(A_{i_{0}}\right)^{\mathrm{C}}: g_{\alpha, s}\left(f_{2}\right)(x)>2^{i_{0}}\right\}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}^{q} \lesssim\|f\|_{H_{\mathrm{atom}, r, d}^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}^{q} . \tag{3.5}
\end{equation*}
$$

Finally, combining (3.1), (3.3), (3.4) and (3.5), we obtain

$$
\left\|g_{\alpha, s}(f)\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \lesssim\|f\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}
$$

The proof is complete.
Now, we recall a discrete variant of the Littlewood-Paley $g$-function $\widetilde{g}^{\lambda}(f)$ from [8]. For any $\lambda \in E, f \in L^{2}\left(\mathbb{R}^{n}\right)$, and $x \in \mathbb{R}^{n}$, let

$$
\widetilde{g}^{\lambda}(f)(x):=\left[\sum_{j \in \mathbb{Z}}\left|f * \psi_{2^{-j}}^{\lambda}(x)\right|^{2}\right]^{1 / 2}
$$

where $\psi^{\lambda}$ is as in (2.8).
Lemma 3.4 ([17, Theorem 2.6]). Let $\lambda \in E$, $s \in \mathbb{Z}_{+}, \alpha \in(0,1]$, and $\varepsilon \in$ $(\max \{s, \alpha\}, \infty)$. Then there exists a positive constant $C$ such that, for any $f$ satisfying

$$
\begin{equation*}
|f(\cdot)|(1+|\cdot|)^{-n-\varepsilon} \in L^{1}\left(\mathbb{R}^{n}\right) \tag{3.6}
\end{equation*}
$$

it holds that

$$
\widetilde{g}^{\lambda}(f)(x) \leqslant C g_{\alpha, s}(f)(x), \quad \forall x \in \mathbb{R}^{n}
$$

The following conclusion follows from Theorem 3.3 and Lemma 3.4, the details being omitted.

Proposition 3.5. Let $\lambda \in E, s \in \mathbb{Z}_{+}, p(\cdot) \in C^{\log }\left(\mathbb{R}^{n}\right), p_{-} \in(n / n+\alpha+s, 1]$ and let $0<q \leq 1$. If $f \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right) \cap L^{2}\left(\mathbb{R}^{n}\right)$, then $\widetilde{g}^{\lambda}(f) \in L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$; moreover, there exists a positive constant $C_{(\lambda)}$, depending on $\lambda$, such that, for any $f \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right) \cap L^{2}\left(\mathbb{R}^{n}\right)$,

$$
\left\|\widetilde{g}^{\lambda}(f)\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \leqslant C_{(\lambda)}\|f\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}
$$

For any $\lambda \in E, j \in \mathbb{Z}, \nu \in(0, \infty), f \in L^{2}\left(\mathbb{R}^{n}\right)$, and $x \in \mathbb{R}^{n}$, we recall a variant of the Peetre type maximal functions from [8] defined by setting,

$$
\psi_{j, \nu}^{\lambda, * *}(f)(x):=\sup _{y \in \mathbb{R}^{n}} \frac{\left|f * \psi_{2^{-j}}^{\lambda}(x-y)\right|}{\left[1+2^{j}|y|\right]^{\nu}}
$$

From some arguments similar to those used in the proof of [10, Proposition 4.8], we obtain the following result.

Proposition 3.6. Let $s \in \mathbb{Z}_{+}, p(\cdot) \in C^{\log }\left(\mathbb{R}^{n}\right), r \in(1, \infty), \nu \in(\max \{1 / 2$, $\left.\left.1 / p_{-}\right\}, \infty\right)$ and $q \in(0,1]$. Then there exists a positive constant $C_{(\lambda, \nu)}$, depending on $\lambda$ and $\nu$, such that, for any $f \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right) \cap L^{r}\left(\mathbb{R}^{n}\right)$,

$$
\left\|\left\{\sum_{j=-\infty}^{\infty}\left|\psi_{j, \nu}^{\lambda, * *}(f)\right|^{2}\right\}^{1 / 2}\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \leqslant C_{(\lambda, \nu)}\|f\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}
$$

Proof. By some arguments similar to those used in the proof of [11, p. 271] in 1 -dimensional case, for any $j \in \mathbb{Z}, f \in L^{r}\left(\mathbb{R}^{n}\right), \nu \in(0, \infty)$, and $x \in \mathbb{R}^{n}$, we find that,

$$
\psi_{j, \nu}^{\lambda, * *} f(x) \lesssim\left[\mathcal{M}\left(\left|f * \psi_{2^{-j}}^{\lambda}\right|^{1 / \nu}\right)(x)\right]^{\nu}
$$

with the implicit positive constant depending only on $\lambda, \nu$, and $n$.
Combining this, the Fefferman-Stein vector-valued maximal inequality, Theorem 3.3 and Lemma 3.4, for any $f \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right) \cap L^{r}\left(\mathbb{R}^{n}\right)$, we obtain

$$
\begin{aligned}
& \left\|\left[\sum_{j \in \mathbb{Z}}\left|\psi_{j, \nu}^{\lambda, * *}(f)\right|^{2}\right]^{1 / 2}\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \\
\lesssim & \left\|\left\{\sum_{j \in \mathbb{Z}}\left[\mathcal{M}\left(\left|\psi_{2^{-j}}^{\lambda} * f\right|^{1 / \nu}\right)\right]^{2 \nu}\right\}^{1 / 2}\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \\
\sim & \left\|\left\{\sum_{j \in \mathbb{Z}}\left[\mathcal{M}\left(\left|\psi_{2^{-j}}^{\lambda} * f\right|^{1 / \nu}\right)\right]^{2 \nu}\right\}^{1 /(2 \nu)}\right\|^{\nu} \|_{L^{\nu p(\cdot), \nu q}\left(\mathbb{R}^{n}\right)}
\end{aligned}
$$

$$
\begin{aligned}
& \lesssim\left\|\left\{\sum_{j \in \mathbb{Z}}\left|\psi_{2^{-j}}^{\lambda} * f\right|^{2}\right\}^{1 /(2 \nu)}\right\|^{\nu} \|_{L^{\nu p(\cdot), \nu q}\left(\mathbb{R}^{n}\right)}^{\nu} \\
& \sim\left\|\widetilde{g}^{\lambda}(f)\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \\
& \lesssim\left\|g_{\alpha, s}(f)\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \\
& \lesssim\|f\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}
\end{aligned}
$$

where $\alpha \in(0,1], s \in \mathbb{Z}_{+}$, and $(p \nu)_{-}=p_{-} \nu>1$. The proof is complete.

## 4. Wavelet characterizations of $\boldsymbol{H}^{p(\cdot), q}\left(\mathbb{R}^{\boldsymbol{n}}\right)$

In this section, we provide several equivalent characterizations of the variable Hardy-Lorentz space $H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ via wavelets.

Let $d$ be an integer satisfying (2.4). By Remark 2.6(ii), it follows that if $f \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ for $0<q \leq 1$, then $f \in\left(B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)\right)^{*}$. Further, from Corollary 2.9, it follows that $\psi_{j, k}^{\lambda} \in B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)$. Let $\langle\cdot, \cdot\rangle$ denote the dual relation between $B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)$ and $\left(B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)\right)^{*}$. Hence, following an idea used in [20, p. 177], it follows that $\left\langle f, \psi_{j, k}^{\lambda}\right\rangle$ is well defined in the sense of the duality between $B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)$ and $\left(B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)\right)^{*}$.

Now we state the main results of this section.
Theorem 4.1. Let $p(\cdot) \in C^{\log }\left(\mathbb{R}^{n}\right), q \in(0,1]$ and $1 \leq r<\infty$. Let $d$ be an integer satisfying

$$
\frac{n}{n+1+d}<p_{-} \leqslant p_{+} \leqslant 1
$$

and suppose $\left\{\psi_{j, k}^{\lambda}\right\}_{(\lambda, j, k) \in \Lambda}$ is a d-order wavelet system.
For any $f \in\left(B M O_{p(\cdot), r}\left(\mathbb{R}^{n}\right)\right)^{*}$, assume that

$$
\begin{equation*}
f=\sum_{(\lambda, j, k) \in \Lambda}\left\langle f, \psi_{j, k}^{\lambda}\right\rangle \psi_{j, k}^{\lambda} \tag{4.1}
\end{equation*}
$$

Then the following statements are mutually equivalent:
(i) $f \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$;
(ii)

$$
\|f\|_{(\mathrm{i})}:=\left\|\left[\sum_{(\lambda, j, k) \in \Lambda}\left|\left\langle f, \psi_{j, k}^{\lambda}\right\rangle\right|^{2}\left|\psi_{j, k}^{\lambda}\right|^{2}\right]^{1 / 2}\right\|_{L^{p(\cdot), q\left(\mathbb{R}^{n}\right)}}<\infty ;
$$

(iii)

$$
\|f\|_{(\text {ii) }}:=\left\|\left[\sum_{(\lambda, j, k) \in \Lambda}\left|\left\langle f, \psi_{j, k}^{\lambda}\right\rangle\right|^{2} \frac{\chi_{Q_{j, k}}}{\left|Q_{j, k}\right|}\right]^{1 / 2}\right\|_{L^{p(\cdot), q\left(\mathbb{R}^{n}\right)}}<\infty ;
$$

(iv)

$$
\|f\|_{(\text {iii })}:=\left\|\left[\sum_{(\lambda, j, k) \in \Lambda}\left|\left\langle f, \psi_{j, k}^{\lambda}\right\rangle\right|^{2} \frac{\chi_{W_{j, k}^{\lambda}}}{\left|Q_{j, k}\right|}\right]^{1 / 2}\right\|_{L^{p(\cdot), q\left(\mathbb{R}^{n}\right)}}<\infty,
$$

where, for any $(\lambda, j, k) \in \Lambda, W_{j, k}^{\lambda} \subset Q_{j, k}$ is as in (2.9) and

$$
\begin{equation*}
\left|W_{j, k}^{\lambda}\right| \sim\left|Q_{j, k}\right| \tag{4.2}
\end{equation*}
$$

with the implicit positive constants independent of $(\lambda, j, k)$.
Moreover, all the quasi-norms $\|\cdot\|_{(\mathrm{i})},\|\cdot\|_{(\mathrm{ii})}$, and $\|\cdot\|_{(\mathrm{iii})}$ are equivalent to $\|\cdot\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}$.

Proof. We observe that (4.2) follows from (2.10). Then, we only need to prove that (i) through (iv) of Theorem 4.1 are mutually equivalent. Indeed, we prove (i) $\Rightarrow$ (ii) $\Rightarrow$ (iv) $\Rightarrow$ (iii) $\Rightarrow$ (i).
(i) $\Rightarrow$ (ii) Notice that $H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$ is a quasi-Banach space and $H^{p(\cdot), q}\left(\mathbb{R}^{n}\right) \cap$ $L^{r}\left(\mathbb{R}^{n}\right)$ with $r$ as in Lemma 2.12 is dense in $H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$. Hence, it suffices to prove that, for any $f \in H^{p(\cdot), q}\left(\mathbb{R}^{n}\right) \cap L^{r}\left(\mathbb{R}^{n}\right)$,

$$
\left\|\mathcal{U}_{\psi,(\mathrm{i})} f\right\|_{L^{p(\cdot), q\left(\mathbb{R}^{n}\right)}} \lesssim\|f\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}
$$

Indeed, from the proof of $[8$, Theorem 1.9], for any $(\lambda, j, k) \in \Lambda$ and $f \in$ $H^{p(\cdot), q}\left(\mathbb{R}^{n}\right) \cap L^{r}\left(\mathbb{R}^{n}\right)$, we have that,

$$
\left|\left(f, \psi_{j, k}^{\lambda}\right)\right| \lesssim 2^{-j n / 2} \sup _{y \in Q_{j, k}^{\lambda}}\left|\widetilde{\psi}_{2^{-j}}^{\lambda} * f(y)\right|
$$

with $\widetilde{\psi}(x):=\overline{\psi(-x)}$ for any $x \in \mathbb{R}^{n}$ and, for almost every $x \in \mathbb{R}^{n}$,

$$
\sum_{k \in \mathbb{Z}^{n}}\left|\left\langle f, \psi_{j, k}^{\lambda}\right\rangle\right|^{2}\left|\psi_{j, k}^{\lambda}(x)\right|^{2} \lesssim\left[\psi_{j, \nu}^{\lambda, * *} f(x)\right]^{2},
$$

where the implicit positive constants depend only on $\nu, m$, and $n$ with $m$ as in (2.5).

By some arguments similar to $[7,(17)]$, we select $\nu \in\left(\max \left\{1 / 2,1 / p_{-}\right\}, \infty\right)$. Combining these facts and Proposition 3.6, we get

$$
\begin{aligned}
& \left\|\mathcal{U}_{\psi,(\mathrm{i})} f\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}^{p_{-}} \\
\lesssim & \left\|\sum_{\lambda \in E}\left\{\sum_{j=-\infty}^{\infty} \sum_{k \in \mathbb{Z}^{n}}\left|\left\langle f, \psi_{j, k}^{\lambda}\right\rangle\right|^{2}\left|\psi_{j, k}^{\lambda}\right|^{2}\right\}^{1 / 2}\right\|_{p^{p_{-}}} \\
\lesssim & \sum_{\lambda \in E}\left\|\left\{\sum_{j=-\infty}^{\infty} \sum_{k \in \mathbb{Z}^{p}}\left|\left\langle f, \psi_{j, k}^{\lambda}\right\rangle\right|^{2}\left|\psi_{j, k}^{\lambda}\right|^{2}\right\}^{1 / 2}\right\|_{\mathbb{R}^{p-}\left(\mathbb{R}^{n}\right)}^{p_{-}}
\end{aligned} \|_{L^{p(\cdot),( }\left(\mathbb{R}^{n}\right)}
$$

$$
\begin{aligned}
& \left.\lesssim \sum_{\lambda \in E}\| \|_{j=-\infty}^{\infty}\left|\psi_{j, \nu}^{\lambda,, *}(f)\right|^{2}\right\}^{1 / 2} \|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}^{p_{-}} \\
& \lesssim\|f\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}^{p_{-}}
\end{aligned}
$$

which completes the proof for (i) $\Rightarrow$ (ii).
(ii) $\Rightarrow$ (iv) It is an easy consequence of the fact (2.11). Moreover, we obtain

$$
\begin{equation*}
\left\|\mathcal{U}_{\psi,(\mathrm{iii})} f\right\|_{L^{p(\cdot), q\left(\mathbb{R}^{n}\right)}} \leqslant\left\|\mathcal{U}_{\psi,(\mathrm{i})} f\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} . \tag{4.4}
\end{equation*}
$$

(iv) $\Rightarrow$ (iii) By $[8,(3.3)]$, we have that, for any $s \in(0, \infty)$ and $(\lambda, j, k) \in \Lambda$,

$$
\begin{equation*}
\chi_{Q_{j, k}} \lesssim\left[\mathcal{M}\left(\chi_{W_{j, k}^{\lambda}}\right)\right]^{1 / s} \tag{4.5}
\end{equation*}
$$

Moreover, choosing $\frac{1}{e} \in\left(\max \left\{1 / 2,1 / p_{-}\right\}, \infty\right)$, combining (4.5) and the Fefferman-Stein vector valued maximal inequality (see Lemma 2.17) with $u$ replaced by $2 / e$ and $\frac{1}{e} \in\left(\max \left\{1 / 2,1 / p_{-}\right\}, \infty\right)$, we get

$$
\begin{aligned}
& \|f\|_{(\mathrm{ii})} \\
\lesssim & \left\|\left\{\sum_{(\lambda, j, k) \in \Lambda} \frac{\left|\left\langle f, \psi_{j, k}^{\lambda}\right\rangle\right|^{2}}{\left|Q_{j, k}\right|}\left[\mathcal{M}\left(\chi_{W_{j, k}^{\lambda}}\right)\right]^{2 / e}\right\}^{1 / 2}\right\|_{L^{p(\cdot), q\left(\mathbb{R}^{n}\right)}} \\
\sim & \left.\| \sum_{(\lambda, j, k) \in \Lambda}\left[\mathcal{M}\left(\left[\frac{\left|\left\langle f, \psi_{j, k}^{\lambda}\right\rangle\right|^{1 / 2}}{\left|Q_{j, k}\right|^{1 / 2}} \chi_{W_{j, k}^{\lambda}}\right]^{e}\right)\right]^{2 / e}\right\}^{e / 2}\left\|_{L^{1 / e}}\right\|_{L^{p(\cdot) / e, q / e}\left(\mathbb{R}^{n}\right)} \\
\lesssim & \left.\sum_{(\lambda, j, k) \in \Lambda}\left[\frac{\left|\left\langle f, \psi_{j, k}^{\lambda}\right\rangle\right|^{2}}{\left|Q_{j, k}\right|^{1 / 2}} \chi_{W_{j, k}^{\lambda}}\right]^{2}\right\}^{e / 2} \|_{L^{p(\cdot) / e, q / e}\left(\mathbb{R}^{n}\right)}^{1 / e} \\
\sim & \left\|\left\{\sum_{(\lambda, j, k) \in \Lambda}\left|\left\langle f, \psi_{j, k}^{\lambda}\right\rangle\right|^{2} \frac{\chi_{W_{j, k}^{\lambda}}}{\left|Q_{j, k}\right|}\right\}^{1 / 2}\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \\
\sim & \|f\|_{(\mathrm{iii})},
\end{aligned}
$$

where

$$
\left(\frac{p}{e}\right)_{-}=\frac{1}{e} p_{-}>1 .
$$

This shows that (iv) $\Rightarrow$ (iii).
(iii) $\Rightarrow$ (i) From some arguments similar to those used in the proof of (iii) $\Rightarrow$ (i) in $\left[8\right.$, Theorem 1.9], we only need to prove that, for any $f \in L^{2}\left(\mathbb{R}^{n}\right)$ with
$\mathcal{U}_{\psi,(\mathrm{ii})} f \in L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$,

$$
\begin{equation*}
\|f\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \lesssim\left\|\mathcal{U}_{\psi,(\mathrm{ii})} f\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \tag{4.7}
\end{equation*}
$$

For any $f \in L^{2}\left(\mathbb{R}^{n}\right)$ with $\mathcal{U}_{\psi,(\mathrm{ii)}} f \in L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$, we aim to show

$$
\begin{equation*}
f=\sum_{Q \in \mathcal{D}}\left(f, \psi_{Q}^{\lambda}\right) \psi_{Q}^{\lambda}=\sum_{k \in \mathbb{Z}} \sum_{i \in \Delta_{k}} b(k, i), \tag{4.8}
\end{equation*}
$$

where $\left\{b(k, i): k \in \mathbb{Z}, i \in \Delta_{k}\right\}$ are some multiples of $(p(\cdot), r, \psi)$-atoms with $p(\cdot)$ and $r$ as in Lemma 2.12, and $\psi$ is the mother wavelet, which will be determined later. For any $k \in \mathbb{Z}$, let

$$
\begin{gathered}
\Omega_{k}:=\left\{x \in \mathbb{R}^{n}: \mathcal{U}_{\psi,(\mathrm{ii})} f(x)>2^{k}\right\}, \\
\mathcal{D}_{k}:=\left\{Q \in \mathcal{D}:\left|Q \cap \Omega_{k}\right| \geqslant \frac{1}{2}|Q|,\left|Q \cap \Omega_{k+1}\right|<\frac{1}{2}|Q|\right\},
\end{gathered}
$$

and $\widetilde{\mathcal{D}}:=\bigcup_{k \in \mathbb{Z}} \mathcal{D}_{k}$.
From the proof of [8, Theorem 1.7], we have that, for any $Q \in \widetilde{\mathcal{D}}$, there exists a unique $k \in \mathbb{Z}$ such that $Q \in \mathcal{D}_{k}$ and, for any $f \in L^{2}\left(\mathbb{R}^{n}\right)$ such that $\mathcal{U}_{\psi,(\mathrm{ii)}} f \in L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$, and $Q \in \mathcal{D} \backslash \widetilde{\mathcal{D}}$,

$$
\begin{equation*}
\left\langle f, \psi_{Q}^{\lambda}\right\rangle=0 \tag{4.9}
\end{equation*}
$$

Observe that, due to the nesting property of dyadic cubes, for any $Q \in \mathcal{D}_{k}$, there exists a unique maximal dyadic cube $\widetilde{Q} \in \mathcal{D}_{k}$ such that $Q \subset \widetilde{Q}$. Let $\left\{\widetilde{Q}_{k}^{i} \in \mathcal{D}_{k}: i \in \Delta_{k}\right\}$ be the collection of all such maximal dyadic cubes in $\mathcal{D}_{k}$. Then

$$
\widetilde{\mathcal{D}}=\bigcup_{k \in \mathbb{Z}} \mathcal{D}_{k}=\bigcup_{k \in \mathbb{Z}} \bigcup_{i \in \Delta_{k}}\left\{Q \in \mathcal{D}_{k}: Q \subset \widetilde{Q}_{k}^{i}\right\}
$$

By (4.8) and (4.9), we find that, for any $f \in L^{2}\left(\mathbb{R}^{n}\right)$ with $\mathcal{U}_{\psi,(i i)} f \in L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)$,

$$
\begin{aligned}
f & =\sum_{\lambda \in E} \sum_{Q \in \mathcal{D}}\left(f, \psi_{Q}^{\lambda}\right) \psi_{Q}^{\lambda} \\
& =\sum_{\lambda \in E} \sum_{Q \in \widetilde{\mathcal{D}}}\left(f, \psi_{Q}^{\lambda}\right) \psi_{Q}^{\lambda} \\
& =\sum_{\lambda \in E} \sum_{k \in \mathbb{Z}} \sum_{i \in \Delta_{k}}\left\{\sum_{\left\{Q \subset \widetilde{Q}_{k}^{i}, Q \in \mathcal{D}_{k}\right\}}\left(f, \psi_{Q}^{\lambda}\right) \psi_{Q}^{\lambda}\right\} \\
& =: \sum_{k \in \mathbb{Z}} \sum_{i \in \Delta_{k}} b(k, i),
\end{aligned}
$$

where, for any $k \in \mathbb{Z}$ and $i \in \Delta_{k}$,

$$
b(k, i):=\sum_{\lambda \in E} \sum_{\left\{Q \subset \widetilde{Q}_{k}^{i}, Q \in \mathcal{D}_{k}\right\}}\left(f, \psi_{Q}^{\lambda}\right) \psi_{Q}^{\lambda}
$$

According to the estimate in [7, p. 754], we know that $b(k, i)$ is a multiple of $(p(\cdot), r, \psi)$-atom.

Let $m \in \mathbb{N}$ satisfy $2 m \geqslant r$. According to the estimate in [7, (24)], we have

$$
\begin{equation*}
\left\|\mathcal{U}_{\psi,(\mathrm{ii})} b(k, i)\right\|_{L^{r}\left(\mathbb{R}^{n}\right)} \lesssim 2^{k}\left|\widetilde{Q}_{k}^{i}\right|^{1 / r} \tag{4.10}
\end{equation*}
$$

By some arguments similar to (4.5) and the fact that

$$
\left|m \widetilde{Q}_{k}^{i}\right| \sim\left|\widetilde{Q}_{k}^{i}\right| \lesssim\left|\widetilde{Q}_{k}^{i} \cap \Omega_{k}\right|
$$

for any $t \in(0, \infty)$, we deduce that,

$$
\begin{equation*}
\chi_{m \widetilde{Q}_{j, k}} \lesssim\left[\mathcal{M}\left(\chi_{\widetilde{Q}_{k}^{i} \cap \Omega_{k}}\right)\right]^{1 / t} \tag{4.11}
\end{equation*}
$$

This, together with (4.10), the Fefferman-Stein vector-valued maximal inequality, some arguments similar to those used in the estimate of (4.6), and the disjointness of $\left\{\widetilde{Q}_{k}^{i}\right\}_{i \in \Delta_{k}}$, implies that

$$
\begin{aligned}
& \sum_{k \in \mathbb{Z}}\left\|\left[\sum_{i \in \Delta_{k}}\left(\frac{|\lambda(k, i)|}{\left\|\chi_{m \widetilde{Q}_{k}^{i}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}}\right)^{\underline{p}} \chi_{m \widetilde{Q}_{k}^{i}}\right]^{1 / \underline{p}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}^{q} \\
= & \sum_{k \in \mathbb{Z}} \|\left[\sum_{i \in \Delta_{k}}\left(\frac{\left\|\mathcal{U}_{\psi,(\mathrm{ii})} b(k, i)\right\|_{L^{r}\left(\mathbb{R}^{n}\right)}}{\left|m \widetilde{Q}_{k}^{i}\right|^{1 / r}}\right)^{\underline{p}} \chi_{m \widetilde{Q}_{k}^{i}}^{1 / \underline{p}}\left\|^{q}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}\right. \\
\lesssim & \sum_{k \in \mathbb{Z}}\left\|\left[\sum_{i \in \Delta_{k}} 2^{k \underline{p}} \chi_{m \widetilde{Q}_{k}^{i}}\right]^{1 / \underline{p}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}^{q} \\
\lesssim & \sum_{k \in \mathbb{Z}}\left\|\left[\sum_{i \in \Delta_{k}} 2^{k \underline{p}} \chi_{\widetilde{Q}_{k}^{i} \cap \Omega_{k}}\right]^{1 / \underline{p}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}^{q} \\
\lesssim & \sum_{k \in \mathbb{Z}}\left\|\left[2^{k \underline{p}} \chi_{\Omega_{k}}\right]^{1 / \underline{p}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}^{q} \\
\sim & \sum_{k \in \mathbb{Z}} 2^{k q}\left\|\chi_{\Omega_{k}}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}^{q} \\
\sim & \left\|\mathcal{U}_{\psi,(\mathrm{ii})} f\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}^{q}
\end{aligned}
$$

Hence, by Lemmas 2.12 and 2.15, we get

$$
\begin{align*}
& \|f\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} \\
\lesssim & \left.\left\{\sum_{k \in \mathbb{Z}}\left\|\left[\sum_{i \in \Delta_{k}}\left(\frac{|\lambda(k, i)|}{\left\|\chi_{m} \widetilde{Q}_{k}^{i}\right\|_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}}\right)^{\underline{p}} \chi_{m \widetilde{Q}_{k}^{i}}\right]^{1 / \underline{p}}\right\|^{q}\right\}_{L^{p(\cdot)}\left(\mathbb{R}^{n}\right)}\right\}^{\frac{1}{q}}  \tag{4.12}\\
\lesssim & \left\|\mathcal{U}_{\psi,(\mathrm{ii})} f\right\|_{L^{p(\cdot), q}\left(\mathbb{R}^{n}\right)} .
\end{align*}
$$

This implies (iii) $\Rightarrow$ (i).
Hence, (i)-(iv) are mutually equivalent.
Moreover, by (4.3), (4.4), (4.7), and (4.12), we find that each of $\|\cdot\|_{(\mathrm{i})}$, $\|\cdot\|_{(\text {ii) }}$, and $\|\cdot\|_{(\text {iii })}$ is equivalent to $\|\cdot\|_{H^{p(\cdot), q}\left(\mathbb{R}^{n}\right)}$. This completes the proof of Theorem 4.1.

From Theorem 4.1, we conclude the following result, see the proof of $[8$, Corollary 1.10] for more details.

Corollary 4.2. Replacing the assumption (4.1) in Theorem 4.1 by $f \in L^{2}\left(\mathbb{R}^{n}\right)$, then all the conclusions in Theorem 4.1 still hold true.

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Yao He
School of Mathematics and Statistics
Central South University
Changsha 410075, P. R. China
Email address: heyao1992@csu.edu.cn


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